WATER WAVE-STRUCTURE INTERACTION FOR SMALL AMPLITUDE STRUCTURAL OSCILLATIONS

A THESIS SUBMITTED FOR THE DEGREE OF DOCTOR OF PHILOSOPHY IN THE DEPARTMENT OF CIVIL ENGINEERING AT THE CITY UNIVERSITY

BY

NASHAT HAMED EL KADDAH BSc, MSc

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To My Parents

A B S T R A C T

The water wave/structure interaction is a complex phenomenon, which affects the prediction of the dynamic response of offshore structures. The problem arises from the interdependence of the structural response and the wave force. This interdependence is usually represented by the effect of the relative velocity and relative acceleration used in the wave force formula.

This work is carried out to investigate the effect of the water wave/structure interaction on circular cylinder for small Keulegan-Carpenter numbers and structural amplitude of oscillation. In order to achieve this an extensive experimental program, in conjunction with theoretical study was undertaken.

The study indicates that the use of modified Morison's equation taking into consideration the velocity and acceleration of the structure is not needed to predict the structural deflection. But significant difference in the force coefficients for free and fixed structure have obtained experimentally.

ACKNOWLEDGEMENTS

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NOTATION

ϕ	Velocity potential
η	Wave surface
P	Pressure
ρ	Water density
$\dot{U}$	Water particle velocity in x direction
v	Water particle velocity in y direction
w	Water particle velocity in z direction
g	Gravitational acceleration
H	Wave height
L	Wave length
h	Water depth
D	Structure diameter
F_m	Total force
F_D	Drag force
F_I	Inertia force
C_D	Drag coefficient
C_M	Inertia coefficient
$\ddot{U}$	Water particle acceleration
K-C	Keulegan and Carpenter number(Um T/D)
R_e	Reynolds number(Um D/γ)
C_L	Transverse (lift) coefficient
X	Structure displacement
$\dot{X}$	Structure velocity
$\ddot{X}$	Structure acceleration
[M]	Mass matrix

$[C]$	Damping matrix
$[K]$	Stiffness matrix
$\{F(t)\}$	Hydrodynamic force matrix
$\{X\}$	Structure displacement matrix
$\{\dot{X}\}$	Structure velocity matrix
$\{\ddot{X}\}$	Structure acceleration matrix
$[C_s]$	Damping in the free structure matrix
$[C_f]$	Damping in the foundation matrix
$[K_s]$	Stiffness in the free structure matrix
$[K_c]$	Stiffness in the foundation matrix
$[C_s^f]$	Damping matrix for fixed based structure
$[C_s^{bb}]$	Damping matrix associated with the base degree of freedom
$[C_s^b]$	The coupling damping matrix
$[M^f]$	Diagonal mass matrix including hydrodynamic added mass
$\{\phi_J\}$	Mode shape vector
W_J	Frequency of the J mode
ξ_J	Damping ratio of the J Mode
M_j	Generalized mass of the J mode
$[K_s^f]$	Stiffness matrix of fixed base structure
$[K_s^b]$	Stiffness-coupling matrix
M	Total mass
M_m	Material structural and platform mass
M_a	Hydraulics added mass
M_F	Flooding and marine growth mass
Ψ	Volume of water displaced by structure
ξ	Damping coefficient

C	Structural damping
C_c	Critical damping
σ_n	Wave angular frequency ($2\pi/T$)
k	Wave number ($2\pi/L$)
c_n	Velocity of wave propagation $\frac{\sigma_n}{k}$
a	Wave amplitude
L_s	Structure length
z	Vertical distance from the still water level
C'_D	Drag coefficient for free structure
C'_M	Inertia coefficient for free structure
F'	Force exerted on free structure
F_m	Measured force
S	Element length
$\{Pt\}$	Force matrix from the modified Morison's equation
$[\tilde{C}]$	Equivalent linear damping matrix
$[C^W]$	Damping matrix using the damping coefficient in water
$[M^*]$	Generalized mass matrix
$[K^*]$	Generalized stiffness matrix
$[C^*]$	Generalized damping matrix
$\{P^*\}$	Generalized force vector
$[\phi]$	Eigen vector
$\{\omega^2\}$	Eigen value
$[\ddot{r}]$	Normal co-ordinate structural displacement
$[\dot{r}]$	Normal co-ordinate structural velocity
$[\ddot{r}]$	Normal co-ordinate structural acceleration
$\{e(x)\}$	Error vector
$\alpha^2 e_n\ e_n$	Ensemble average

α_0	Proportional factor multiplied by the mass matrix
α_1	Proportional factor multiplied by the stiffness matrix
f_{NY}	Nyquist frequency ($1/2\Delta t$)
τ	Shear distribution
ψ	First mode shape of the structure
FN	Natural undamped frequency of vibration
FD	Natural damped frequency of vibration
W_n	Natural frequency of Nth mode $W_n = 2\pi FN$
E	Modulus of elasticity
I	Second moment of area
W	Axial load
C_{DP}	Percentage ratio of the drag coefficient for a free structure to that for fixed structure
C_{MP}	Percentage ratio of the inertia coefficient for a free structure to that for fixed structure
C_{DD}	Percentage ratio of the drag coefficient calculated by modified Morison's equation to that calculated by Morison's equation
C_{MD}	Percentage ratio of the inertia coefficient calculated by modified Morison's equation to that calculated by Morison's equation
VD	Reduced velocity (U_m/D FN)
U_m	Maximum water partical velocity
$\bar{M}$	Effective mass
DM	Response parameter ($2\bar{M}\xi/\rho D^2 L_s$)CD FN/U_m)

CHAPTER ONE

INTRODUCTION

At present, the exploration for energy resources has resulted in offshore structures being built in deep seas, where severe conditions of wind and waves exist; under these conditions the classical design process produces uneconomical and sometimes unsafe structures.

It has been shown in the inquiry into the collapse of Texas Tower No 4[1] that dynamic analysis will have to be used for designing of such towers.

In deep water, the interaction of a time-dependent ocean environment with a dynamically responsive structure leads to complex resonance conditions and gives rise to larger stresses than would be predicted by a quasi-static analysis. For example, Brannon et al 1974[2] have shown that the dynamic response can double the static wave load.

Laird 1962 and 1966[3, 4] has shown that the drag and inertia coefficients may vary rather widely when the structures are oscillating.

From the above it can be concluded that the investigation of the water wave/structure interaction is important for both the dynamic analysis and for the determination of the water wave force.

The existing knowledge in the fields of mechanical vibration and fluid mechanics can be used to determine the structural response. By knowing the system input of the wave profile and by applying the hydrodynamic theories, such as Airy's linear wave theory, the flow characteristics of the wave in the vicinity of the structure can be evaluated. Morison's semi-empirical force formula enables the hydrodynamic forces on the structure to be evaluated. The theory of mechanical vibration is used to compute the structural response due to these forces. The complication arises here due to the water wave/ structure interaction which is the interdependence of the fluid forces and structural response.

In general the problem of the interdependence of the fluid forces and structural response is caused by the structures significant velocity and acceleration only. This may be overcome by using the relative velocity and relative acceleration in the Morison's equation which is called the modified Morison's equation.

In this thesis an investigation of this interdependence, including the effect of the structure's vibration on the force coefficients (C_D and C_M) used in the force formula and the effect of this interdependence in the dynamic response (deflection), was carried out for the condition of low Keulegan-Carpenter numbers (1.87 - 10.69) and amplitudes of structure oscillation.

For this purpose extensive experimental work was done. Two sets of structures in different material were tested. The structures were circular cylindrical piles. The first set was PVC material and consisted of two groups. Each group had two structures with different heights. The first group was relatively stiff (the structure diameter was 0.11 metre), the second group was relatively flexible (the structure diameter was 0.0605 metre). The second set was aluminium and consisted of two structures of the same diameter (0.0574 metre), but with different height; in this set each structure was tested three times, once with no added load on top and later with different added loads on top, to give different flexibility properties of the sturcture. The dynamic properties of the tested structures, such as the damping coefficient, the fundamental natural frequency and the stiffness constant, were determined experimentally.

The above structures were tested for different wave conditions. At each wave condition and at each level of the structure where the pressure can be measured, the structure was tested twice, once when it was free to vibrate and once when it was prevented from vibrating. At each test, the wave profile, the pressure distribution around a certain level of the structure, the tip displacement and the bending moment at the base of the structure were measured.

The data obtained from the experiments were subjected to extensive analysis in order to predict the most appropriate wave theory for the tested wave condition and the method used for the determination of the average value of the force coefficients C_D and C_M.

For the determination of the wave velocity and acceleration, the fast Fourier's analysis was utilized to obtain the free higher harmonic components in the tested wave. The least square method was used in the determination of the average value of the C_D and C_M.

For the analysis of the structural response, the structure system was idealized. The equivalent structure was assumed to be fixed at the base and the continuous member was assumed to consist of a series of pipe finite elements with discrete springs

of mass at the two end nodes. The vertical wave force was ignored and the horizontal wave force was discretized at the node point. The analysis was restricted to two dimensions.

Two methods were used in the formulation of the dynamic equation of structural response in order to find whether the water wave/structure interaction represents a significant contribution in the prediction for the case of the condition studied.

The first method included the water wave/structure interaction which will cause a nonlinear term in the partial differential equation of motion, the direct approach method was used to overcome this nonlinearity.

In the second method, the water wave/structure interaction was ignored by assuming that the wave force is not modified by the motion of the structure and substituted for this assumption by using the value of the hydrodynamic damping obtained in still water for the dynamic analysis.

The numerical analyses were done by finite element method using computer program SAP IV.

The dynamic response was calculated by the mode superposition technique taking only the first three modes and using time domain description.

This work gives a clear insight into the relative influence of the structure vibration on the fluid force coefficients, (the inertia and drag coefficient). It also gives the possibility of substituting for the nonlinear hydrodynamic damping that arise from the terms of relative fluid-structure velocities by the damping coefficient in still water for obtaining the structural response.

CHAPTER TWO

SURVEY OF LITERATURE RELATED TO DYNAMIC RESPONSE

The flow induced vibrations are a relatively complex and diverse phenomena and heterogeneous body of research and analyses in related fields is available but is scattered through a wide variety of journals.

One of these phenomena is the water wave/structure interaction. The term, water wave/structure interaction, however, denotes only these phenomena in which an interdependence develops between the fluid dynamic forces acting on the structure and the structural response.

The inclusion of the dynamic response of an offshore structure to the water wave force complicates the determination of the structure's behaviour and water wave force.

In order to study the behaviour of these structures under dynamic loading and wave conditions, the following aspects of the problem will have to be understood and properly applied to the structure in question.

(1) The hydrodynamic forces on the structure caused by waves
(2) The idealization used to model the structure's system should simulate the system's response.

(3) The applicable numerical procedure must be used.

2.1 HYDRODYNAMIC FORCES

The determination of the hydrodynamic forces exerted by waves on structure is complex, even for slender members because of the assumptions and approximations required to predict the wave force. This requires the determination of the time histories of the loading corresponding to the design wave (i.e. velocity and acceleration time histories of wave particular with the region of each structure).

2.1.1 WATER WAVE THEORIES

If the flow field is assumed incompressible and irrotational, the flow field should satisfy the Laplace's equation

$$\nabla^2\phi = 0 \quad \ldots \ldots \ldots \ldots \ldots \ldots \quad (2.1.1)$$

where ϕ is the velocity potential.

The flow field will have to satisfy the following boundary conditions

- A - at the surface

let $\eta(x, y, z, t)$ describe the surface S

$\therefore$ the kinematic boundary $\dot{u}\eta_x + V\eta_y + W\eta_z = -\eta_t$

$$\ldots \ldots \ldots \ldots \ldots \ldots \quad (2.1.2)$$

free surface boundary $\frac{P}{\rho} = 0 \quad \ldots \ldots \ldots \ldots \quad (2.1.3)$

where P/ρ = Pressure at the free surface

then the generalized Bernoalli equation reduce to

$$\frac{\partial\phi}{\partial t} + \frac{1}{2}(\dot{u}^2 + v^2 + w^2) + gY = F(t) \quad \ldots \ldots \quad (2.1.4)$$

- B - at the bottom

Let G(x, y, z, t) represent the bottom surface

$$\therefore \dot{U}G_x + VG_y + WG_z = - G_t \quad \ldots \quad \ldots \quad \ldots \quad (2.1.5)$$

where the suffix mean differentiation with respect to the subscript.

All wave theories find close solution to the above equations by using certain assumptions to match with the physical condition of the wave (5, 3, 7).

The wave in the sea can be classified periodic, aperiodic and translatory. The regions of applicability of the various wave theories depend on the following parameters:

(a) The wave parameters, wave height, wave length and period (H, L, T)

(b) The position parameter water depth (h).

The limitation conditions for sinusoidal and stakes wave were shown by Laitone 1962[8]. Wilson 1957[9], Dean[10,11,12] and more recently Nishimuru and Isobe 1978[13] have discussed the validity of these various theories. Figure (2.1.1) shows the relation between the validities of the theories to evaluate the kinematic properties for water wave with the relation to wave and position parameter after Dean[12], also Figure (2.1.2) after Wilson.

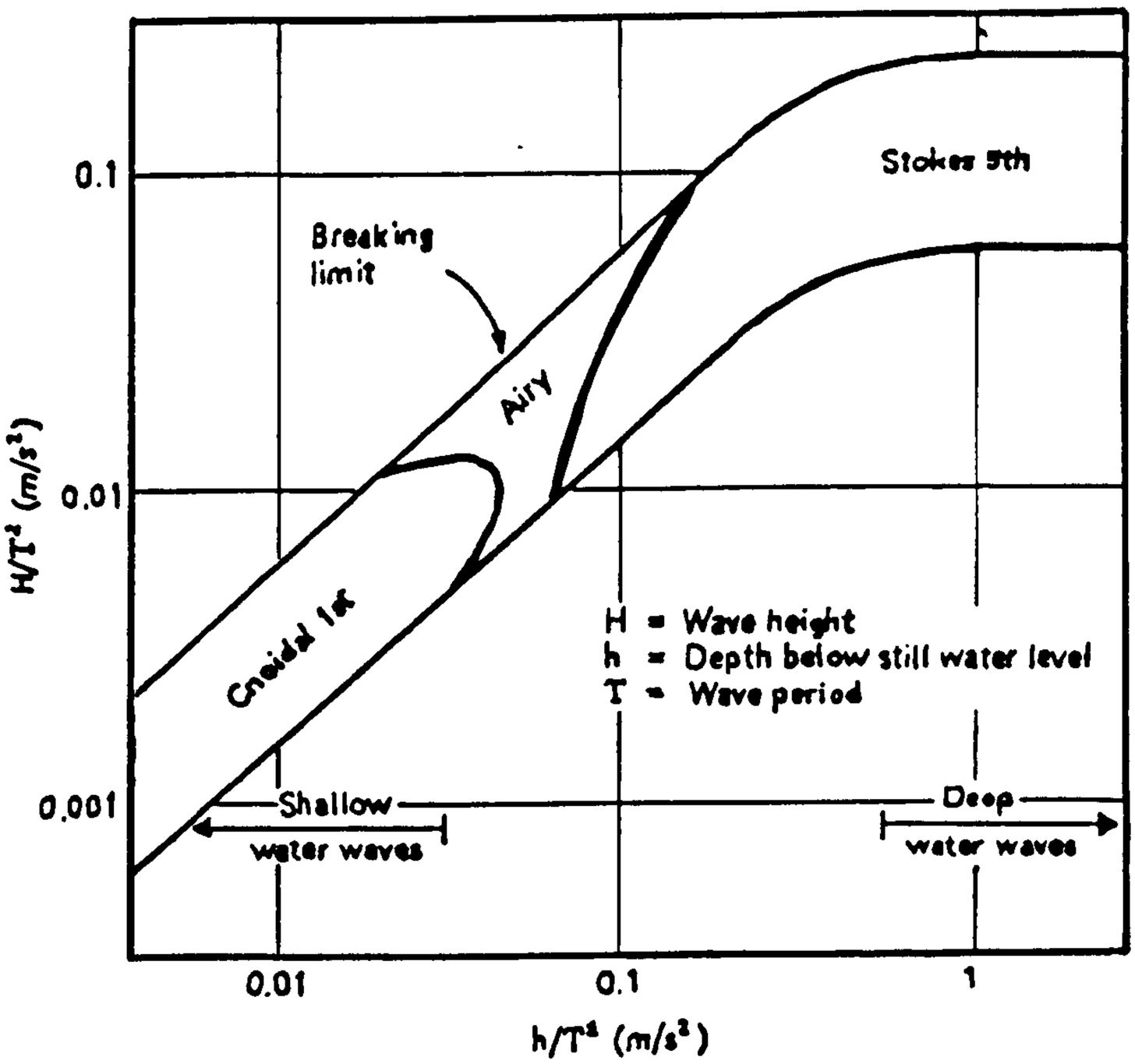


Figure (2.1.1) - The Range of Validity of Various Wave Theories.

In irregular random waves, the statistical properties will have to be determined. The random wave spectrum consist of waves having different wave-lengths and frequencies, Wiegel 1964[14] and more recently 1978[15] used several linear and non-linear representations of the ocean system.

The wave theory for the appropriate wave region is then chosen. From this, velocity and accelerations are calculated and related to the force.

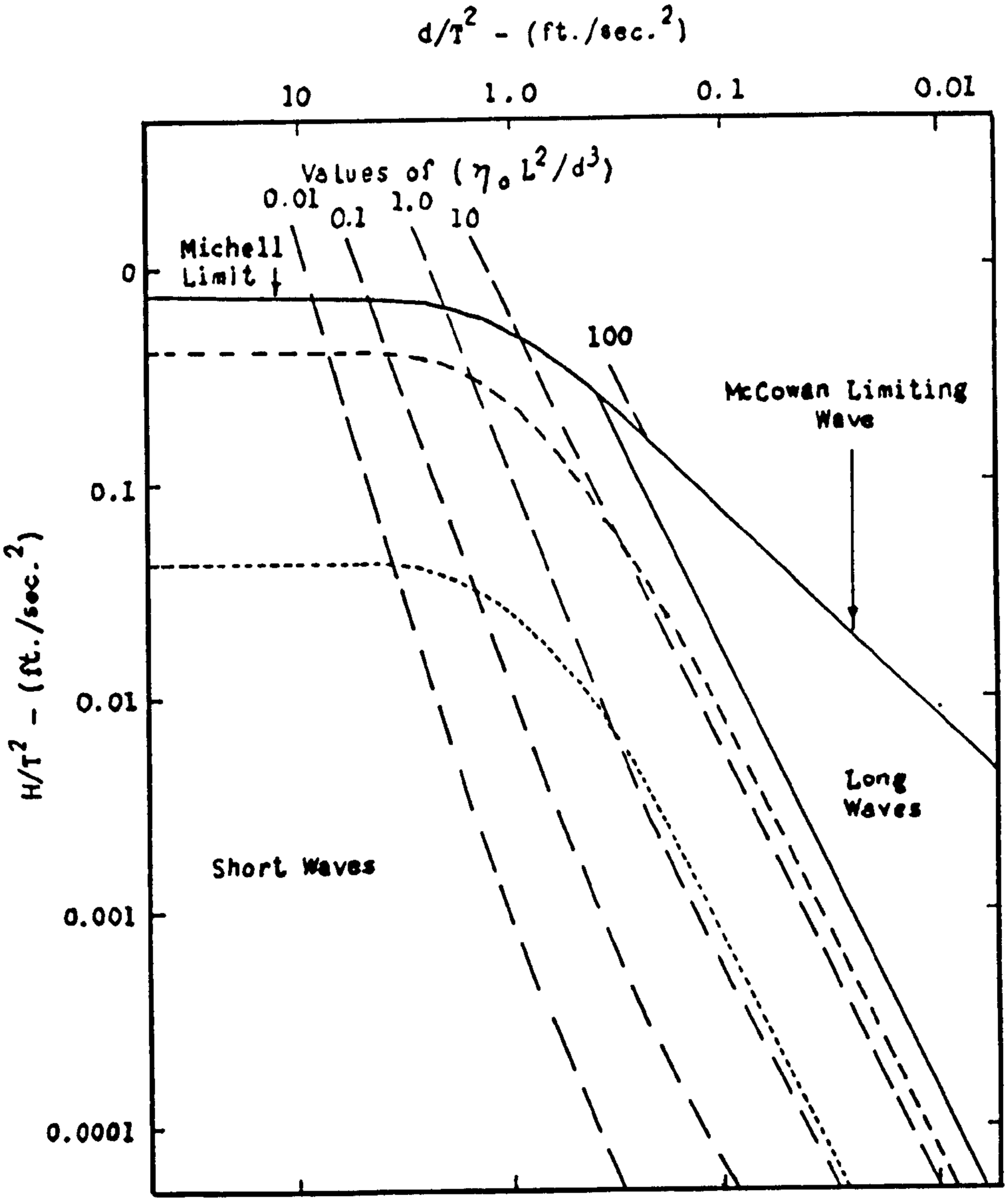


Figure (2.1.2) - Representation of Zones of Validity of Various Wave Theories.

2.1.2 HYDRODYNAMIC FORCE

Several approaches have been developed for evaluation of the wave forces on offshore structures. The applicability of these methods depends on the relative magnitudes of the typical dimension of the structure, D, with respect to the wave-length, L, and wave height, H. The dimensionless ratios H/2D (the wake parameter) and $2\pi D/L$ (the scattering parameter) are key parameters which dictate the choice of the method used for wave force evaluation on offshore structure. See Figure (2.1.3) by Garrison and Rao 1977[16].

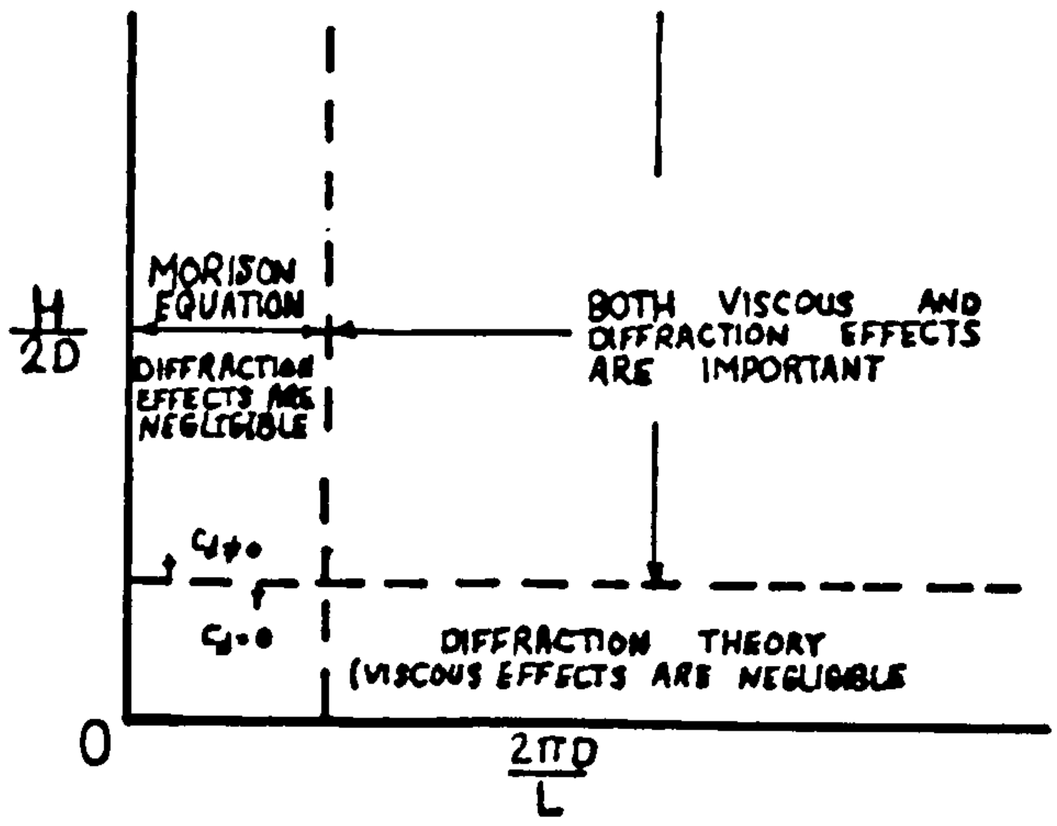


Figure (2.1.3) - Wave Force Theory.

The scattering parameter $2\pi D/L$ controls the ratio between the incident wave and the reflected wave. For large value of $2\pi D/L$ the structure is largely in a deflection regime. For smaller value of $2\pi D/L$ the structure will be in the drag regimes where the incident wave plays a predominant part.

The following ratios of D/L give the approximate regime of flow

$D/L > 1$ pure reflection

$D/L \geqslant 0{\cdot}2$ diffraction effects predominate

$D/L < 0{\cdot}2$ incident wave predominate

The wake parameter H/2D controls the flow regime of the wave round the body as H/2D becomes sufficiently large, flow separation and energy dissipation in eddys are important, where the viscous effects become important and the following ratios of D/H give approximate regimes, Verley[17] 1975.

$D/H > 0{\cdot}3$ inertia increasingly predominant

$D/H = 0{\cdot}6$ incipience of lift (and drag)

$D/H < 0{\cdot}3$ drag increasingly predominant

2.1.2.1 FORCE ON LARGE BODIES

2.1.2.1A DIFFRACTION THEORY

For the case of diffraction where viscous effects are not important, MacCamy 1954[18] calculated the forces on cylinder subjected to regular wave using linear wave theory and Bessel function.

Spring and Monkmeyer 1974[19] outlined an analytical procedure of solution for any number of cylinders (Linear Theory) and obtained results for the case of two cylinders. Chakrabarti 1978[20] applied with slight modification, Spring and Monkmeyer (1974) method to the case of more than two cylinders.

Garrison 1974, 1978[21-22] presented the linear wave theory based on Green's function, for general structural geometries. Raman et al 1975[23] used a higher order theory to obtain analytical studies for circular cylinder. They used the perturbation method to obtain a solution for the second order theory. This resulted in a slight improvement when compared with the linear theory.

For the case of diffraction where viscous and diffraction effects were important, for the case of an isolated cylinder, Chakrabarti 1973[24], used a fifth order waves theory. However there is an inconsistency in his solution as the fifth order wave theory was used without satisfying the free surface non-linear condition.

2.1.2.2 FORCE ON SMALL BODIES

For small values of the scattering parameter "2πD/L" and the wake parameter "H/2D" where Morison's equation and diffraction theory both apply, the problem can be treated as radiated waves.

2.1.2.2A MORISON'S EQUATION

As mentioned above that when scattering parameter is small Morison's equation is applied to evaluate the wave force, this equation has been formulated by Morison et al 1950[25] for the vertical cylinder The main assumptions were,

(a) the body has negligible effect on the waves (the wave field does not change due to the presence of the structure),

(b) the total force on the body has two independent components.

These two component of forces are:

The drag force which results from the flow separation induced by the relative velocity of the fluid around the structure, the drag force on the body is generally made up of two terms, one a viscous and the other a pressure term, In the vast majority of cases the viscous term is negligible compared to the pressure as the viscous drag is proportional to velocity and not to the velocity squared as for the pressure drag force.

The other component is the inertia force which is due to the pressure gradient associated with the relative acceleration of the ambient fluid and is proportional to the particle acceleration.

Therefore the total force on a vertical cylindrical structure is assumed to be

$$F_T = F_D + F_I \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (2.1.6)$$

Each of these time varying components has been formulated in terms of (i) Geometrical properties of the structure, (ii) fluid properties describing the flow field and (iii) some "variable coefficients"

$$\therefore \quad d\,F_D(t) = C_D\, 1/2\, \rho D\, \dot{U}(t)|\dot{U}(t)|\, ds \quad \ldots \quad (2.1.7)$$

$$d\,F_I(t) = C_M \frac{\rho \pi D^2}{4} \ddot{U}(t)\, ds \quad \ldots \quad \ldots \quad (2.1.8)$$

where D = diameter of the cylinder

ρ = water density

$\dot{U}$ = fluid particle velocity

$\ddot{U}$ = fluid particle acceleration

C_D = drag coefficient

C_M = inertia coefficient

In the above equations the geometrical properties of the structure appear as exposed area in the drag force term and as volume in the inertia term. But this geometry of the structure also affects the fluid flow field around the structure as shown in Blesvins 1977[26].

The fluid properties represented in the wave velocity and acceleration in the above equations are calculated from the wave theories. Hence it is important to use the sequential wave theory in order to estimate accurately the wave force. They also affect indirectly the estimation of C_D and C_M as shown below.

The variable coefficients are the drag coefficient and the inertia coefficient. The drag force on structure is due to the pressure difference across the structure. This pressure difference is due to separation of the flow from some point on the structure creating a low pressure area behind the structure. The position at which separation occurs, the way in which it occurs and the resultant wake width all influence the pressure behind the structure. Thus the drag coefficient is used taking into account the unknown pressure difference across the structure.

The inertia force on the body is due to the fluid in the wave accelerating and as the fluid in the wave must move round the object additional accelerations are involved, related to the curvature of flow. The accelerations are dependent on the form of the wake. These accelerations induce a force field which is expressed in terms of the fluid mass it displaces.

It is clearly shown that the drag and the mass (inertia) coefficient is dependent on the geometrical properties of the structure and the fluid properties describing the flow field.

The experimental values of C_d and C_m obtained by Morison et al show some scattering and no trend as function of dimensionless parameter, Reynolds number (U_{max} D/ν), where "U_{max} is the maximum surface orbital velocity, D is the structure diameter and ν is the kinematic viscosity of the fluid.

Because the wave flow reverses every half wave cycle as shown in Figure (2.1.4). This means that there is only a relatively short time for the wake to form before it is destroyed. Keulegan and Carpenter 1968[27] formed a dimensionless parameter which compares the time for a wake to form with the time available for it to form.

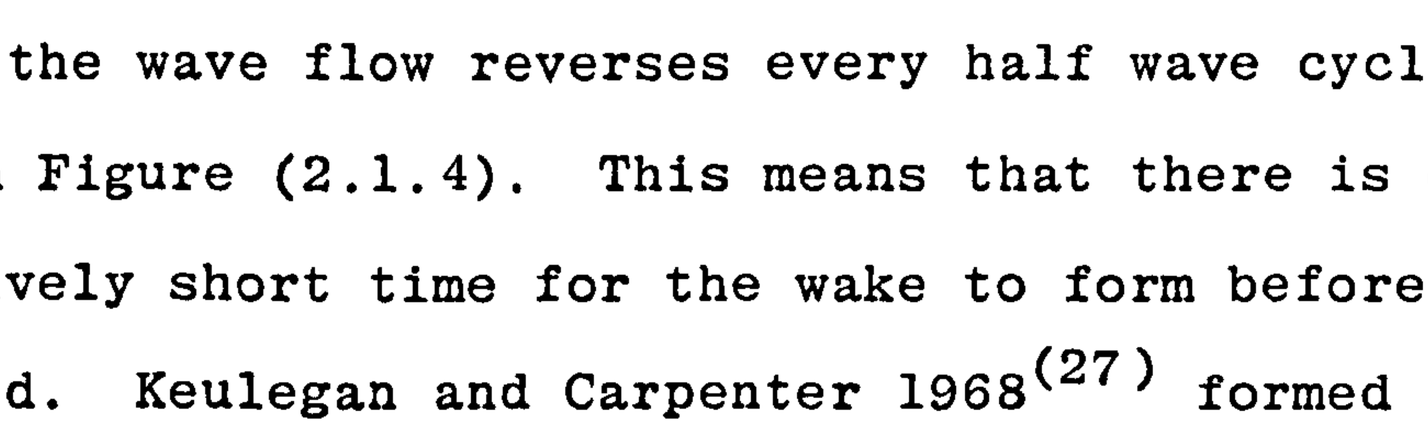


Figure (2.1.4) - Direction of Velocity and Acceleration at Various Points of a Wave Cycle.

They carried out studies both theoretically and experimentally on the force acting on cylinders in an oscillating flow field. They observed that C_D and C_m varied over a wave cycle. They correlated the cycle averaged values of C_D and C_m with the period parameter K-C = U_m T/D called the Keulegan-Carpenter number where "U_m is the maximum orbital velocity, T is the wave period and D is the diameter of the cylinder."

The variation of the average of the inertia coefficient values per cycle with the period parameter is shown in Figure (2.1.5). Similarly for the drag coefficient is shown in Figure (2.1.6).

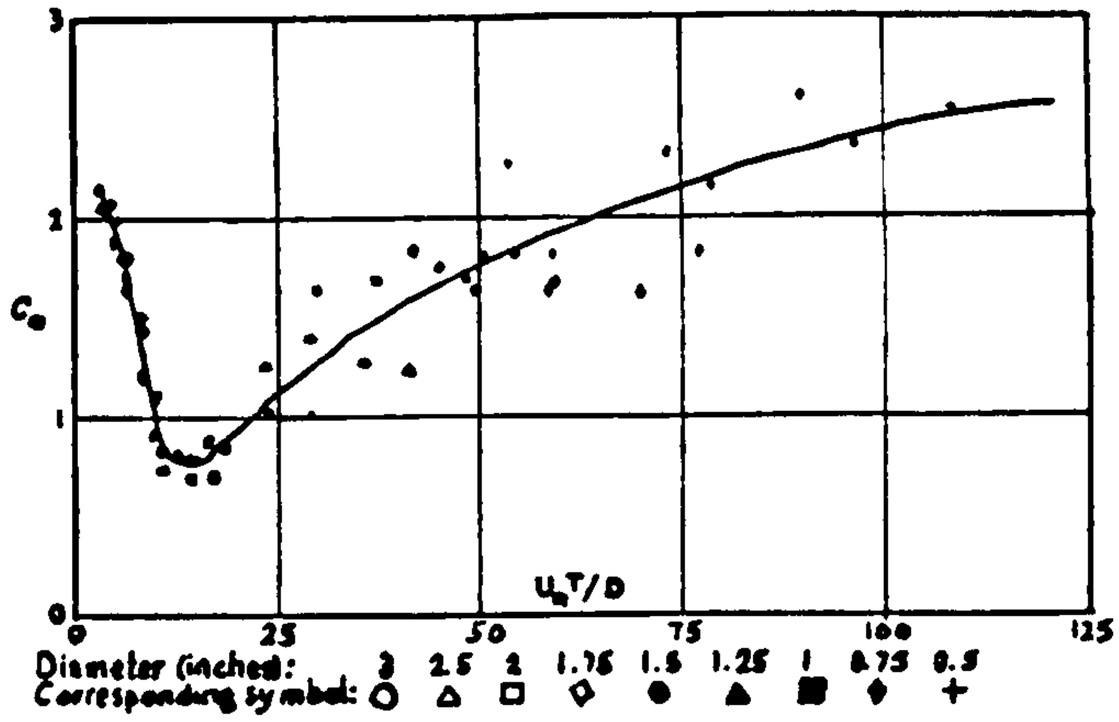


Figure (2.1.5) - C_M Versus Keulegan-Carpenter Number.

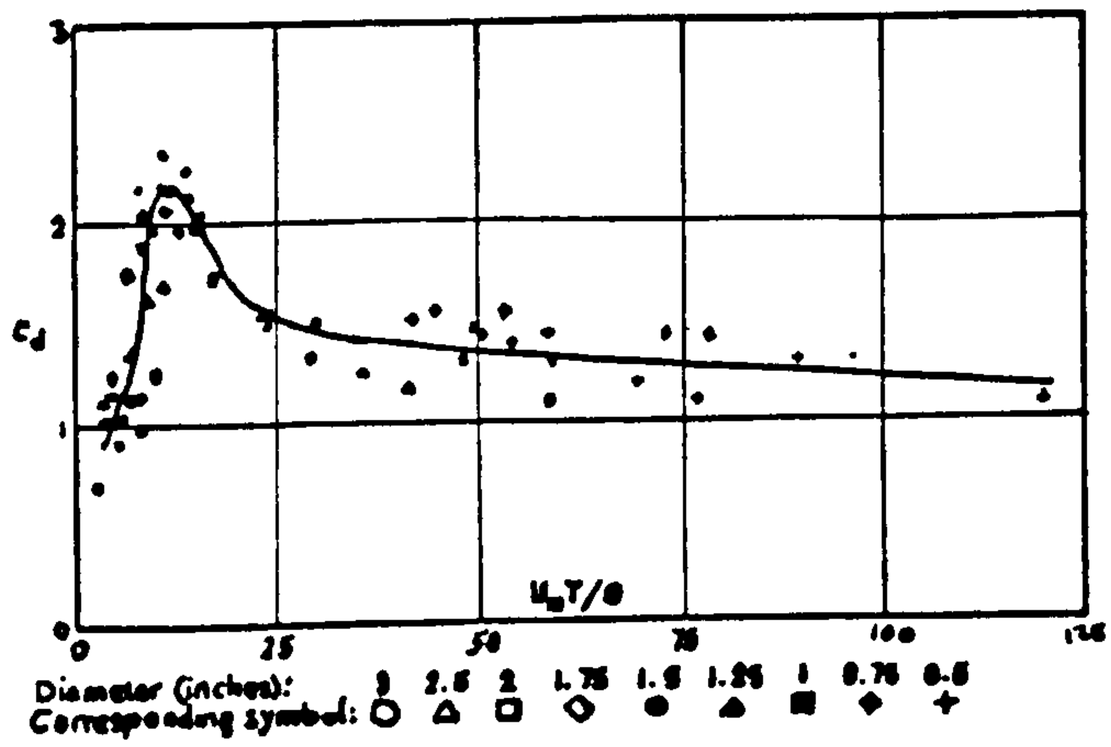


Figure (2.1.6) - C_M Versus Keulegan-Carpenter Number.

The period parameter K-C is associated with the flow separation process and eddy formation around the cylinder. When K-C is small, no separation occurs. As K-C increase, separation is initiated and eddies are formed.

In the work of Keulegan and Carpenter they did not correlate C_D and C_M with respect to Reynolds number R_e.

In order for Morison's equation to be applicable and be useful to engineers in the design of offshore structure, it was necessary to undertake a comprehensive laboratory and field tests for the purpose of determining the coefficient of drag and inertia which appear in the equation. Some of the data came from the measurements carried out in the actual sea condition (field test).

Wiegel, Beebe and Moon 1957[28] made field measurements at Pacific Coast (Davenport, California) on various sections of a 6·625 inch cylinder. They used linear wave theory, average values of C_D and C_M were obtained. They calculated C_M at zero velocity and C_D at zero acceleration, the results show considerable variations in the mean values of C_D and C_M as well as scatter.

Agerschau and Edens 1965[29] used the data obtained from Wiegel et al and used Stokes fifth-order theory. The values of C_D and C_M obtained were also scattered and show that the fifth order approach was not superior to the first order "linear theory."

Reid 1957[30] measured force on a section of an 8·625 inch cylinder in water of 30 ft depth in the Gulf of Mexico. The kinematics of the flow were calculated from the wave profile and the drag and inertia coefficient were obtained through the use of the least squares technique. Measured wave force records and those calculated using constant mean C_D and C_M values were in good agreement. Structural vibration was observed and allowed for in the analysis.

Wilson 1965[31] presented the results of wave force data from an experiment conducted with a 30 inch diameter pile in confused sea conditions in the Gulf of Mexico. He developed a numerical filter with which unwanted high frequency effects were removed from force records. Large scatter in C_D and C_M values from different wave force record analyses.

Two wave projects were undertaken at the Gulf of Mexico and the results were evaluated by several researchers; Aagaard and Dean 1970[32] by using stream function theory in the estimation of C_D and C_M with C_M = 1·33 and C_D varying with Reynolds number as obtained from the data, they found agreement between calculated and measured local force maximum within 50%.

Evans 1970[33] performed his work in 100 ft depth of water with 3·7 ft diameter pile. The average pressure distribution was measured over a one foot section. The forces and their directions were calculated from it. The wave profile was recorded simultaneously with pressure. He used Stokes fifth order wave theory to calculate C_D and C_M. Calculated total forces were generally within 10% of the measured forces, and usually conservative. None of the evaluations of the data of the two projects considered the effect of currents.

Wheeler 1970[34] used another technique to represent the data of the Gulf of Mexico by predicting the velocities and accelerationsby linear digital filter acting on the wave profile records. No correlation of C_M and C_D coefficient with critical parameter was reported; agreement between measured and calculated maximum local forces was within 40%.

Kim and Hibbard 1975[35] presented the results of data obtained from measurements in Australia. The test pile was 38 ft long and 12·75 inch in diameter, it was subjected to rather small amplitude waves, the agreement between the measured and calculated force was good in the drag dominated part of the wave cycle and fair in the inertia dominated region.

Heideman, Olsen and Johansson 1979[36] used two methods to evaluate the drag and inertia coefficients from the large scale experiments of space frame structure in the Gulf of Mexico. The first was the least squared error procedure for each half wave cycle. The second method consisted of evaluation of C_D over short segments of wave where drag force was dominant and of C_M over short segments in which inertia force was dominant. The force coefficients exhibited large scatter particularly for K-C < 20. The current was taken into consideration.

Bretschneider 1967[37] predicted the probability distributions of peak wave drag and inertia force from field tests carried out in California Coast at Devenport for vertical cylinder using linear wave theory, the probabilistic approach considering peak drag and inertia forces were unsatisfactory.

The reasons for discrepancies in the evaluation of C_D and C_M in field tests may be many; among the more important are:

1. Representation of the irregular wave

2. Turbulence around the structure by which the test structures are supported

3. Inaccuracy in measuring force

4. Vibration of test structure

5. Inability of wave theories to describe actual water particle motion especially if there is the possibility of a steady ocean current.

In order to eliminate most of the above discrepancies, to obtain for a wide range of data and flow conditions and to segregate each parameter which can affect the calculation of C_D and C_M, laboratory experimental tests were conducted. In the laboratory tests several experimental methods have been used in the evaluation of the drag and inertia coefficient which include:

(a) Force measured on structure in laboratory wave

(b) Force measured on structure in an oscillating fluid, where the motion of the particle is in straight line rather than orbital

(c) Force measured on a cylinder constrained to move with an oscillatory motion in a stationary fluid.

The flow field characteristics of (a) are essentially different from that of (b) and (c).

Some of the experimental work which follow pattern (a) were carried out by Morison et al 1950[25] using measured moments and using linear wave theory, in calculating C_M and C_D they did not correlate well with d/λ, D/λ or R_e for linear sinusoidal waves.

Susbielles et al 1971[38] used various methods of derivation of coefficients which included linear, stokes third and fifth order wave theory and stream function wave theory. The actual C_M and C_D range of values varied with the method of derivation. Several C_M and C_D pairs were shown to predict the same force. By using the values of the coefficients from Keulegan and Carpenter's curve the results obtained differed by 10% between the measured and the calculated.

Chakrabarti et al 1975[39] found large scatter in C_D and C_M values using Morison's equation in three-dimensional vector form with component normal to the axis of the cylinder. Measured and calculated mean forces agreed to within 10% using C_D and C_M values for individual waves. The tested structures were held at various inclinations in line with and normal to the waves.

Gaston and Ohmart 1979[40] measured the total wave force and overturning moment on a smooth and roughened 14 ft long, 1ft diameter, vertical cylinder under conditions of periodic and random waves. Drag and

inertia coefficients have been determined by the least squares method, using the measured in-line moment and predicted kinematics from the irregular stream function theory. Typical Reynolds number of the experiments were 2×10^5 to 3×10^5 based on the r.m.s water-particle velocity data. The results gave the average values of C_D and C_M as:

$$C_D = 0{\cdot}77 \text{ and } C_M = 1{\cdot}81$$

for smooth cylinder.

Experiments which seem to have shed the most light upon the fundamental behaviour of C_D and C_M are those which have been conducted in straight line oscillatory motion. The best known early experiments were those of Keulegan and Carpenter 1958[27] who utilized the linear oscillatory flow field under a node of standing wave. Their experiments showed that C_D and C_M are not constant throughout a motion cycle but, in fact, varied substantially with motion phase. More recently in a series of papers Sarpkaya et al have presented the results of an extensive series of oscillating fluid experiments using vertical U-shaped water tank. Sarpkaya (1976a, 1976b)[41-42] conducted a series of experiments for smooth and sand-roughened cylinders to evaluate the drag and inertia coefficients introducing "frequency parameter " β in which $\beta = D^2/\nu T$ or β is equal to R_e/K-C parameter to represent the data of C_D and C_M. The dependence of C_D and C_M on β has already

been noted in connection with the discussion of the Stokes 1851[43] sphere problem. Sarpkaya 1977a[44] carried out his series of experiments for high Renolds numbers. Sarpkaya et al 1977[45] expressed their results for C_D and C_M as function of Reynolds number, Keulegan-Carpenter number and relative roughness. For large Keulegan-Carpenter number, they obtained a different trend than Keulegan-Carpenter[27]. In general the scatter in the data is small.

Maull and Milliner 1978[46] examined the production and motion of vortices in a sinusoidally-oscillating flow in a small U-tube at Reynolds numbers smaller than about 4,000. They proposed that the variation of the drag coefficient during a cycle may be considered as the addition of two terms, the inertia term with $C_M = 2.0$, and a further term which is a function of the movement of the vortices produced.

Bearman and Graham 1979[47] measured the in-line force on several cylindrical bodies in plane oscillatory flow in a small U-tube over a range of K-C number from 3 - 70 at relatively small Reynolds number. They had noted large cycle to cycle variations in computed values of C_D and C_M even though the bulk flow in the U-tube was closely repetitive.

Force measured on a structure constrained to move with an oscillatory motion in a stationary fluid (Pattern C) help to add insight into the role played time dependence.

Most of the work of unidirectional acceleration of a body in stationary fluid conducted to establish a single force coefficient combines the effects of drag and inertia. Kiem 1956[48] showed that the single coefficient as a function of $(d\dot{U}/dt)D/\dot{U}^2$. This correlated the data fairly well. Also the experimental work of Laird and Johnson 1956[49] and Laird et al 1959[50] expressed the result in terms of total resistance coefficient.

Dalton et al 1976[51] and Sarpkaya and Garrison 1963[52] tried to use dimensional analysis to show that the single coefficient can be used to transform the results from the oscillating cylinder into the problem of a stationary cylinder in an oscillating fluid otherwise from Pattern C to Pattern A and B.

From the above we can see the differences in the test condition, methods of measurement and data evaluation do not permit a critical and comparative assessment of the drag and inertia coefficients obtained in each investigation. A comprehensive summary of the data on force-transfer coefficients has been presented by Hogben et al 1977[53]. The relationship between C_D and C_M has shown that there is not a unique relationship between them, independent of K - C and R_e. Hogben 1976[54] suggested

a conceptual modelling of the interaction leading to an explicit formula for the inter dependence of C_D and C_M, it is based on highly simplified and somewhat initiative reasoning.

2.1.2.2B TRANSVERSE FORCE (LIFT FORCE)

Although the transverse force is not concerned in these studies it is important to mention it because this phenomenon may cause increase in in-line force and it has an important effect closely related to drag forces. The transverse force can be calculated from the equation:

$$dF_L(t) = D_L \tfrac{1}{2} \rho D \dot{U}_m^2 \, ds \quad \ldots \quad \ldots \quad \ldots \quad (2.1.9)$$

where C_L = transverse or (Lift) coefficient. The rest of the symbols have been identified previously.

The phenomenon of the transverse force is caused by the eddy shedding and so it is likely to be influenced by the way in which the eddies are shed (this also applied to the drag force). Some of the work concern the transverse force in wave flow have been looked at by Bidde 1971[55], Wiegel and Delmonte 1972[56], Isaacson 1974[57] and Zedan and Rajabi 1981[58]. Some of the work concerning the transverse force in harmonically oscillating flows has been measured by Mercier 1973[59], Sarpkaya 1975a, 1976a[60, 41] and Maull and Milliner 1978, 1979[46, 61].

2.1.2.3 MORISON EQUATION FOR A FLEXIBLE CYLINDER

The original Morison equation has been modified by several investigations to cover in some sense the effect on the in-line forces of the flexibility of the structure. In particular the absolute value of the fluid velocity has been replaced by the relative one with respect to the structure's velocity, and the added mass term associated with the acceleration of the structure has been included (see Berge and Penzien 1976[62]; Moan, Haver and Vinje 1975[63]).

The modified relation is:

$$dF(t) = ((C_M-1)\frac{\pi D^2}{4}\rho(\ddot{U}(t) - \ddot{X}(t)) + \frac{\rho\pi D^2}{4}\ddot{U} + \frac{1}{2}C_D\rho D|\dot{U}(t) - \dot{X}(t)|(\dot{U}(t)- \dot{X}(t))\,ds \quad \ldots\ \ldots\ \ldots\ \ldots \quad (2.1.10)$$

where $\ddot{X}$ = The acceleration of the structure at the point under consideration

$\dot{X}$ = The velocity of the structure at the point under consideration.

The rest of the symbols have been identified previously.

Although Equation (2.1.10) is a reasonable extension of the original Morison equation, it requires experimental verification as the original form of Morison's equation has been used for the estimation of C_D and C_M in field studies.

2.2 THE IDEALIZATION USED TO MODEL THE STRUCTURE SYSTEM

In order to solve the static or dynamic response of any structure to the existing force it is necessary to model structure-foundation accurately to be able to establish the structure mechanism behaviour and the interaction between the superstructure and the supporting foundation.

As the offshore structure is a very complex structure and problem; Penzein and Tseng 1976[64] refer to separate the modelling of the structure foundation system into (a) a foundation system, (b) a structure system.

2.2.1 FOUNDATION SYSTEM

The foundation of offshore structure can be either raft foundation as those commonly used for the gravity type of structure; see Figure (2.2.1) or a piles foundation as those normally used for the framed steel type of structure, see Figure (2.2.2).

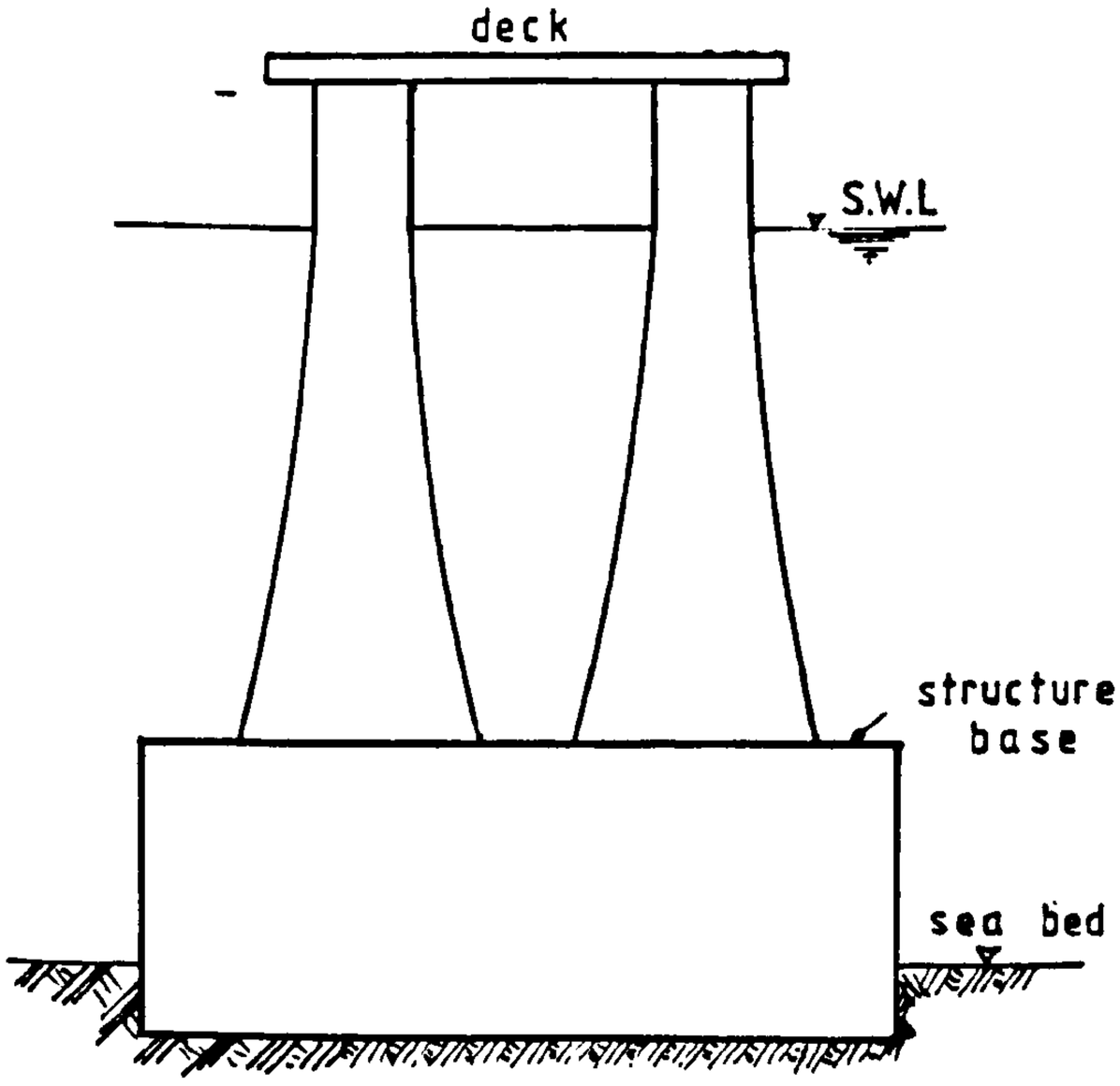


Figure (2.2.1) - Gravity Structure

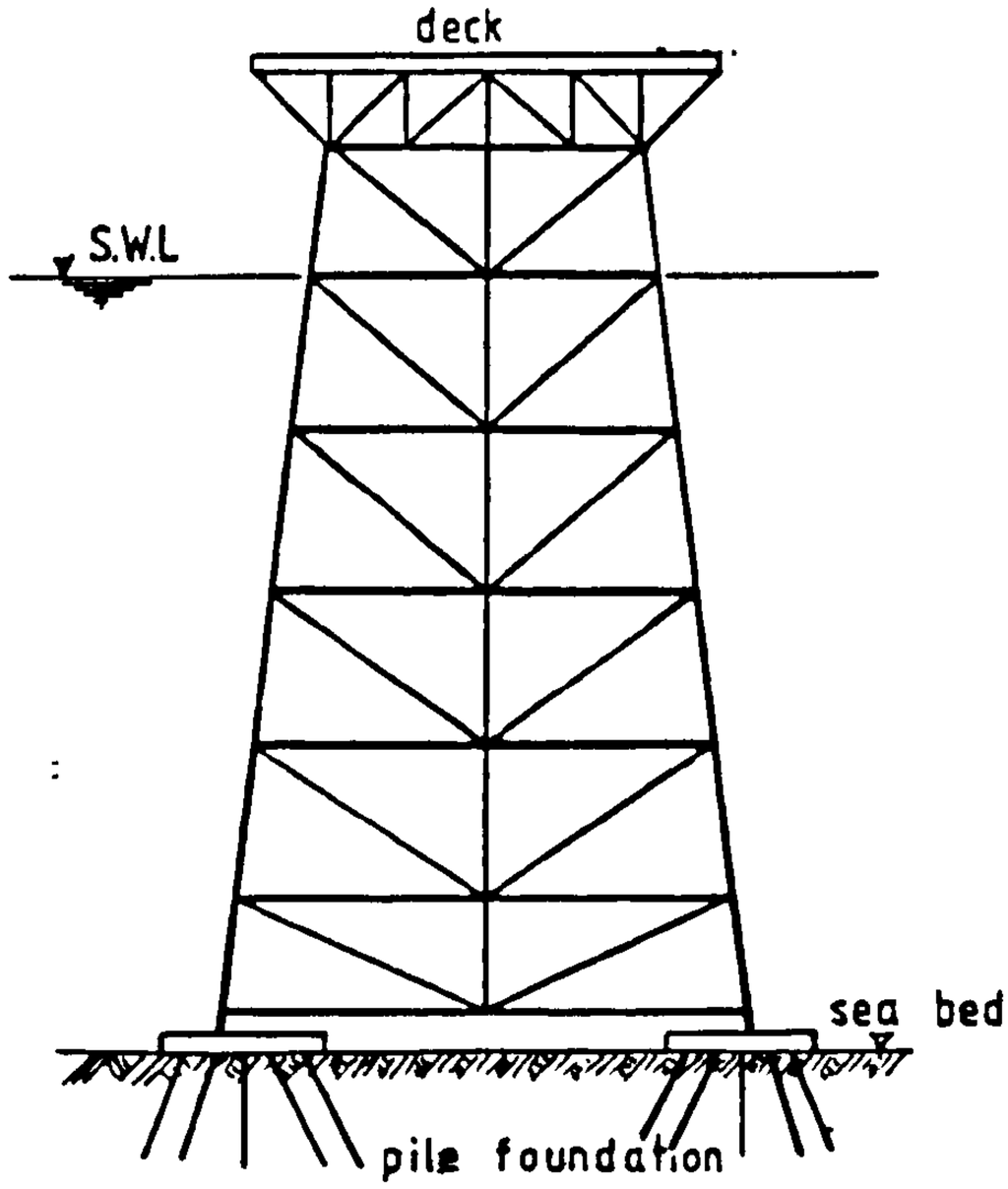


Figure (2.2.2) - Framed Steel Structure

The foundation impedence function stiffness coefficient and the damping coefficient of the foundation must be derived by performing a steady state foundation sub-structure dynamic analysis under harmonic excitation at its interface boundary. For the gravity structure Veletsos and Wei 1978[65], Luco and Westmann 1971[66] have obtained the foundation impedence assuming a uniform elastic half space. Luco 1974[67] obtained it assuming a layered viscoelastic half-space.

Due to the presence of piles, realistic model of the foundation subsystem for framed steel structure is more complicated. Novak 1974[68] generate the foundation impedence for single pile in a uniform layer of soil founded in rigid rock. Kausel and Roessel 1975[69]

applied the finite element representation to obtain the impedence function. Tsai and Housner 1970[70], Lysmer and Kuhlemeyer 1969[71] developed the "exact" wave transmitting boundary element.

As the foundation impedence and damping which consist of rediation damping causing loss of energy of the motion due to the soil properties hence it is frequency dependent.

A type of the soil modelling of the foundation for gravity type and framed steel type are shown in (Figure 2.2.3) and (Figure 2.2.4) respectively.

The effect of the soil parameter on the mathematical model of the foundation subsystem for the dynamic response of the structure had two principal effect according to (Taylor 1975[72] and Angelides and Conner 1979[73]).

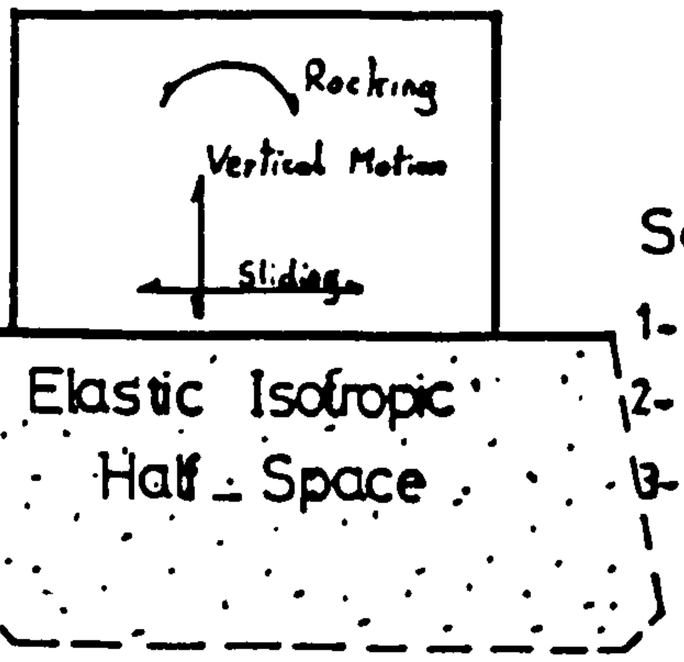


Figure (2.2.3) - Soil Model for Gravity Structure.

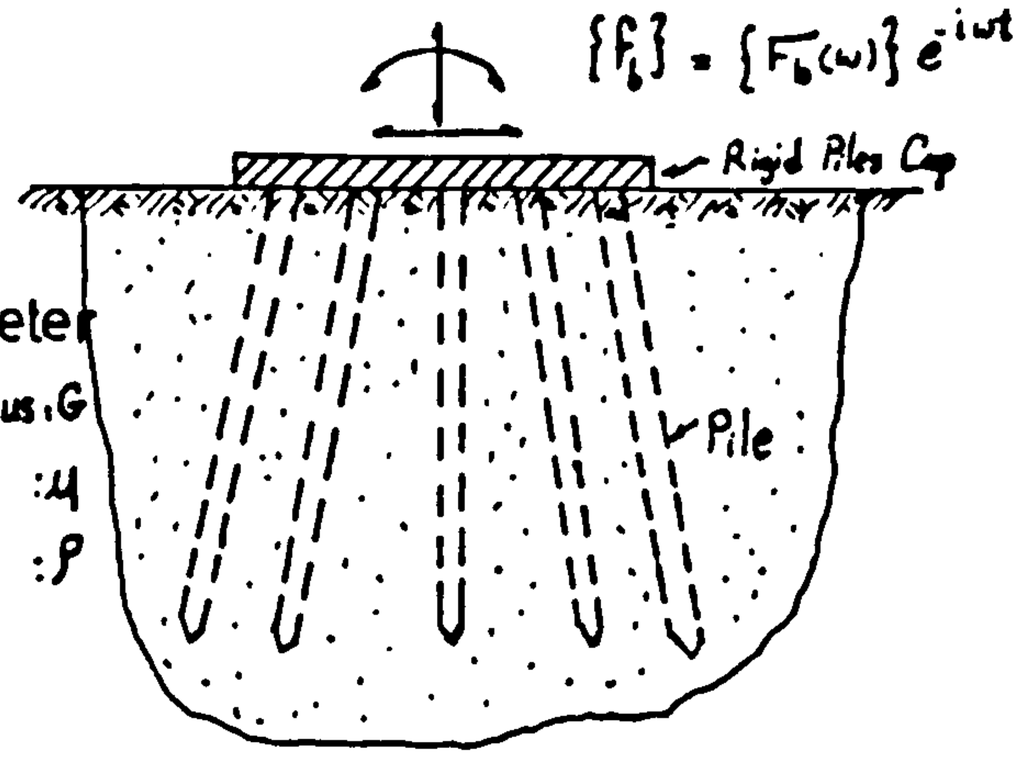


Figure (2.2.4) - Soil Model for Framed Steel Structure.

1. Significant shifts in the response were observed for reasonable modification in the soil properties

2. Reduction in foundation stiffness due to soil degradation increases the foundation period of the structure and the dynamic amplification of the structure response

The first point has been investigated by Moan et al 1975[63] for the shear modulus variation. The combination of the soil properties expressed in terms of damping and stiffness indicate that smaller soil damping will give higher dynamic amplification at resonnance (Bell et al 1976[74]). For the soil stiffness a higher dynamic amplification is found for a higher soil stiffness (Taylor 1974[75, 72]) and that is clear due to the small amount of energy absorped by the soil.

The second point has been investigated (e.g. Walt 1978[76]). It is important to represent the mathematical model for modelling the structure foundation interaction using the foundation impedence and damping. Figure (2.2.5) shows some of the modelling used.

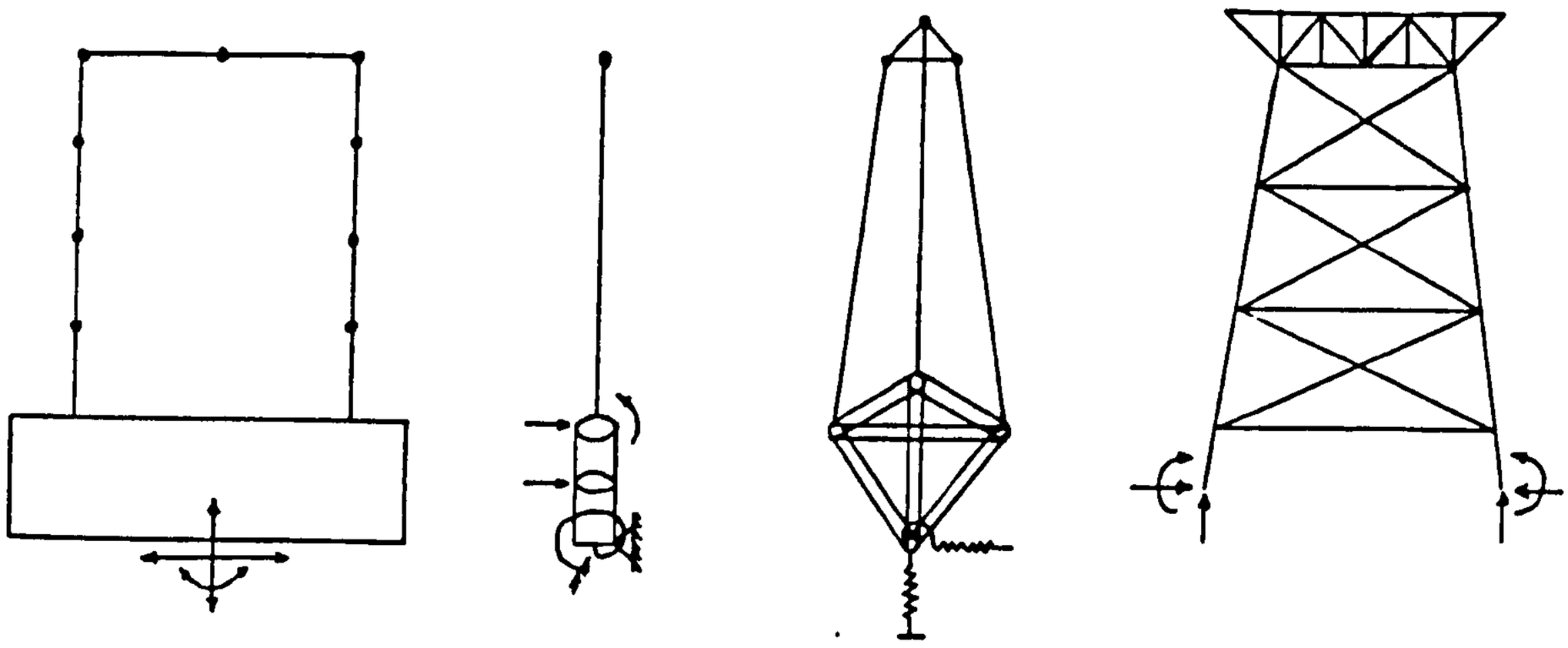


Figure (2.2.5) - Foundation Modelling.

As the foundation has been represented completely in the offshore structure the vital points now are to use the information from the foundation modelling in the analysis using the foundation structure-interaction, this point has been looked at by many professional investigators e.g. (Penzin 1976[77], Taylor[75]). The interaction can be represented by forming the stiffness and damping matrix which are used in the equation of motion governing the dynamic analysis of the structure.

$$[M]\{\ddot{X}\} + [C]\{\dot{X}\} + [K]\{X\} = \{F(t)\} \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (2.2.1)$$

where $[M]$ = Mass matrix

$[C]$ = Damping matrix

$[K]$ = Staffness matrix

$\{\ddot{X}\},\{\dot{X}\}$ and $\{X\}$ = Acceleration, velocity and displacement of the structure respectively

$\{F(t)\}$= Hydrodynamic force

by using the interaction between structure-foundation, the damping and the stiffness matrix will be

$$[C] = [C_s] + [C_f] \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (2.2.2)$$

$$[K] = [K_s] + [K_f] \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (2.2.3)$$

where $[C_s]$ = Damping in the free structure

$[C_f]$ = Damping in the foundation

$[K_s]$ = Stiffness in the free structure

$[K_c]$ = Stiffness for the foundation

One suitable method for evaluating $[C_s]$ is to use the dynamic properties of the fixed base structure along with the pseudo/static influence coefficient associated with its base degrees of freedom, i.e.

Let $$[C_s] = \begin{bmatrix} [C_s^f] & [C_s^b] \\ [C_s^b] & [C_s^{bb}] \end{bmatrix} \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (2.2.4)$$

where $[C_s^f]$ is the damping matrix for the fixed base structure,

$[C_s^{bb}]$ is the damping matrix associated with the base degree of freedom of the structure, and

$[C_s^b]$ is the coupling damping matrix,

Matrix $\left[C_s^f\right]$ can be formulated for the fixed base structure using the relation formulated by Clough 1975[78]

$$\left[C_s^f\right] = \left[M^f\right] \left[\sum_{J=1}^{m} \frac{2\zeta_J W_J}{M_J} \{\phi_J\}\{\phi_J\}^T\right] \left[M^f\right] \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (2.2.5)$$

where $\left[M^f\right]$ is the diagonal mass matrix, including hydrodynamic add mass,

m is the number of normal modes to be considered and

$\{\phi_J\}$, W_J, ζ_J and M_J are the mode shape vector, frequency, damping ratio and generalized mass respectively of the J^{th} normal mode.

Matrices $\left[C_s^b\right]$ and $\left[C_s^{bb}\right]$ can be formulated as follows:

$$\left[C_s^b\right] = - \left[C_s^f\right]\left[b_s\right] \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (2.2.6)$$

$$\left[C_s^{bb}\right] = \left[b_s\right]^T \left[C_s^f\right]\left[b_s\right] \quad \ldots \quad \ldots \quad (2.2.7)$$

where $\left[b_s\right]$ = pseudo-static influence coefficient matrix which is equal to

$$\left[b_s\right] = - \left[K_s^f\right]^{-1} \left[K_s^b\right] \quad \ldots \quad \ldots \quad \ldots \quad (2.2.8)$$

In which $\left[K_s^f\right]$ = Stiffness matrix of fixed base structure

$\left[K_s^b\right]$ = Stiffness-coupling matrix expressing the force developed in the degrees of freedom of the fixed base structure caused by the pseudo-static displacement of the base degrees of freedom.

From the above we can conclude that the understanding and evaluating the dynamic analysis of the fixed base structure should be carried out before establishing the analysis of the offshore structure with its foundation.

2.2.2 THE STRUCTURE SYSTEM

The structure system must be idealized so that the dynamic analysis can be made with simple systems forming the dynamic properties Mass, Stiffness and Damping of the structure.

2.2.2.1 THE MASS MATRIX

The mass matrix can be formulated either by consistent mass matrix or by the lumped mass matrix, the simplest procedure for defining the mass matrix of any structure is to assume that the entire mass is concentrated at the points of which the translation displacement are defined.

The usual procedure for defining the point mass to be located at each node is to assume that the structure is divided into segments, the modes serving as connection points.

The mass matrix of the offshore structure can be represented as

$$M = M_m + M_a + M_F \quad \ldots \quad \ldots \quad (2.2.9)$$

where M_m = Material structural mass and the platform mass

M_a = Added mass due to the displaced fluid replaced by the submerged part of the structure

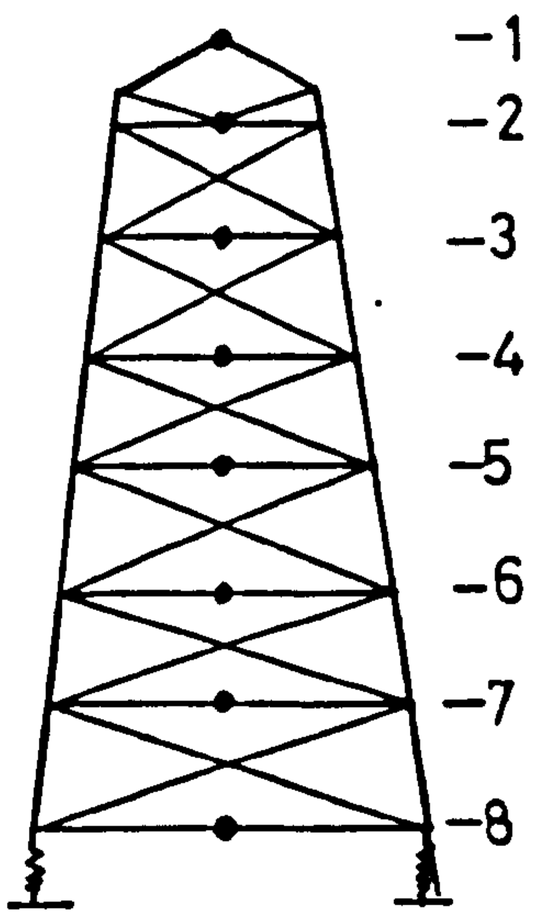
M_F = Flooding and marine growth mass

The part of M_F due to the flooding mass can be easily calculated and it is the mass of water occupying certain volume of structure element to make it flood. The part of M_F due to marine growth are not easy to predict due to uncertainties of the marine growth thickness over the whole structure. Heaf 1979[79] shows the effect of increased marine growth on the mass matrix. This effect appears in two ways; firstly, the increase in mass of the structure and secondly, the increase in hydrodynamic added mass due to an increase in displaced volume.

M_m is the mass of the structure material plus the mass of the platform deck including all the machinery and equipment. It is easy to calculate and form the mass matrix for M_m. But the only problem associated with M_m is to choose the position and number of mass points to be taken in the analysis.

Maddox 1974[80] has conducted a study to find the best position and number of discrete mass points to be considered by using two different representation of the mass points for the same structure, as shown in figures (2.2.6A and 2.2.6B) which represent model 1 and model 2 respectively.

A comparison between the two models has shown that model 1 can be considered to be an adequate representation of the structure to be accounted for.



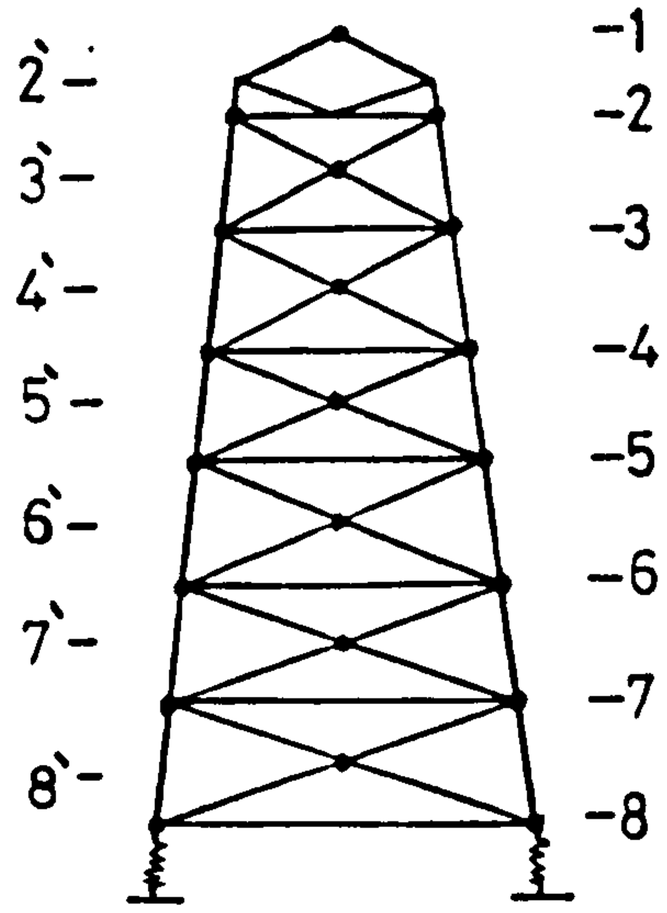


Figure (2.2.6a) - Model 1. Figure (2.2.6b) - Model 2.

The number of discrete mass points must be considered in the analysis, because at each discrete mass point the dynamic equation of motion is set. By using a large number of discrete mass points we get a similar number of equations of motion to be solved simultaneously which requires too much computer time and a high computer capacity. This is not the only disadvantage of choosing too many lumped masses but also it gives an inaccurate solution because of the higher frequency modes which will be concluded in the solution while not being required.

From the above it is important to choose the exact number of lumped mass to be taken in the mathematical modelling of the structure. Nath and Harleman 1969[81] used one lumped mass at the platform for two massless leg structures, the experimental result shows that there is quite good agreement with the predicted analysis for deep water wave and with different space between structures. Harleman et al 1963[82] carried out experiments for four massless leg structures using one lumped mass at the platform. They show general agreement between experimental and analytical prediction. Experimental and analytical studies for one-degree of freedom of one, two and four massless towers subjected to regular and random water waves had been performed by Nath and Harleman 1967[83] taking into consideration the stress and the moment of the structures and the effects of the spacing of vertical supports on the platform deflection and also the effect of wave direction emphasis, the general agreement with the use of the one-degree of freedom. The above experiments used structure model fixed at the base and the wave structure interaction were ignored. It does not show whether the model represent any full scale real structure.

Nolan and Honsinger 1962[84] found that for the single degree of freedom approximation for the structures assuming that all masses to be concentrated at the platform level, the ratio of the maximum theoretical to experimental displacement varied from 0·5 - 1·7.

Molhotra and Penzin 1970[85] mentioned that the behaviour of a tower in water subjected to random wave forces can be represented fairly accurately by using the first and second mode or first, second and third lateral mode of vibration. Also Bege and Penzien 1974[62] concluded that at least six normal modes of vibration should be included in the three-dimensional dynamic analysis.

Wu 1976[86] concluded that the number of discrete mass points necessary to satisfy the displacement analysis is less than that required for the analysis of shear force and bending moment.

The use of a large number of lumped masses is unnecessary.

M_a represent the mass of water attached to the structure and move with it as one body. The evaluation of this mass of water represented by the equation

$$M_a = (C_m - 1)\rho\forall \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (2.2.10)$$

where $\forall$ = Volume of water displaced by structure

ρ = Water density

C_m = Mass coefficient (inertia coefficient)

Equation (2.2.10) has a numerical coefficient which attribute uncertainties to the evaluation of the added mass. Many works assumed that C_M = 2·00 which means that the added mass will be equal to the mass of water displaced by the structure.

King 1971[87] support the idea of C_M = 2·00 and also mentioned that the added mass is unaffected by streaming flow and vortex shedding, as well as the added mass function was seen to be independent of frequency, amplitude and model shape. Jenssen et al 1977[88] observe that the above statement of King is not absolute as they found a variation in added mass near the surface. Sarpkaya 1981[89] dealt with the question of instantaneous value of inertia coefficient which lead to negative added mass in some parts of the wave cycle, which prove that the added mass is not a constant value during the wave cycle and also it varies due to wave and structure condition.

2.2.2.2 THE STIFFNESS MATRIX

The stiffness matrix is independent of the type of the structure whether offshore or inshore. It is also independent of the solution type i.e. whether static or dynamic, although some papers state that there are changes in stiffness matrix due to the difference in the dynamic Young's modulus. (See Nath and Harleman 1969[81] and Hallam et al 1978[90]).

2.2.2.3 DAMPING MATRIX

It is important to understand the general description of damping. Damping is basically a dissipation of energy due to motion. There are several possibilities for energy dissipation:

(i) MATERIAL DAMPING

The material damping is thought to be very small. This component may mathematically be represented as viscous damping. The damping coefficient may be assumed to be constant throughout the structure and only depend on the structure material.

(ii) STRUCTURAL DAMPING

The structural damping is caused by the way the structure is assembled (bracing, welding) and on the structure mechanism. The assessment of material damping is equally valid for structural damping.

(iii) ENVIRONMENTAL DAMPING

For offshore structure the environmental damping is mainly of hydrodynamic origin. Damping forces due to movement of air above the water line are negligible compared with the hydrodynamic damping.

The above physical possibilities of damping summed together to represent and use of damping in computations. In calculations, the damping are represented by the damping coefficient which can be calculated from the equation:

$$\zeta = \frac{C}{C_c} \times 100 \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (2.2.11)$$

the damping coefficient is usually imagined to be constant throughout the structure which is not a good physical description (see Vugts and Hayes 1979[91]), the value of this coefficient is given in the form of a percentage of critical damping, as critical damping cannot be defined for a general N-degree of freedom, the natural period of vibration is important and to get the damping coefficient in the equation

$$C_c = 2\ M\ W_n \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (2.2.12)$$

where M = mass of structure

W_n = natural frequency of structure

The most vital part is to establish the appropriate damping coefficient. Several papers have been published concerning this point, some used model structures (Skep et al 1976[92] and King 1972[93]) and others used prototype structures.

The damping coefficient of the structure can be measured by one of the three possible methods

(a) Log-decrement measurements of decay vibration resulting from a free vibration due to tip displacement

(b) Response measurements under known imposed excitation

(c) Response measurements under natural excitation for offshore structure

For offshore structure methods (a) and (b) seem practical for establishing the damping coefficient, the snag in both methods has been pointed out [91], also in the mathematical representation of the physical realities Three main factors affect the choice of a preferable method: nonlinear effects, local damping effect and frequency dependent coefficients.

Ogilive 1964[94] found that the choice of a damping coefficient has strong effect upon forces and displacement but it had hardly any effect on the resonant frequency.

In reviewing the work done by some offshore professional in estimating the damping coefficient, Cherry and Brady 1965[95] used the concept of outocorrelation technique for random vibration of experimental data as a method for predicting structure dynamic properties (Natural period and damping coefficient), the result concluded that reasonable estimation of structure period can be obtained from an analysis involving small amounts of data, relatively large sample sizes are required if meaningful estimates of structure damping coefficient are expected.

A method of measuring structural damping based on time delay correlation has been established by Jeary and Winney 1972[96] for random vibration experiment data has shown that the results obtained with this method is in good agreement with the results obtained from the standard decay tests. Vanmarke 1971[97] outlines a new method based on power spectral density analysis for determining period and damping values of offshore structure from field measurements of their response to wave-induced random excitation. Ruhl and Berdahl 1979[98] carried out forced vibration tests to find the damping coefficient they found that the damping coefficient is less than two percent of the fundamental group of modes and is three to five percent range for the second group of modes. They used the steady state test which show that these tests are extremely useful for determination of coefficient of damping. Also Ruhl

1976[90] conducted free vibration tests and spectral analysis of random response data to find the fundamental periods and damping coefficient using the partial spectral moment method for the random vibration.

Another way of getting the damping of offshore structures is by separating it into two parts, first the material and structure damping and second the environmental "Hydrodynamic Damping. Skop et al 1972[92] evaluated a damping empirical formula for circular cylinder, but it had not been used or proven. King 1972[93] tried the hydrodynamic damping in still water, concludes that it is pure viscous and may be described by Stokes element damping theory. His work also shows that the damping increases rapidly with water depth. Verley 1978[100] compared between the independent flow fields, i.e. by taking the damping as being the still water and the hydrodynamic added damping due to the relative velocity, his result of the analysis confirms the suggestion of using the independent flow fields for low values of the reduced velocity ($U_r = U_m/f_nD$) and low expected vibration amplitude as the relative velocity assumption for calculating the hydrodynamic damping will overestimate the damping, and underestimate vibration. Verley and Moe 1979[101] carried out experimental investigation of oscillating cylinder in current and the results indicate the same conclusion of those of Verley 1978[102],[100]. Dean et al 1979[103] in his report tried to establish values to be added, called

(implied damping) to substitute the hydrodynamic damping but in their report they could not generalize the values for the sea state.

2.3 NUMERICAL PROCEDURE

A realistic dynamic analysis of the fixed offshore structure is complex. Hence approximation analogous to the single-degree of freedom system is adequate. But a full and comprehensive dynamic analysis appears to be the only satisfactory way of dealing with the problem.

Two general aspects of the problem must be briefly discussed.

2.3.1 TIME DOMAIN AND FREQUENCY DOMAIN DESCRIPTION

A distinction in time and frequency domain analyses is governed by the manner in which the excitation is handled; (The right hand side of the equation of motion). With a description in the time domain the full time history of every quantity is involved, certain initial conditions are specified and with the loading distribution over the whole structure known as a function of time, the response of the structure may be obtained at each and every instant.

2.3.1.1 MATHEMATICAL FORMULATION OF THE DYNAMIC PROBLEM IN FREQUENCY DOMAIN DESCRIPTION

- Constant and frequency depended coefficient can both simply and equally be handled. These coefficients may arise from e.g. hydrodynamic and soil interaction (added mass, damping coefficient). Time dependent coefficients are strictly not admissible.

- Non-linearities in system properties (stiffness, damping) or in loading wave force can only be modelled by linearization.

- The equation of motion have the appearance of differential equation but are in fact algebraic equations (see Ogilvie 1964[94]).

- The steady state solution only is obtained, and it is obtained directly.

- For N degree of freedom and harmonic excitation, a system of 2N linear algebraic equations must be solved at each frequency. The response to a random excitation can be obtained through spectral analysis.

- An arbitrary though prescribed excitation over a short period cannot be handled in the sense that the response include transients but it can be analysed using Fourier series.

2.3.1.2 MATHEMATICAL FORMULATION OF THE DYNAMIC PROBLEM IN TIME DOMAIN DESCRIPTION

In the time domain with constant coefficients the equations are differential equations. Frequency dependent coefficients lead to the inclusion of retardation functions, the equations become integro-differential equations (see Ogilvie). Coefficients that are functions of time can be handled.

- The equations are differential or integro-differential equations.

- Transient and/or steady state solutions can be obtained. Initial conditions have to be specified. Number of cycles required for steady state to be reached may be large depending on initial conditions and damping.

- For N degree of freedom and harmonic excitation of a linear system under certain assumptions the set of N second order differential equations may be replaced by N simultataneous equation for every time step response, to random excitation may be obtained directly through numerical integration.

- Any non-linearity can in principle be included provided that an adequate mathematical formulation and solution procedure is available.

- The differential equations must be solved numerically with small time steps, initial conditions and a criterion to define steady state condition have to be specified.

- Numerical integration under random excitation is a time consuming process. For different excitation the whole process must be repeated. In general the random loading distributions are required.

- Any prescribed excitation can be dealt with by direct numerical integration given appropriate conditions and a reliable integration technique.

2.3.2 MODEL SUPERPOSITION AND DIRECT INTEGRATION

The distinction between model superposition and direct integration techniques lies in the manner by which the equation describing the behaviour of the system are dealt with.

The model superposition technique makes use of the fact that the response of a complex multi-degree of freedom system can be described as the sum of the responses of a number of one-degree of freedom system. By appropriate coordinate transformations the degrees of freedom are uncoupled, so it improves the understanding of the system properties, also the excitation has also to be

transformed to refer to the same coordinate system. The system may be continuous or discretized.

The direct integration approach, the coupled equation which in practical cases would refer to a discretized system are integrated as they stand.

2.3.2.1 MATHEMATICAL FORMULATION OF THE DYNAMIC PROBLEM BY MODEL SUPERPOSITION

- The basic features of this technique is that N degree of freedom system are uncoupled into N equations each describing a single degree of freedom system. The technique is exploited when the eigenvalues of the coupled equation is real.

- The essential difficulty with model superposition is the number of modes that have to be considered.

- The system non-linearities cannot be included.

- The first requirement in an eigenvalue solution in order to find the undamped natural frequencies and mode shapes.

- The transformation to generalized coordinates is lengthy but needs to be done only once. Only the transformation of the loading function is to be represented.

- Relatively small computer can handle the problem.

2.3.2.2 MATHEMATICAL FORMULATION OF THE DYNAMIC PROBLEM BY DIRECT INTEGRATION

- The coupled equation is dealt with directly.

- There are no questions of truncation of a series, the full solution is obtained.

- Non-linearities can in principle be incorporated.

- Although not a part of the solution procedure, definition of damping may be facilitated by an eigenvalue solution of the first few natural frequencies. Knowledge of the principal mode shapes also improve physical insight.

- No transformation is required.

- Large computer is required.

Two practical examples of the way in which dynamic analysis have been incorporated into particular design problems in both the time and frequency domain had been carried out by Schumm 1978[104] for steel jacket platform. It showed that the advantage of using time domain analysis are as follows:

1. The correct structure submergence for members in the region of the still water surface.

2. Drag force can be included without linearizations.

3. Relative velocity between wave and structure can be correctly accounted for.

Also in frequency domain it shows that despite the necessity to perform linearization, the frequency domain approach provides an extremely powerful tool for stochastic analysis for long term effects as it reduces the amount of data processing to a manageable size and significantly reduces computing cost over an equivalent time domain approach. It shows also that the best suited technique of the solution of mathematical model for steel platform is the model superposition technique in which case the major response of the structure can be adequately represented by the first ten or so mode shapes without a significant loss of accuracy.

Vugts 1979[105] made a comparison between the solution of mathematical model, also alternative model superposition analysis using varying numbers of modes to account for dynamic effects, supplemented by the full static solution which replaces the corresponding static component contained in the solution of the truncated model series has been used. All calculations were performed in frequency domain for two derivatives of the structure:

the first is made artificially stiff in order to make it dynamically insensitive to the wave load and the second is made artificially flexible in order to produce significant dynamic response to the prescribed wave excitation. This work indicates that pure model superposition is unacceptable if a realistic prediction of detailed member end forces and moment (and hence stresses) is the primary objective. The comparison with exact direct solutions confirm that model analysis provides reasonably reliable estimates of gross platform behaviour, expressed in terms of global horizontal displacement even when relatively few modes are used.

A considerable improvement in accuracy observed when static response of higher modes, which were truncated in the model series, has a large influence upon the number force contribution and firmly recommended that however stress recovery is the primary objective of a dynamic analysis of an offshore structure a direct solution of the equation of motion in the physical coordinates should be performed. In the stiff structure it is not necessary to perform the direct dynamic solution as it more costly than model superposition.

CHAPTER THREE

THEORETICAL FORMULATION

3.1 INTRODUCTION

The problem of determining structural interaction of cylinders in periodic waves will be examined theoretically. The model posed is a circular cylinder fixed at the base and either free or fixed at the top (i.e. free to vibrate or not) subjected to periodic waves.

The periodic wave form will be assumed to be linear. The force on the structure due to these waves will be obtained by the Modified Morison's equation as well as Morison's equation.

The dynamic sturctural response to waves will be calculated using mode superposition in the time domain. A finite element method is used to solve the resulting dynamical equations.

3.2 WAVE DYNAMICS

The accuracy of the determination of the water wave force depends on the accurate description of waves. This is needed to obtain the appropriate wave velocity and accelerations necessary for the prediction of wave forces.

3.2.1 WAVE DESCRIPTION

The wave flow field used in this study is assumed to be in two dimensions. Also it is assumed to be incompressible and irrotational which will satisfy Laplace's equation (2.1.1)

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = 0 \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.2.1)$$

The bottom surface is assumed to be horizontal, rigid and impermeable and hence Equation (2.1.5) reduces to

$$\frac{\partial \phi}{\partial y} = 0 \qquad \text{at } y = -h \qquad \ldots \quad \ldots \quad \ldots \quad (3.2.2)$$

The Equation (2.1.4) is linearized by neglecting the velocity components, also by assuming small amplitude wave

$$\therefore \quad \eta = \frac{1}{g} \frac{\partial \phi}{\partial t} \qquad \text{at } y = 0 \qquad \ldots \quad \ldots \quad \ldots \quad (3.2.3)$$

To satisfy the linearization of the boundary condition, the ratio between the neglected terms and terms which have been retained should be small. This is found to give:

$$ak \ll 1, \ (\tanh Kh \text{ and } \cosh Kh \ll 1) \qquad \ldots \qquad (3.2.4)$$

if the velocity potential is considered to be the first term ϕ_{11} in a perturbation expansion

$$\phi = \phi_{11} + \phi_{12} + \phi_{13} + \ldots \phi_{n} \qquad \ldots \quad \ldots \qquad (3.2.5)$$

where $\phi_{n-1} >> \phi_{n}$ for all n

The ratio of ϕ_{12}/ϕ_{11} should be small where (ϕ_{12} is the Stokes second order term)

$$\phi_{12} = 3/8 \ \frac{a^2 g}{C} \ \frac{\cosh 2k\,(h + z)}{\sinh^3 Kh \cosh Kh} \ \sin 2(kx - \sigma t)$$

$$\ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \qquad (3.2.6)$$

To satisfy the above requirement the following condition must hold

$$a/K^2h^3 << 1 \qquad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \qquad (3.2.7)$$

The wave generated contained small amplitude of higher order free harmonic waves. The wave profile was periodic

$$\therefore \quad \eta(t + \tau) = \eta(t) \qquad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \qquad (3.2.8)$$

The experimental wave profile obtained for one wave cycle was analysed using the Fast Fourier Transfor to obtain the magnitude of each of the individual free harmonic waves

$$\eta_{total} = \sum_{n} a_n \cos(K_n x - \sigma_n t) + b_n \sin(K_n x - \sigma_n t)$$

... ... (3.2.9)

$$\therefore \quad \eta_{total} = \sum_{n} A_n \cos(K_n x - \sigma_n t + \alpha_n) \quad \ldots \quad (3.2.10)$$

where A_n = amplitude of the nth harmonic in the wave

$$= \sqrt{a_n^2 + b_n^2}$$

α_n = the phase shift relationship between the various waves (nth harmonic) and the measured from the original (kx - σt)and equal to $\tan^{-1} \frac{a_n}{b_n}$

Since Laplace's equation is linear, the velocity potential of a wave system ϕ_{total} is given by the sum of the potential of the individual waves

$$\phi_{total} = \phi_1 + \phi_2 + \phi_3 + \ldots \phi_n \quad \ldots \quad \ldots \quad (3.2.11)$$

in which

$$\phi_n = \frac{A_n g}{\sigma_n} \quad \frac{\cosh K_n(h + z)}{\cosh K_n h} \quad \sin(K_n x - \sigma_n t + \alpha_n)$$

... ... (3.2.12)

and the dispersion relation

$$C_n^2 = (g/K_n) \tanh K_n h \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.2.13)$$

where g = gravitational acceleration

σ_n = the nth wave angular frequency = $2\pi/T_n$

K_n = the nth wave number = $2\pi/L_n$

h = distance from the still water level to bottom

C_n = velocity of the nth wave propagation (phase velocity) = L_n/T_n

Z = vertical distance measured positive upward from the still water level

Since the fluid particle velocity and acceleration are first derivative of ϕ, they can be added because ϕ is linear

$$\dot{U}_{total} = \dot{U}_1 + \dot{U}_2 + \dot{U}_3 + \ldots + \dot{U}_n \quad \ldots \quad \ldots \qquad (3.2.14)$$

where

$$\dot{U}_n = \frac{\partial \phi_n}{\partial x} = \frac{A_n g K_n}{\sigma_n} \frac{\cosh K_n(h + z)}{\cosh K_n h} \cos (K_n x - \sigma_n t + \alpha_n)$$

$$\ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \qquad (3.2.15)$$

and

$$\ddot{U}_n = \ddot{U}_1 + \ddot{U}_2 + \ddot{U}_3 + \ldots + \ddot{U}_n \quad \ldots \quad \ldots \qquad (3.2.16)$$

where

$$\ddot{U}_n = \frac{\partial \dot{U}_n}{\partial t} = A_n \, gK \quad \frac{\cosh K_n(h + z)}{\cosh K_n h} \quad \sin(K_n x - \sigma_n t + \alpha_n)$$

... ... (3.2.17)

3.2.2 WAVE FORCES

The wave forces on structure will be determining using Morison's equation. This force varies with time during the wave cycle. It also varies with the position of different element along the structure. The force on the element is calculated using the average values of the drag and the inertia coefficients during the wave cycle (over the cycle C_D and C_M are assumed constant).

$$dF_{(t,\,z)} = \left[C_{D(z)} \frac{1}{2} \rho D(\dot{U}_{total(t,\,z)} \left| \dot{U}_{total(t,\,z)} \right|) + C_{M(z)} \frac{\rho \pi D^2}{4} \ddot{U}_{total(t,\,z)}\right] dz \quad \ldots \quad (3.2.18)$$

In the case where the structure is free to vibrate the force on the element will be represented by

$$dF'_{(t,\,z)} = \left[C'_{D(z)} \frac{1}{2} \rho D(\dot{U}_{total(t,\,z)} - \dot{X}_{(t,z)}) \left| \dot{U}_{total(t,\,z)} - \dot{X}_{(t,z)} \right| + (C'_M - 1) \frac{\rho \pi D^2}{4} (\ddot{U}_{total(t,\,z)} - \ddot{X}_{(t,\,z)}) + \frac{\rho \pi D^2}{4} \ddot{U}_{total(t,\,z)}\right] dz \quad \ldots \quad \ldots \quad (3.2.19)$$

The corresponding bending moment at the base of the structure due to the wave force can be obtained by integrating of Equations(3.2.18) and (3.2.19) along the submerged part of the structure for the case of small structure's deflection.

For the case of free structure (cantilever)

$$M_B = \int_{z_o}^{h} dF_{(t,z)}\, zdz$$

or

$$M_B = \int_{z_o}^{h} dF'_{(t,z)}\, zdz \quad \ldots \quad \ldots \quad \ldots \quad \ldots \qquad (3.2.20)$$

For the case of fixed structure

$$M_B = \int_{z_o}^{h} dF_{(t,z)}\, \frac{z(L_s - z)^2}{L_s}\, dz \quad \ldots \quad \ldots \qquad (3.2.21)$$

where L_s = structure length.

The appropriate drag and inertia coefficients are obtained using either Equation (3.1.15) or (3.1.16) in conjunction with the measured wave force. The method of obtaining the measured wave force, the parameter $\dot{X}_{(t,\ z)}$ and $\ddot{X}_{(t,\ z)}$ will be shown later in this chapter.

The average values of the drag and inertia coefficients during the wave cycle can be calculated by one of the following methods:

1. Fourier analysis

2. Cross point(calculated of the drag coefficient at zero inertia force and calculating of the inertia coefficient at zero drag force)

3. Least square method

The Fourier analysis for an oscillating flow represented by

$$\dot{U} = \dot{U}_m \cos \sigma t \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.2.22)$$

the average value of C_D and C_M are given by Keulegan and Carpenter(27)

$$C_D = - 0.75 \int_0^{2\pi} (F_m \cos \sigma t / \rho \dot{U}_m^2 \, SD) dt \quad \ldots \quad (3.2.23)$$

and

$$C_M = (2\dot{U}_m T/\pi^3 D) \int_0^{2\pi} F_m \sin \sigma t / \rho \dot{U}_m^2 \, SD) dt \quad (3.2.24)$$

in which F_m is the measured force and S is the length of the segment over which the force was measured. This method yields the same value of C_M as that of the least square method but with slight difference in

the value of C_D. This method is not applicable if there are free higher order harmonic waves present in the wave. The omission of the second harmonic wave will affect the value of the force coefficients (see D. I. Maull and M. G. Milliner 1979[61]).

The cross point will give discrepances due to the difficulty in determining the region of the zero drag and zero inertia during the wave cycle.

The least square method is more applicable, as it can deal with higher harmonic component in the wave. The least square method used is as follows

$$F_m = (f_I)\, C_M + (f_D)\, C_D \quad \ldots \quad \ldots \quad \ldots \quad \ldots \qquad (3.2.25)$$

where

$$(f_I) = \frac{\rho \pi D^2}{4} \ddot{U}_{total(t,\ z)} S \quad \ldots \quad \ldots \quad \ldots \qquad (3.2.26)$$

$$(f_D) = \frac{1}{2} \rho D\, \dot{U}_{total(t,\ z)} \left|\dot{U}_{total(t,\ z)}\right| S \qquad (3.2.27)$$

$$\Sigma(f_D)^2\, C_M + \Sigma(f_I)(f_D) C_D = \Sigma(f_I\, F_m) \qquad \ldots \qquad (3.2.28)$$

$$\Sigma(f_I)(f_D) C_M + \Sigma(f_D)^2 C_D = \Sigma(f_D\, F_m) \qquad \ldots \qquad (3.2.29)$$

$$\therefore C_{D(z)} = \frac{\Sigma(f_I)^2 \Sigma(f_D F_m) - \Sigma(f_I F_m) \Sigma(f_I)(f_D)}{\Sigma(f_I)^2 \Sigma(f_D)^2 - \Sigma((f_I)(f_D))^2} \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.2.30)$$

$$C_{M(z)} = \frac{\Sigma(F_m f_I) \Sigma(f_D)^2 - \Sigma(f_D F_m) \Sigma(f_I)(f_D)}{\Sigma(f_I)^2 \Sigma(f_D)^2 - \Sigma((f_I)(f_D))^2} \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.2.31)$$

where F_m = measured force for the case either (fixed or free).

In the case of the modified Morison's equation is used Equation (3.2.25) becomes

$$F_m = (F_I) + (f_I')C_m' + (f_D')C_D' \quad \ldots \quad \ldots \quad \ldots \quad (3.2.32)$$

where

$$F_I = -\frac{\rho\pi D^2}{4} \ddot{X}_{(t,\ z)}\, S \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.2.33)$$

$$(f_I') = \frac{\rho\pi D^2}{4} (\ddot{U}_{total(t,\ z)} - \ddot{X}_{(t,\ z)})S \quad \ldots \quad (3.2.34)$$

$$(f_D') = \frac{1}{2} \rho D(\dot{U}_{total(t,\ z)} - \dot{X}_{(t,z)}) |\ddot{U}_{total(t,\ z)} - \dot{X}_{(t,\ z)}|S \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.2.35)$$

by arranging Equation (3.2.29)

$$F_m' = (f_I')C_M' + (f_D')C_D' \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.2.36)$$

the values of $C_{D(z)}$ and $C_{M(z)}$ can be obtained from Equations (3.2.30) and (3.2.31) using these new values of F'_m, f'_I, f'_D.

3.3 STRUCTURAL DYNAMICS

The problem of vibrating structure is solved by equilibrating the externally applied force with the inertia forces resulting from the acceleration of the structure, the elastic resistance to displacement and the energy-loss due to the structure mechanism (damping). The difficulties arises in the prediction of structural dynamic of offshore structure is due to the dependence of external applied force on the structure oscillation.

The prediction of the dynamic response of offshore structure will depend upon the idealization used to model the structure system, the formulation of the equation of structure system response and the method of solution of the structural system equations with its numerical analysis.

3.3.1 THE STRUCTURE MODELLING

The structural system used in this study is modelled as follows:

- The continuous member is assumed to be series of discrete masses connected by springs (with no mass)

- The continuous wave forces are represented by point loads acting at the mass centres

Figure (3.3.1(a)) shows the actual structure and the wave force while Figure (3.3.1(b)) shows the modelled structure

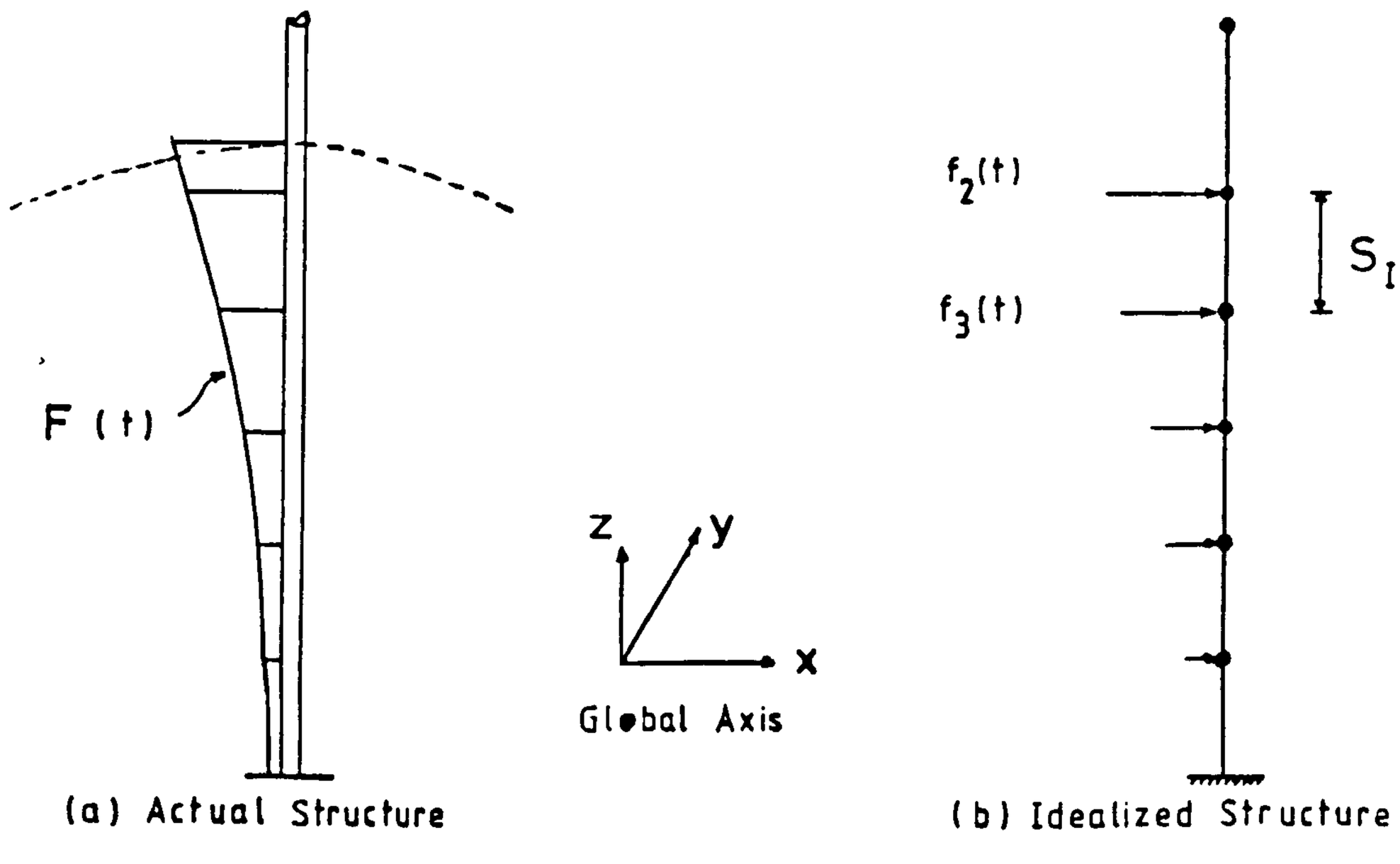


Figure (3.3.1) - Modelling of the Structures.

Further idealizations upon which the analysis is based are (a) stress is proportional to strain, (b) small deflection theory, (c) all motion is measured from the position of static equilibrium.

In this study, it is assumed that the waves travel in the positive x direction. All structural and loading variations in the y direction are therefore neglected. Thus, the number of degrees of freedom at each node is reduced to three; namely, translational displacements in the horizontal x direction and in the vertical direction, and rotational dizplacement about the y-axis.

3.3.2 FORMULATION OF THE EQUATION OF STRUCTURAL SYSTEM RESPONSE

The force on the structure including the interaction effect and noting that $\dot{U}_o \simeq \dot{U}_x$ and $\ddot{U}_o \simeq \ddot{U}_x$, where substript 'o' indicates that these quantities are taken at the undeflected structure location, the complete dynamic equation for the idealized structure shown in Figure (3.3.1(b)) may be written as follows

$$\begin{bmatrix} M_1 & 0 & 0 & 0 & 0 \\ 0 & M_2 & 0 & 0 & 0 \\ \cdot & & & & \\ \cdot & & & & \\ \cdot & & & & \\ 0 & 0 & 0 & 0 & M_n \end{bmatrix} \begin{Bmatrix} \ddot{X}_1 \\ \ddot{X}_2 \\ \cdot \\ \cdot \\ \cdot \\ \ddot{X}_n \end{Bmatrix} + \begin{bmatrix} C_{11} & C_{12} & C_{13} & \cdots & C_{1n} \\ C_{21} & C_{22} & C_{23} & \cdots & C_{2n} \\ & & & & \\ & & & & \\ C_{n1} & C_{n2} & C_{n3} & \cdots & C_{nn} \end{bmatrix} \begin{Bmatrix} \dot{X}_1 \\ \dot{X}_2 \\ \cdot \\ \cdot \\ \cdot \\ \dot{X}_n \end{Bmatrix}$$

$$+ \begin{bmatrix} K_{11} & K_{12} & K_{13} & \cdots & K_{1n} \\ K_{21} & K_{22} & K_{23} & \cdots & K_{2n} \\ \cdot & & & & \\ \cdot & & & & \\ \cdot & & & & \\ K_{n1} & K_{n2} & K_{n3} & \cdots & K_{nn} \end{bmatrix} \begin{Bmatrix} X_1 \\ X_2 \\ \cdot \\ \cdot \\ \cdot \\ X_n \end{Bmatrix}$$

$$= \begin{Bmatrix} P_1(t) \\ P_2(t) \\ \cdot \\ \cdot \\ \cdot \\ P_n(t) \end{Bmatrix} \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.3.1)$$

where the column matrix {P(t)} can be obtained from modified Morison's equation (3.2.19)

$$
\begin{Bmatrix} P_1(t) \\ P_2 \\ \\ \\ \vdots \\ P_n(t) \end{Bmatrix} = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & F_{m2} & 0 & 0 & 0 \\ 0 & 0 & F_{m3} & 0 & 0 \\ \vdots & & & & \\ 0 & 0 & 0 & 0 & F_{mn} \end{bmatrix} \begin{Bmatrix} 0 & -\ddot{X}_1 \\ \ddot{U}_{o2} & -\ddot{X}_2 \\ \ddot{U}_{o3} & -\ddot{X}_3 \\ \vdots & \\ \ddot{U}_{on} & -\ddot{X}_n \end{Bmatrix} +
$$

$$
+ \begin{bmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & F'_{m2} & 0 & 0 & 0 \\ 0 & 0 & F'_{m3} & 0 & 0 \\ \vdots & & & & \\ 0 & 0 & 0 & 0 & F'_{mn} \end{bmatrix} \begin{Bmatrix} 0 \\ \ddot{U}_{o2} \\ \ddot{U}_{o3} \\ \vdots \\ \ddot{U}_{on} \end{Bmatrix} +
$$

$$
+ \begin{bmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & F_{D2} & 0 & 0 & 0 \\ 0 & 0 & F_{D3} & 0 & 0 \\ 0 & 0 & 0 & 0 & F_{Dn} \end{bmatrix} \begin{Bmatrix} (0 \quad -\dot{X}_1) \mid 0 \quad -\dot{X}_1 \mid \\ (\dot{U}_{o2} - \dot{X}_2) \mid \dot{U}_{2} - \dot{X}_2 \mid \\ (\dot{U}_{o3} - \dot{X}_3) \mid \dot{U}_{3} \quad \dot{X}_3 \mid \\ (\dot{U}_{on} - \dot{X}_n) \mid \dot{U}_{on} - \dot{X}_n \mid \end{Bmatrix}
$$

$$
\ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \qquad (3.3.2)
$$

where

$$F_{mn} = \frac{\rho \pi D^2}{4} (C_{Mn} - 1) S_n$$

$$F'_{Mn} = \frac{\rho \pi D^2}{4} S_n$$

$$F_{Dn} = \frac{1}{2} \rho D C_{Dn} S_n$$

Here and thereafter, the agreement t of these time dependent quantities is dropped for convenience, except when a time dependent quantity first appears.

Equations(3.3.1) and (3.3.2) can be written as

$$[\sim M \sim]\{\ddot{X}\} + [C]\{\dot{X}\} + [K]\{X\} = [\sim F_m \sim]\{\ddot{U}_o - \ddot{X}\} + [\sim F_m \sim \{\ddot{U}_o\} +$$

$$+ [\sim F_D \sim]\{(\dot{U}_o - \dot{X}) |\dot{U}_o - \dot{X}|\}$$

... ... (3.3.3)

By rearranging the above equation

$$[\sim M + M_m \sim]\{\ddot{X}\} + [[C] + [\sim 2\dot{U}_o F_D \sim]]\{\dot{X}\} + [K] \{X\} +$$

$$+ \{E(\dot{X})\} = \{F(t)\}$$

... ... (3.3.4)

where

$$\{F(t)\} = \left[\left[\sim F_m \sim\right] + \left[\sim F'_m \sim\right]\right] \{\ddot{U}_o\} +$$

$$+ \left[\sim F_D \sim\right] \{\dot{U}_o|\dot{u}_o|\} \quad \ldots \quad \ldots \quad \ldots \qquad (3.3.5)$$

From Equation (3.3.4) it can be seen that the water wave/structure interaction produces a hydrodynamic damping due to the drag force, which is $\left[\sim 2\dot{U}F_o \sim\right] \{\dot{X}\}$ and a nonlinear vector $\{E(\dot{X})\}$ which has the form

$$\{E(\dot{x})\} = - \left[\sim F_D \sim\right]\{\dot{X}|\dot{x}|\} \quad \ldots \quad \ldots \quad \ldots \qquad (3.3.6)$$

The nonlinear nature of Equation (3.3.4) necessitates an examination of solution methods for systems of nonlinear differential equation. The following methods are used for the solution of nonlinear differential equation

1. Perturbation method
2. Semilinear method

The above methods depend upon the condition that the nonlinearity in the system is small.

PERTURBATION METHOD

In this method the nonlinear term $\{E(\dot{X})\}$ expressed

$$\{E(\dot{x})\} = \varepsilon\{E'(\dot{x})\} \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.3.7)$$

where $\{E(\dot{X})\}$ is the nonlinear damping function and $\{E'(\dot{X})\}$ described the shape of the nonlinearity and the relative magnitude is indicated by the parameter ε

when $\varepsilon = 0$ the Equation (3.3.4) reduces the equation to linear oscillations. The perturbation method is based on the assumption that, the solution of the Equation (3.3.4) permits an expansion in power of ε. It is only expected to be valid for motions in which the nonlinear part of the damping force remain small in comparison with the linear part.

$$\therefore \quad \begin{Bmatrix} X_1 \\ X_2 \\ \cdot \\ \cdot \\ \cdot \\ X_n \end{Bmatrix} = \begin{Bmatrix} X_{10} + \varepsilon X_{11} + \varepsilon^2 X_{12} + \ldots \\ X_{20} + \varepsilon X_{21} + \varepsilon^2 X_{22} + \ldots \\ \\ \\ \\ X_{n0} + \varepsilon X_{n1} + \varepsilon^2 X_{n2} + \ldots \end{Bmatrix}$$

$$\ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.3.8)$$

Equation (3.3.8) can be written as

$$\{X_n\} = \sum_{I=D}^{NP} \varepsilon^I \{X_{nI}\} \quad \ldots \quad \ldots \quad \ldots \quad \ldots \qquad (3.3.9)$$

where NP = an arbitrarily large integer

n = number of node point

The Ith term of the quantity represents the Ith order correction to the (I-1)th order solution for that quantity . The expansion of the shape function $\{E'(\dot{X})\}$ by using the truncated Taylor series with respect with $(\varepsilon\Delta X)$ is

$$\begin{Bmatrix} E'(\dot{X})_1 \\ E'(\dot{X})_2 \\ \cdot \\ \cdot \\ \cdot \\ E'(\dot{X})_n \end{Bmatrix} = \begin{Bmatrix} E'(\dot{X})_{1O} + \frac{(\varepsilon\Delta\dot{X})}{1!} \frac{d'E(\dot{X})_{1O}}{d\,\dot{x}_{1O}} + \frac{(\varepsilon\Delta\dot{X})^2}{2!} \frac{d'^2E(\dot{X})_{1O}}{d\,\dot{x}_{1O}^2} + \cdots \\ \\ \\ \\ \\ E'(\dot{X})_{nO} + \frac{(\varepsilon\Delta\dot{X})}{1!} \frac{d'E(\dot{X})_{nO}}{d\,\dot{x}_{nO}} + \frac{(\varepsilon\Delta\dot{X})^2}{2!} \frac{d'^2E(\dot{X})_{nO}}{dX^2_{nO}} + \cdots \end{Bmatrix}$$

$$\ldots \quad \ldots \quad \ldots \quad \ldots \qquad (3.3.10)$$

$\cdot\dot{\cdot}\cdot$ equation (3.3.10) can be written as

$$\{E'(\dot{X})_n\} = \sum_{I=0}^{NP} \frac{(\varepsilon\Delta\dot{X}_n)^I}{I!} \left\{\frac{d^I E'(\dot{X})_n}{d(\dot{X})_n^I}\right\} \quad \ldots \quad \ldots \qquad (3.3.11)$$

where

$$\Delta\dot{X}_n = \dot{X}_{n1} + \varepsilon\dot{X}_{n2} + \varepsilon^2\dot{X}_{n3} + \ldots \qquad \ldots \quad \ldots \qquad (3.3.12)$$

Therefore insert Equation (3.3.11) into Equation (3.3.8) and substitute in Equation (3.3.4) with collecting terms having the same power of ε.

The first four of these would appear as

$$\left[\left[\tilde{\ } M + M_a \tilde{\ }\right]\{\ddot{X}\}_o + \left[[C] + \left[\tilde{\ } 2\dot{U}F_D \tilde{\ }\right]\right]\{\dot{X}\}_o +\right.$$

$$\left. + [K]\{X\}_o - \{F(t)\}\right] +$$

$$\varepsilon\left[\left[\tilde{\ } M + M_a \tilde{\ }\right]\{\ddot{X}\}_1 + \left[[C] + \left[\tilde{\ } 2\dot{U}F_D \tilde{\ }\right]\right]\{\dot{X}\}_1 + [K]\{X\}_1 +\right.$$

$$\left. + \{E'(\dot{X})_o\}\right] +$$

$$\varepsilon^2\left[\left[\tilde{\ } M + M_a \tilde{\ }\right]\{\ddot{X}\}_2 + \left[[C] + \left[\tilde{\ } 2\dot{U}F_D \tilde{\ }\right]\right]\{\dot{X}\}_2 + [K]\{X\}_2 +\right.$$

$$\left. + \{(\dot{X})_1 \frac{dE'}{d\dot{x}}(\dot{X})_o\}\right] +$$

$$\varepsilon^3\left[\left[\tilde{\ } M + M_a \tilde{\ }\right]\{\ddot{X}\}_3 + \left[[C] + \left[\tilde{\ } 2\dot{U}F_D \tilde{\ }\right]\right]\{\dot{X}\}_3 + [K]\{X\}_3 +\right.$$

$$\left. + \frac{1}{2}(\{(\dot{X})_1^2 \frac{d^2E'}{d\dot{x}^2}(\dot{X})_o\} + 2\{(\dot{X})_2 \frac{dE'}{d\dot{x}}(\dot{X})_o\})\right] + \ldots = 0$$

$$\ldots \quad \ldots \quad \ldots \quad \ldots \qquad (3.3.13)$$

If Equation (3.3.12) is to be satisfied identically in ε, each square bracket must separately vanish. This provides a chain of linear problems in which $X_{I+1}(t)$ may be considered as a linear response to an excitation that is a nonlinear function of the previously determined $X_I(t)$.

Applying the perturbation method to the given problem leads to difficulty of deciding how many terms the solution, (3.3.9) should retain in order to fully represent the nonlinearity by a series of the type (3.3.11). Also it is necessary to establish how the computation for higher-order terms is to be executed, and what the permissible values of ε should be.

SEMILINEAR METHOD

This method of analysis involves replacing the actual nonlinear system equation of motion (3.3.4), by a system of linear differential equation. It is possible to have the linear system to be an equivalent system by selecting the constants such that the errors arising from its use are minimal. The efficiency of this method depends, as did the perturbation method, upon the nonlinear strength of the system. If a highly nonlinear system were to be replaced by a linear one with constants selected to minimize the mean-square errors, the minimized errors might be large enough

to cast doubt upon the results. However, if the system is weakly nonlinear in the sense that its equations of motion possess linear portion which dominate the nonlinear portions (as anticipated in Equation (3.3.4)), the minimized replacement errors can be much smaller.

This method was applied to a nonlinear system with deterministic inputs by Krylov and Bogaliubov 1947[106]. The semilinear approach for the problem of offshore structure is developed as follows.

The equation of motion (3.3.4) shows that the nonlinear effects are velocity-dependent which implies that the linear replacing system of equations should be of the form

$$[\tilde{M} + M_a\tilde{}]\{\ddot{X}\} + [\tilde{C}]\{\dot{X}\} + [K]\{X\} + \{e(\dot{X})\} = \{F(t)\} \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.3.14)$$

In Equation (3.3.14) $[\tilde{C}]$ is an equivalent linear damping matrix and $\{e(\dot{x})\}$ is the vector of velocity dependent errors introduced by having the linear system. It is the objective of the semilinear approach to minimize the elements of $\{e(\dot{x})\}$ in some manner.

Prior to discussing the manner in which elements of $\{e(\dot{x})\}$ will be minimized, it is necessary to study $[\tilde{C}]$ in more detail to establish the parameters of this matrix which affect $\{e(\dot{x})\}$. Selection of

$\left[[C] + [\overleftarrow{2\dot{U}_o F_D}]\right]$ is aided by examining Equation (3.3.4) and substituting for $\{E(\dot{X})\}$ from Equation (3.3.6) to obtain an alternative expression of Equation (3.3.4) as

$$[\overleftarrow{M + M_a}]\{\ddot{X}\} + \left[[C] + [\overleftarrow{2\dot{U}_o F_D}] + [\overleftarrow{F_D}][\overleftarrow{\dot{X}}]\right]\{\dot{X}\} + [K]\{X\} = \{F(t)\} \quad \ldots \quad \ldots \quad \ldots \quad (3.3.15)$$

where the matrix $[\overleftarrow{F_D}]$ $[\overleftarrow{\dot{X}}]$ will produce a diagonal matrix. Equation (3.3.15) shows that, the diagonal element of the damping matrix(premultiplying $\{\dot{X}\}$) is a function of $\dot{X}_n$ while the off-diagonal element is a constant, namely C_{mn}. Accordingly, it becomes sensible to take $[\tilde{C}]$ to be given by

$$\tilde{C}_{Mn} = \begin{cases} C_{Mn} & m \neq n \\ \\ \tilde{C}_{Mn} & m = n \end{cases} \quad \ldots \quad \ldots \quad \ldots \quad (3.3.16)$$

where the diagonal element $\tilde{C}_{mn}$ remain to be determined such that $\{e(\dot{x})\}$ is minimized.

By equating Equations (3.3.4) and (3.3.14) and solve for

$\{e(\dot{x})\}$

$$\{e(\dot{x})\} = \left(\left[[C] + \left[\overleftarrow{\ 2\dot{U}_o F_D}\right]\right] - [\tilde{C}]\right|) \{\dot{X}\} + \{E(\dot{X})\}$$

... ... (3.3.17)

As a consequence of Equation (3.3.16) at the nth node point the $\{e_n(\dot{x})_n\}$ depends only upon $\dot{X}_n$, C_{nn}, $2\dot{U}_{on}$ F_{Dn}, $E_n(\dot{X})_n$ and the unknown constant of $\tilde{C}_{nn}$

$$\therefore\ e_n(\dot{X})_n = (C_{nn} + 2\dot{U}_{on}\ F_{Dn} - \tilde{C}_{nn})\dot{X}_n + E_n(\dot{X})_n$$

... ... (3.3.18)

Taking the ensemble average of $e_n^2(\dot{x})_n$, given as

$$\sigma^2_{e_n\ e_n} = < e_n^2(\dot{X}(t))_n > \quad \ldots \quad \ldots \quad \ldots \quad \ldots \qquad (3.3.19)$$

provides N mean-square expressions of the nth node point $\{e_n(\dot{x})_n\}$ which can be minimized by N expressions of the form

$$\frac{\partial \sigma^2_{e_n\ e_n}}{\partial C_{nn}} \qquad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \qquad (3.3.20)$$

Since Equation (3.3.18) shows that $\tilde{C}_{nn}$ is the only unknown parameter affecting $e_n(\dot{x})_n$, substituting from Equations (3.3.18) and (3.3.19) into (3.3.20) leads to N equation for $\tilde{C}_{nn}$ by performing the operations

$$\frac{\partial}{\partial \tilde{C}_{nn}} < ((C_{nn} + 2\dot{U}_{on}\ F_{Dn} - \tilde{C}_{nn})\dot{X}_n + E_n(\dot{X})_n)^2 > = 0 \qquad (3.3.21)$$

Differentiating inside the < > operator is permissible, given that the ensemble average exists, and provides the N optimizing equations for $\tilde{C}_{nn}$ as

$$\tilde{C}_{nn} = C_{nn} + 2\dot{U}_{on} F_{Dn} \frac{\sigma^2 E_n(\dot{X})_n}{\sigma^2 \dot{X}_n \dot{X}_n} \quad \ldots \quad \ldots \quad (3.3.22)$$

with

$$\sigma^2 E_n(\dot{X})_n = < E_n(\dot{X}(t))_n \cdot (\dot{X}(t))_n > \quad (a)$$

$$\sigma^2 \dot{X}_n \dot{X}_n = < (\dot{X}^2(t)_n > \quad (b) \quad (3.3.23)$$

That the coefficients $\tilde{C}_{nn}$ define a true minimum is apparent in the computation leading to

$$\frac{\partial^2 \sigma^2 e_n e_n}{\partial \tilde{C}_{nn}^2} = < 2 \sigma^2 \dot{X}_n \dot{X}_n > \quad 0 \quad \ldots \quad (3.3.24)$$

The value of $\sigma \dot{X}_n \dot{X}_n$ should be provided in order to obtain the new parameter for $[\tilde{C}]$. Obviously this cycle must have a beginning point, however, the selection of it is arbitrary.

The process initiates with a selection of the N diagonal terms ${}^{o}\tilde{C}_{nn}$ of the equivalent damping matrix ${}^{o}[\tilde{C}]$ the pre-superscript o indicates the initial approximation of a quantity, the second approximation value of ${}^{1}\tilde{C}_{nn}$ will be

$$ {}^{1}\tilde{C}_{nn} = {}^{o}\tilde{C}_{nn} + \sqrt{\frac{8}{\pi}}\, F_D^o\, \sigma_{\dot{X}_n}\, \dot{X}_n \quad \ldots \quad \ldots \quad \ldots \quad (3.3.25) $$

The nonlinear hydrodynamic damping force $\{E(\dot{X})\}$ is a monotonically increasing function of $\dot{X}$, which means that if the initial value of ${}^{o}\tilde{C}_{nn}$ underestimates the true value of $\tilde{C}_{nn}$, ${}^{o}\sigma_{\dot{x}_n \dot{x}_n}$ will be in excess of its true value so that provides ${}^{1}\tilde{C}_{nn}$ as a magnitude overestimate of the true value of $\tilde{C}_{nn}$. The reverse would be true. This process continues until convergence is defined, for all the N main-diagonal members of $[\tilde{C}]$. The number of cycles requires to establish convergence depends upon the initial value of ${}^{o}\tilde{C}_{nn}$ and the strength of the nonlinearity.

This method requires changing of damping matrix at each step of solution until the convergence is defined.

For the present study, the investigation was done for small amplitudes of vibration so that nonlinearity in Equation (3.3.4) is very small. Therefore the problem is solved by the following two methods.

METHOD A

In this method the governing nonlinear equation of motion is solved by the direct approach method. This method is as follows:

1) The governing equation of motion (3.3.4) can be rearranged as follows:

$$[M + M_a]\{\ddot{X}\} + \left[[C] + [2\dot{U}_o\ F_D]\right]\{\dot{X}\} +$$

$$[K]\ \{X\} = \{F(t)\} - \{E(\dot{X})\} \quad \ldots \quad \ldots \qquad (3.3.26)$$

2) Neglecting the nonlinear terms $\{E(\dot{X})\}$ for the first estimate. A set of linearized equations are solved.

3) Use first solution and calculate the nonlinear term $\{E(\dot{X})\}$

4) Resolve Equation (3.3.26) with the calculated nonlinear term obtained from the first solution. Obtain a second solution and recalculate the nonlinear term $\{E(\dot{X})\}$ using the second estimate of $\dot{X}$

5) Repeat the process until convergence of successive iteration is obtained.

The value of C_D and C_M obtained from the test of free structure were used in forming of Equation (3.3.26).

METHOD B

This method is based upon the assumption that for small displacement of the structure, whether the object moves in still water or the water moves about the object does not affect the value of the hydrodynamic damping.

The equation (3.3.4) will be

$$[M]\{X\} + [C^w]\{\dot{X}\} + [K]\{X\}$$

$$= [F''_m] + [F'_m]\{\ddot{U}\} + [F''_D]\{\dot{U}|\dot{U}\}$$

... ... (3.3.27)

where $[C^w]$ is the damping matrix obtained by using the damping coefficient measured in water

$[F''_m]$ and $[F''_D]$ are calculated by using the value of C_M and C_D obtained from the test of fixed structure.

3.3.3 SOLUTION OF THE SYSTEMS

The solution of the linearized system (Method (A) and Method (B)) is based upon the possibility of treating the system as having linear input-output characteristic in the sense that the principle of superposition is valid. The normal mode superposition approach is used to calculate the response of the structure. The response of the structure is given by

$$\{X\} = [\phi] \{Y\} \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.3.28)$$

where the modal matrix $[\phi]$ is deduced for the undamped free vibration given by the eigen value problem.

$$[\overline{M + M_a}] \{-\omega^2 X\} + [K] \{X\} = \{O\}$$

$$[\underline{\omega^2}][\overline{M + M_a}][\phi] = [K]\{\phi\} \quad \ldots \quad (3.3.29)$$

Equation (3.3.26) or (3.3.27) is premultiplied by the transpose of the modal matrix, and the displacement $\{X\}$ expressed in terms of the normal co-ordinate vector $\{Y\}$, to yield the relations

$$[\overline{M^*}] \{\ddot{Y}\} + [\hat{C}] \{\dot{Y}\} + [\overline{K^*}] \{Y\} = \{P^*\}$$

$$\ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.3.30)$$

Since the system equations comprise a self-adjoint eigen value problem and orthogonality property of the normal modes exists, where

$$[M^*] = [\phi]^T [M + M_a][\phi] \quad \text{(For Method A)}$$

$$\text{or} \quad [\phi]^T [M] [\phi] \quad \text{(For Method B)} = \text{Generalized Mass Matrix}$$

... ... (3.3.31)

$$[K^*] = [\phi]^T [K] [\phi] \quad \text{(For Both methods)} = \text{Generalized Stiffness Matrix}$$

... ... (3.3.32)

$$[\hat{C}] = [\phi]^T [[C] + [2\dot{U}_o F_D]][\phi] \quad \text{(For Method A)}$$

$$\text{or} \quad [\phi]^T [C''] [\phi] \quad \text{(For Method B)}$$

$$= \text{Generalized Damping Matrix} \quad (3.3.33)$$

$$\{P^*\} = [\phi]^T ([[F_m] + [F'_m]] \{\ddot{U}_o\} + [F_D] \{\dot{U}_o\}$$

$$+ [F_D] \{\dot{X}|\dot{X}|\}) \quad \text{(For Method A)}$$

$$\text{or} = [\phi]^T ([[F''_m] + [F'_m]] \{\ddot{U}_o\} +$$

$$+ [F''_D] \{\dot{U}_o\}) \quad \text{(For Method B)}$$

... ... (3.3.34)

$\{P^*\}$ = Generalized Force vector.

The coupled damping matrix $[\hat{C}]$ is symmetric but is not a diagonal matrix. The coupled damping matrix is decoupled by assuming that the damping matrix $[C]$ satisfied the modal orthogonality condition

$$[\phi]^T [C] [\phi]_m = O \quad (n \neq m) \quad \ldots \quad \ldots \quad (3.3.35)$$

$\therefore$ Equation (3.2.27) become

$$[M^*] \{\ddot{Y}\} + [C^*] \{\dot{Y}\} + [K^*] \{Y\} = \{P^*\} \quad (3.3.36)$$

Equation (3.3.36) therefore represents N uncoupled linear second order differential equations. These equations can also be put in the familiar form.

$$\ddot{Y}_n + 2\omega_n \,'\xi_n \dot{Y}_n + \omega_n^2 Y_n = \frac{P_n^*}{M_n^*} \quad n = 1, 2, \ldots, N \quad (3.3.37)$$

where

$$'\xi_n = \frac{C_{nn}^*}{2M_n^*\omega_n} \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.3.38)$$

In this derivation of the normal- coordinate equations of motion it has been assumed that the normal-coordinate transformation serves to uncouple the damping forces in the same way that it uncouples the mass and stiffness.

The conditions under which this uncoupling will occur is the Rayleigh damping matrix, that is, the form of damping matrix in which Equation (3.3.35) applied.

Rayleigh showed that a damping matrix of the form

$$[C] = \alpha_o[M] + \alpha_1 [K] \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.3.39)$$

in which α_o and α_1 are arbitrary proportionality factors, will satisfy the orthogonality condition. The method of determination of α_o and α_1 will be shown later in this chapter.

The solution of Equation (3.3.37) can be obtained in the time domain or the frequency domain. Having determined $Y_n(t)$, n = 1, 2, 3, ..., N the displacement {X} can be found through Equation (3.3.28).

3.3.4 NUMERICAL ANALYSIS

The analysis of the structures which used in this study, were carried out by using finite element method. The computer program SAPIV was available to analyse the dynamic response of the structures in the time domain. The structure's response was calculated by the mode superposition technique, using the first three modes only.

The structural system was analysed by using the tangent pipe element. The tangent pipe element (Figure (3.3.2) can represent a straight segment; the elements require a uniform section and uniform material properties. The member stiffness matrix account for bending, torsional, axial and shearing deformations. The types of structure loads contributed by the pipe elements include gravity loading in the global directions, forces and moments acting at the member ends (i, j).

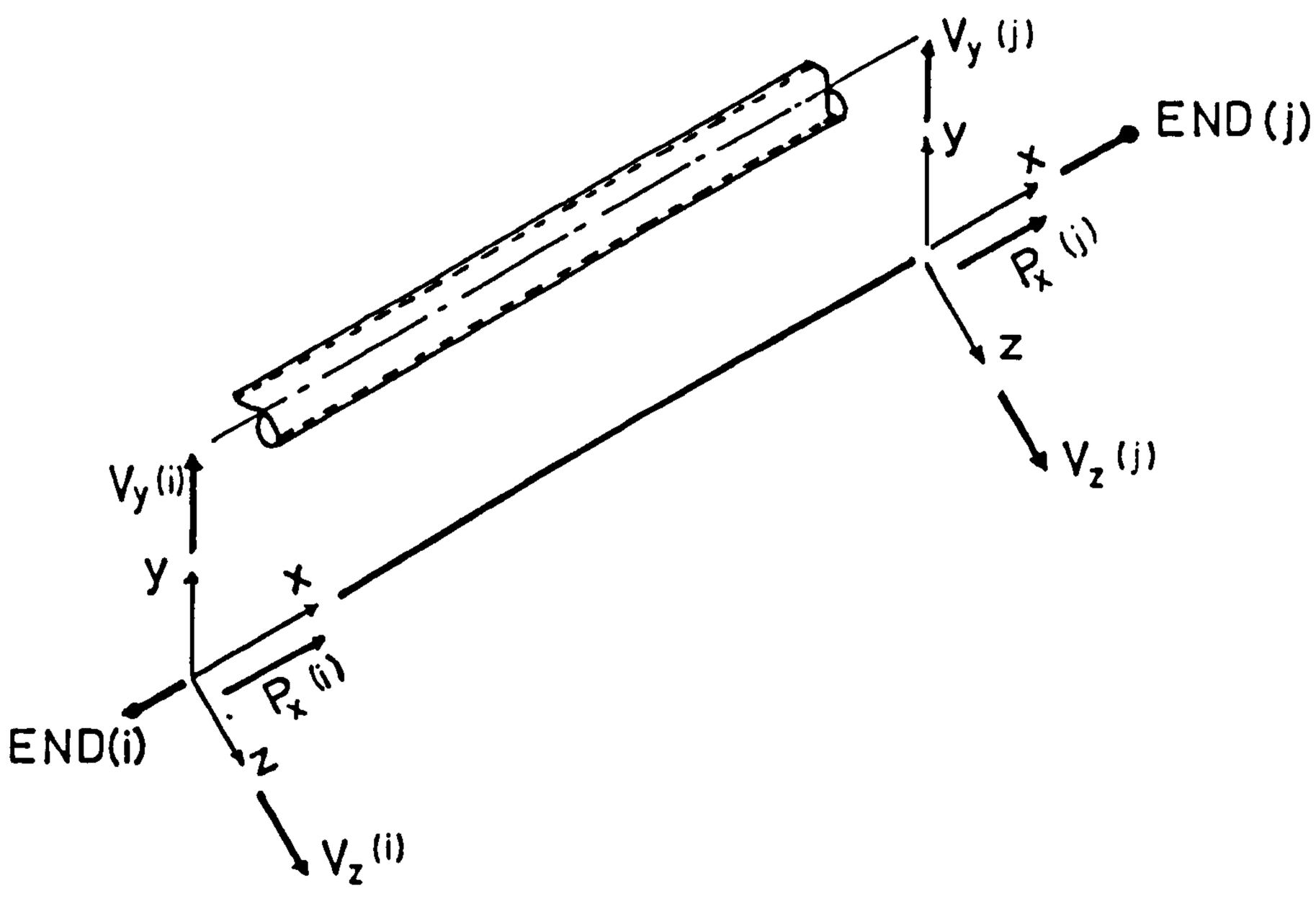


Figure (3.3.2) - The Tangent Pipe Element.

3.4 THEORETICAL FORMULATION OF THE EXPERIMENTS

The parameters required in the analysis of the problem were obtained experimentally. The theory and assumption behind the analysis of experimental data are shown below.

3.4.1 MODELLING OF THE TESTED STRUCTURES

All the tested structures were divided into six elements. The first element had the length from the still water level to the tip of the structure and the other five elements were obtained by dividing the length from the still water level to the base of the structure. The element mass was lumped at the element ends equally. Each of these lumped mass has two degrees of freedom: (1) horizontal translation and (2) rotation. The response of the element is discretized in step-wise fashion, i.e. the response of each lumped mass represents the response of the element half way to the upper and lower lumped masses (see Figure (3.4.1)).

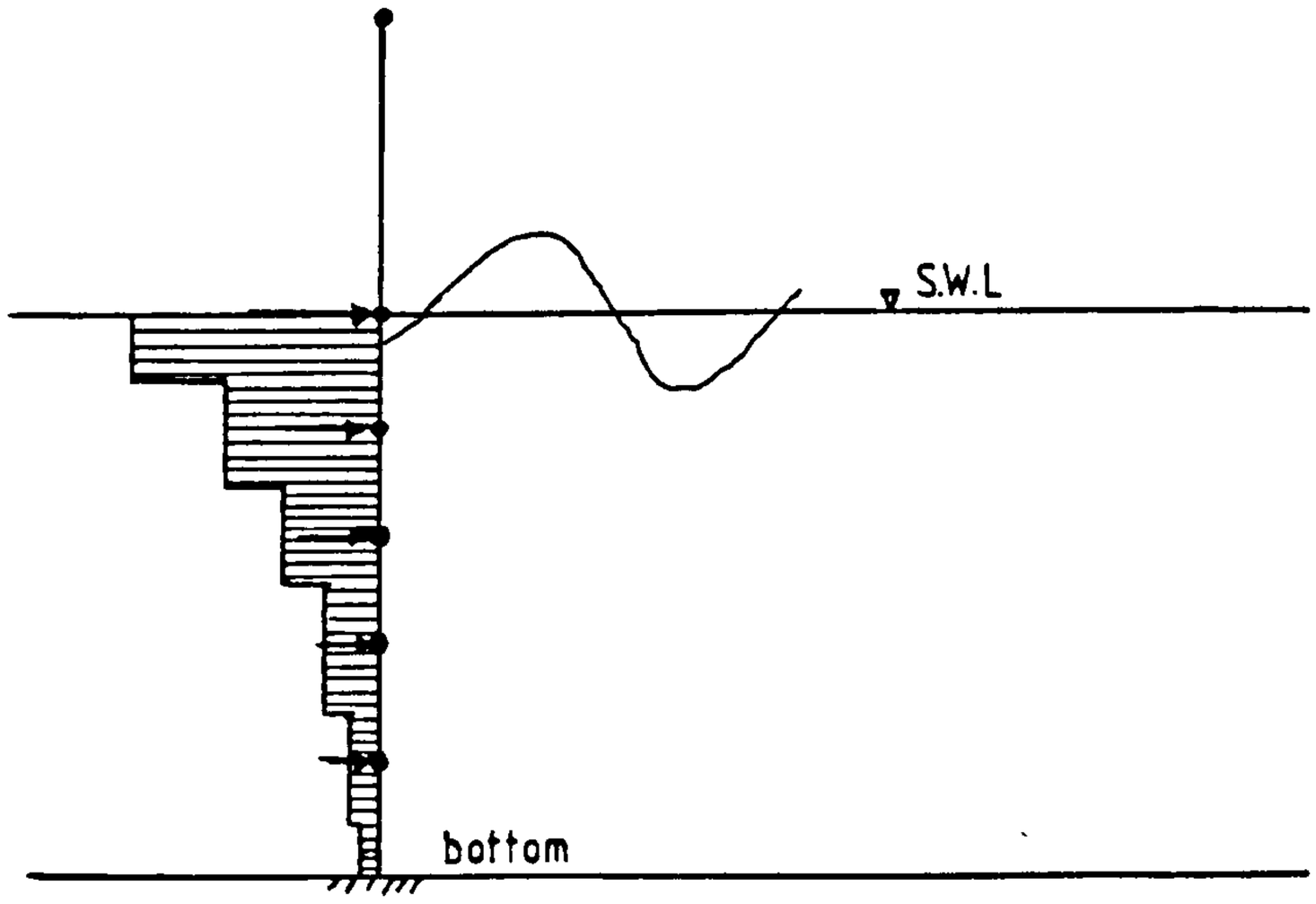


Figure (3.4.1) - Discretization of Response and Force.

The fluid orbital velocities and acceleration are discretized in the same stepwise manner as the response. Morison's equation and the modified Morison's equation are used for the evaluation of the wave force which are discretized at the lumped mass positions. Each force is evaluated for the real structure configuration and represent the contribution from half of the panels below and above that position. The net moment created by the forces about a lumped mass centre is neglected, as it will be relatively small.

3.4.2 DIGITIZATION OF THE DATA

The experimental data obtained were recorded in analog format. For numerical solutions this data need to be digitized in a way that it simulate the obtained analog data.

In order to maintain a proper relationship between the phase and amplitude of a given variable a suitable frequency of digitization or digitization interval had to be chosen. Otherwise aliasing could occur which is a potential source of error in an analog to digital data conversion (see Bendat and Piersol 1971[107]). The maximum frequency reasonable is called Nyquist frequency and is defined as

$$f_{NY} = \frac{1}{2\Delta t} \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.4.1)$$

where Δt is the interval of digitization. The frequencies in the original data above this cut off frequency are folded back into the frequency range from zero to f_{NY}. Therefore, in order to choose a suitable Δt the maximum frequency encountered in the experiments had to be known. The observation of the analog records showed that the highest frequencies in the range of 3-4 Hz.

In general, it is a good rule to select f_{NY} to be greater than the maximum anticipated frequency. Therefore a Nyquist frequency of 25 Hz was chosen from which an interval of digitization of 0.02 seconds resulted. Any digitization below this value would have given correct results. The smaller of this interval will require more computing time.

3.4.3 DETERMINATION OF THE MEASURED WAVE FORCE

At any time t, the force acting on the measured level of the structure can be obtained by integration of the pressure and the shear distribution on the object surface (see Figure (3.4.2)).

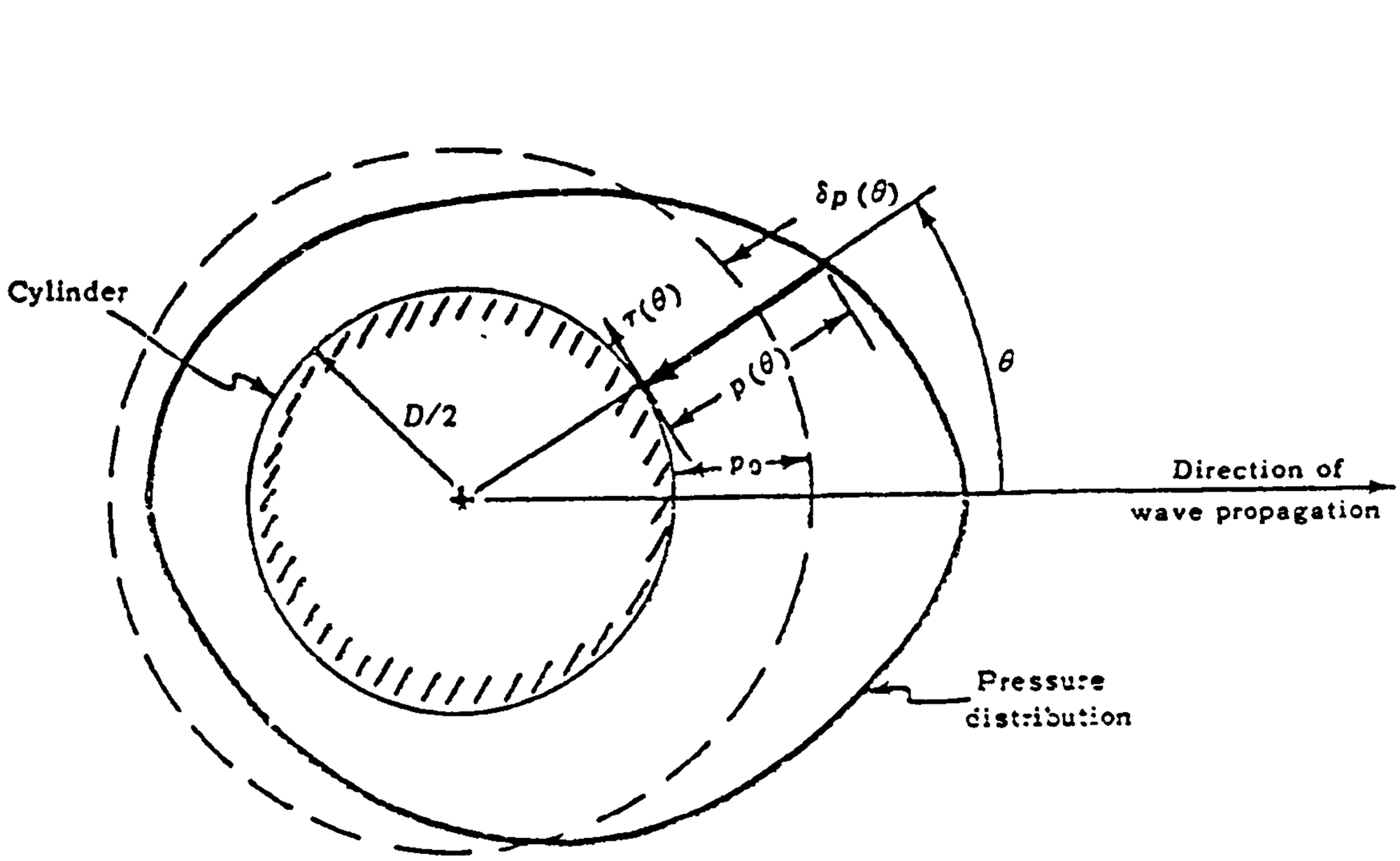


Figure (3.4.2) - Pressure and Shear Distribution about a Circular Cylinder.

$$dF_m = \Big[\int_0^{2\pi} \frac{D}{2}(P_o + \delta P) \cos\theta \; d\theta$$

$$+ \int_0^{2\pi} \frac{D}{2} \tau(\theta) \sin\theta \; d\theta \Big] ds$$

$$dF_m = \Big[\frac{DP_o}{2} \int_0^{2\pi} \cos\theta \; d\theta + \int_0^{2\pi} \frac{D}{2} \delta P(\theta) \cos\theta \; d\theta$$

$$+ \int_0^{2\pi} \frac{D}{2} \tau(\theta) \sin\theta \; d\theta \Big] ds$$

$$dF_m = \Big[\text{zero} + \int_0^{2\pi} \frac{D}{2} \delta P(\theta) \cos\theta \; d\theta$$

$$+ \int_0^{2\pi} \frac{D}{2} \tau(\theta) \sin\theta \; d\theta \Big] ds \qquad \ldots \quad \ldots \qquad (3.4.2)$$

As the structure has a smooth surface the force attributed to the shear distribution is negligible

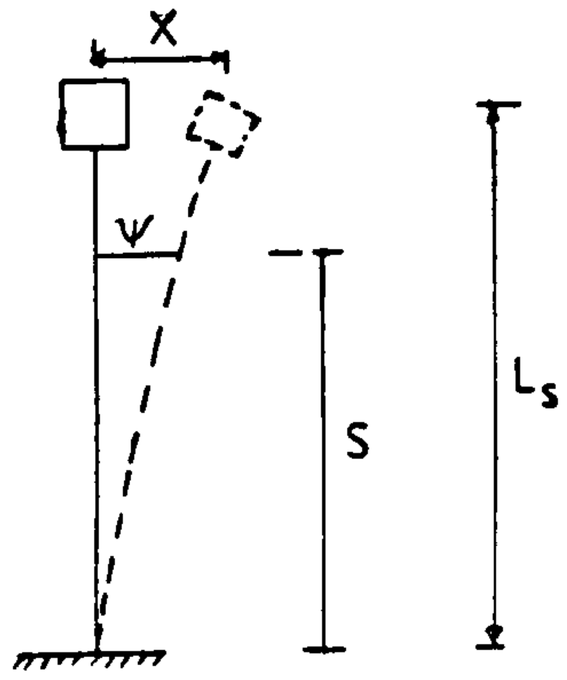
$$\therefore \quad dF_{m(t,z)} = \int_0^{2\pi} \frac{D}{2} \delta P(\theta) \cos\theta \; d\theta \; ds \qquad \ldots \qquad (3.4.3)$$

3.4.4 DETERMINATION OF THE VELOCITY AND ACCELERATION OF THE STRUCTURE AT ANY LEVEL

The velocity and acceleration of the structure at any level (the level where the pressure were measured) were required for the determination of the force coefficients C_D' and C_M' when the modified Morison's equation was used.

The calculation of the velocity and acceleration of the structure at any level by knowing the tip displacement is based upon the assumption that, if the load distribution is similar to the deflected shape of the first mode shape and if the load frequency is smaller or equal to the first mode frequency of the structure, then only the first mode of vibration exists.

Timoshenko 1956[108] shows a solution for elastic curve, ψ as an approximation formed from a truncated series of the sum of sinusoidal shapes, each contributing to ψ for the first mode shape. The result is



$$\psi = \frac{PL_S^3}{3EI} \quad \frac{1}{1-\beta'} \left(1 - \cos \frac{\pi s}{2L_s}\right) \quad \ldots \quad \ldots \quad (3.4.4)$$

where I = sectional moment of inertia

$$\beta' = \frac{4WL_s^2}{\pi^2 EI}$$

when $S = L_S$

$$\psi = \frac{PL_s^3}{3EI} \quad \frac{1}{1-\beta'}$$

which represent the tip displacement

$$\psi = X\left(1 - \cos \frac{\pi s}{2L_s}\right) \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.4.5)$$

Equation (3.3.5) represent the deflection at any level of the structure. As the tip displacement of the structure due to the wave forces is periodic function, it is possible to represent X by the fast Fourier transforms.

$$X = \sum_n a_n'' \sin\left(\frac{2\pi n}{T} \Delta t\right) + b_n'' \cos\left(\frac{2\pi n}{T} \Delta t\right)$$

$$\ldots \quad \ldots \quad \ldots \quad \ldots \quad (3.4.6)$$

The velocity $\dot{X}$ at any level of the structure can be obtained by differentiation of the displacement with respect to the time

$$\dot{X} = \frac{\partial X}{\partial t}$$

$$\dot{X} = \left(\Sigma a''_n \frac{2\pi n}{T} \cos\left(\frac{2\pi n}{T} \Delta t\right) - b''_n \frac{2\pi n}{T} \sin\left(\frac{2\pi n}{T} \Delta t\right)\right)\left(1-\cos \frac{\pi s}{2L_s}\right.$$

... ... (3.4.7)

Also the acceleration $\ddot{X}$ will be obtained by

$$\ddot{X} = \frac{\partial \dot{X}}{\partial t}$$

$$\ddot{X} = \left(\Sigma - a''_n\left(\frac{2\pi n}{T}\right)^2 \sin\left(\frac{2\pi n}{T} \Delta t\right) -\right.$$

$$\left. - b''_n\left(\frac{2\pi n}{T}\right)^2 \cos\left(\frac{2\pi n}{T} \Delta t\right)\right)\left(1 - \cos \frac{\pi s}{2L_s}\right)$$

... ... (3.4.8)

3.4.5 EVALUATION OF THE CONSTANT α_o AND α_1 FOR THE DAMPING MATRIX

The Rayleigh damping matrix (Equation (3.3.39)) should satisfy the orthogonality condition. This can be demonstrated by applying the orthogonality operation which is

$$[\phi]_m^T \quad [M] \quad [\phi]_n = O \quad m \neq n \quad (a)$$

$$[\phi]_n^T \quad [K] \quad [\phi]_n = O \quad m \neq n \quad (b) \qquad \ldots \qquad (3.4.9)$$

in both side of the equation

$$[\phi]^T [C][\phi] = \alpha_o [\phi]^T[M] [\phi] + \alpha_1 [\phi]^T [K] [\phi]$$

$$\ldots \quad \ldots \quad \ldots \quad \ldots \qquad (3.4.10)$$

as each term at the right hand side of the above equation satisfy the orthogonality condition, therefore the left hand side satisfy the orthogonality condition.

However it can be shown that an infinite number of matrices formed from the mass and stiffness matrices also satisfy the orthogonality condition:

$$[\phi]_m^T \left[[M]\ [M]^{-1}\ [K]\right]^b [\phi]_n = 0 \quad -\infty < b < \infty$$

... ... (3.4.11)

where $[M]^{-1} = \frac{1}{[M]}$

The two basic relationships Equations (3.4.9a) and (3.4.9b) are given by exponents b = 0 and b = 1 in Equation (3.4.11).

.

Thus the damping matrix can also be made up of combination of these. In general, then, the orthogonal damping matrix may be of form

$$[C] = [M] \sum_b \alpha_b \left[[M]^{-1}\ [K]\right]^b = \sum_b [C_b] \quad \ldots \quad (3.4.12)$$

With this type of damping matrix it is possible to compute the damping influence coefficients necessary to provide a decoupled system having any desired damping ratios in any specified number of modes. For each mode n, the generalized damping is given by Equation (3.3.8).

$$[C_n^*] = [\phi]_n^T [C]\ [\phi]_n = 2\ \xi_n \omega_n\ [M^*]_n \quad \ldots \quad (3.4.13)$$

but if [C] is given by Equation (3.4.12), the contribution of b in the series to the generalized damping is

$$[C_b]_n = [\phi]_n^T [C_b] [\phi]_n = \alpha_b [\phi]_n^T [M] [[M]^{-1} [K]]^b [\phi]_n$$

... ... (3.4.14)

$$\therefore [\phi]_n^T [M] [[M]^{-1} [K]]^b [\phi_n] = [\omega_n^2]^b [M^*]_n$$

... ... (3.4.15)

therefore

$$[C_b]_n = \alpha_b [\omega_n^2]^b [M^*]_n$$

on this basis, the damping matrix associated with any mode n is

$$[C_n^*] = \sum_b [C_b]_n = \Sigma \alpha_b [\omega_n^2]^b [M^*]_n \quad \therefore \ldots \quad (3.4.16)$$

Equating the right hand side of Equations (3.4.13) and (3.4.16)

$$\therefore \xi_n = \frac{1}{2\omega_n} \Sigma \alpha_b \omega_n^{2b} \quad \ldots \ldots \ldots \quad (3.4.17)$$

Equation (3.4.17) in matrix form with the first two modes only is

$$\begin{Bmatrix} \xi_1 \\ \xi_2 \end{Bmatrix} = \frac{1}{2} \begin{bmatrix} 1/\omega_1 & \omega_1 \\ 1/\omega_2 & \omega_2 \end{bmatrix} \begin{Bmatrix} \alpha_o \\ \alpha_1 \end{Bmatrix} \quad \ldots \ldots \ldots \quad (3.4.18)$$

$$\therefore \quad \alpha_1 = \frac{\frac{\omega_1}{\omega_2}\xi_1 - \xi_2}{\frac{1}{2}(\frac{\omega_1}{\omega_2} - \omega_2)}$$

$$\therefore \quad \alpha_o = \frac{\frac{\omega_2}{\omega_1}\xi_1 - \xi_2}{\frac{1}{2}(\frac{\omega_2}{\omega_1^2} - 1/\omega_2)}$$

The parameter required to solve the above equations for α_1 and α_o are known except ξ_2 (the damping coefficient of second mode). ξ_2 was assumed to be equal to ξ_1 (see Wilson and Pensica 1972[109]).

CHAPTER FOUR

EXPERIMENTAL TECHNIQUE

The formula used in the estimation of the hydrodynamic forces on structure with relative small diameter has empirical parameters C_D and C_M which vary with the flow field and structure characters, any investigation dealing with water wave-structure interaction involves extensive experimental work before a concrete solution can be obtained.

The aim of this experimental work is to find the relation between the force coefficients when the structure is prevented from vibration and when it is free to vibrate, under different wave and structure characteristics.

4.1 EXPERIMENTAL FACILITIES

4.1.1 WAVE TANK

The experiments were preformed in a rectangular water tank. The dimensions of the cross section are 0.75 x 0.75 metre and 18.0 metre long. The side walls are 0.0095 metres thick glass which is suitable for visually observing the model. The water tank had a regular wave generator installed at one end with a sloping beach wave observer at the other end. At this end an extension of dimension 3.00 metre long and 1.25 x 1.25 cross section connected to the tank. This extension was built in order to help the beach to absorb the waves.

This regular wave generator is a wedge type, the wedge is connected to a 0·75 horse power motor with different speeds of rotation. The motor has a digital counter to control the speed of rotation. By changing the speed of rotation, the frequency of the displaced water by the wedge changes giving a different wave period and length. Also the wedge arm can be adjusted to give a different submerged volume of the wedge which effects the volume of displaced water, consequently effecting the wave amplitude.

From the above and also as the water depth in the tank can be changed, the wave characteristics (d, L, T, H) can be changed in order to produce waves of differing characteristics for the tests.

It is desirable to eliminate cross-waves at all times and all frequencies of the tests; so two-dimensional analysis can be carried out without discrepancies, there is a good case for mounting fins against the face of the wavemaker. The fin was a mesh, in the shape of Λ placed at a distance 1·00 metre from the wavemaker the width of the channel, the angle between the Λ shape is adjustable to give the best elimination for each wave frequency. This fin improve the efficiency in eliminating the cross-waves for the range of wave frequency 0·75 - 1·25.

Reflected waves also prove undesirable though they are often not detected or allowed for. Their energy tends to increase in time to a limiting state, allowing them to exert a corresponding influence on the experimental results. The function of the beach is simply to absorb all wave energy incident upon it as reflections will produce problems. The fact that no beach is a fully efficient absorber implies that the beach design must be related to the wave climate required for the experimental programme, and the amount of reflection that can be tolerated. The beach used in the tank was made of wood sheet having the width of the tank and maximum slope of $9 \cdot 925^{o}$. The beach was efficient in absorbing the waves as defined by a "reflection coefficient", this being the ratio of the amplitudes of the primary reflected and primary incident waves.

The position of the tested structure was half way between the wave generator and beach, approximately 8·5 metre away from either of them to minimize any disturbance caused either by wave generator or the beach.

Figure (4.1.1) shows the wave tank dimension and the position of the beach and wavemaker.

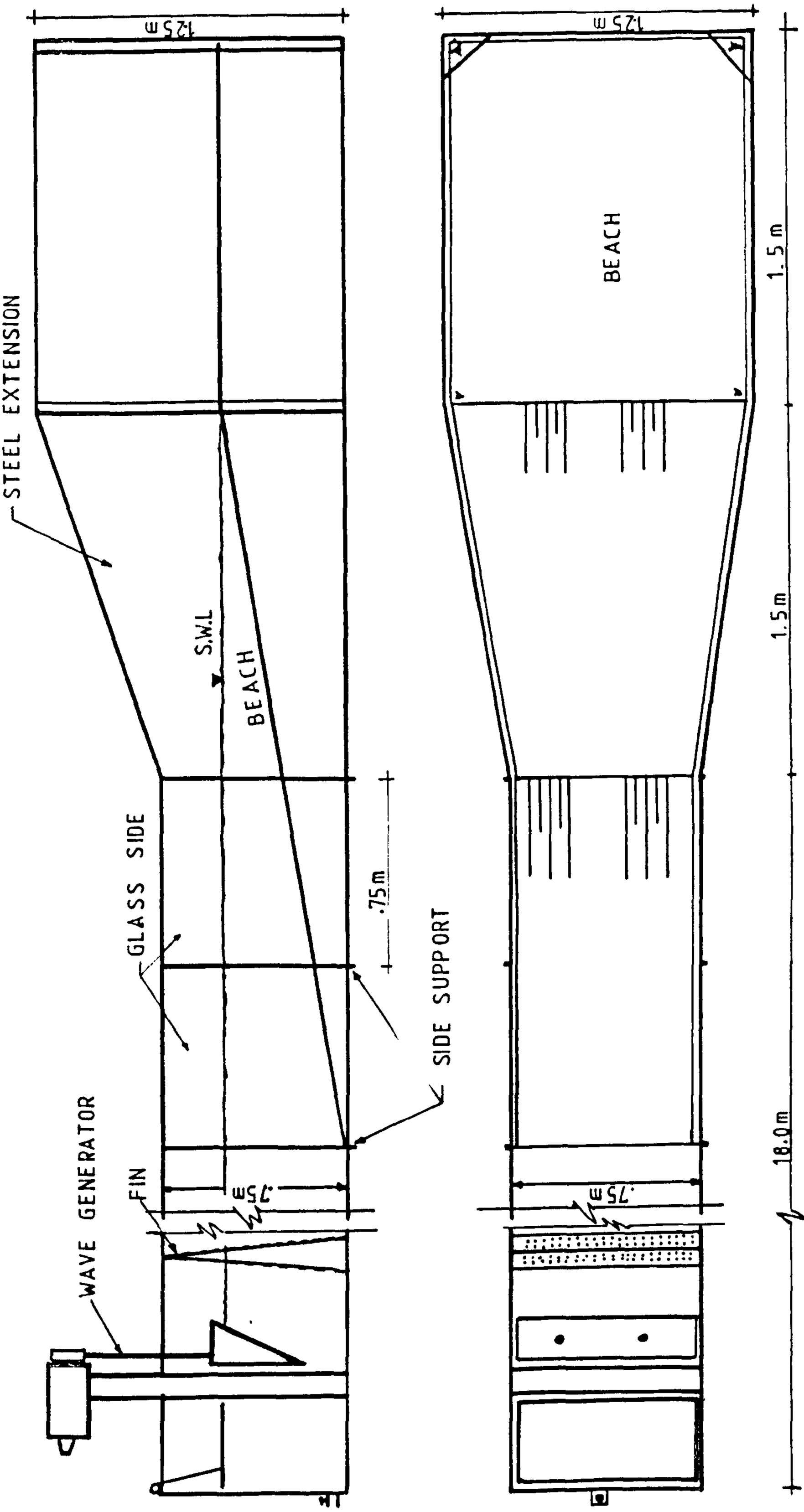


Figure (4.1.1) - Wave Tank.

4.1.2 TESTED STRUCTURES

Two sets of single circular cylinder cantilever structure were tested. A cantilever was chosen because it represents the structure mechanism of the prototype structure, the cross-sectional shape of the structure was chosen to be circular in shape as it is the most common shape in the offshore structures.

The first set consists of two group of structures. The structures of this set were plastic Acrylonitril Butadiene Styrene (Durapipe ABS) tube material. This material was chosen because it afforded a cylinder that could be fixed at the bottom and also has a low modulus of elasticity which decreases the natural frequency of vibration.

The second set consists of one group of structures. The structures were made of Aluminium tube material. This material was chosen to compare between the behaviour of the first and second sets i.e. to compare the behaviour of two different dynamic properties.

4.1.2.1 DURAPIPE ABS STRUCTURES

Four different structures were tested. These structures can be divided into two groups according to the structure diameter and modulus of elasticity i.e. relatively stiff and relatively flexible.

Group one consists of two structures (A) and (B), the structures consist of a 0·11 metre outside diameter tube. The tube was glued at the bottom end to a cap which was rested and glued to a heavy pipe flange with a leader. The pipe flange has six 0·0125 metre bolt holes, equally spaced to fit the imbedded six 0·0125 metre bolts in the heavy perspex sheet of a dimension of 0·75 x 1·5 metre with 0·0125 metre thickness placed on the tank floor at the mid part of the tank. Silicone vacuum grease was applied between the contact surfaces to obtain a water-tight fit.

At one level of the structure eight 0·006 metre diameter holes at equal space of the cross section ($\theta = 45^{o}$) were made. At each of these holes a plastic PVC (hard) tube having L shape of 0·03 x 0·01 metre in dimension where the short leg were fixed to the cylinder wall and filed to give smooth surface while the long leg of the tube were connected to one end of 0·75 metre long rubber tube of 0·005 metre inside diameter, the other end of the rubber tube come out from the structure top so it can be connected to pressure transducer.

A perspex top cap with eight holes to pass the rubber tubes was placed in the top of the structure. It was used to prevent the structure from vibrating by the use of the clamp. The parts of the structure before assembly were shown in Figure (4.1.2).

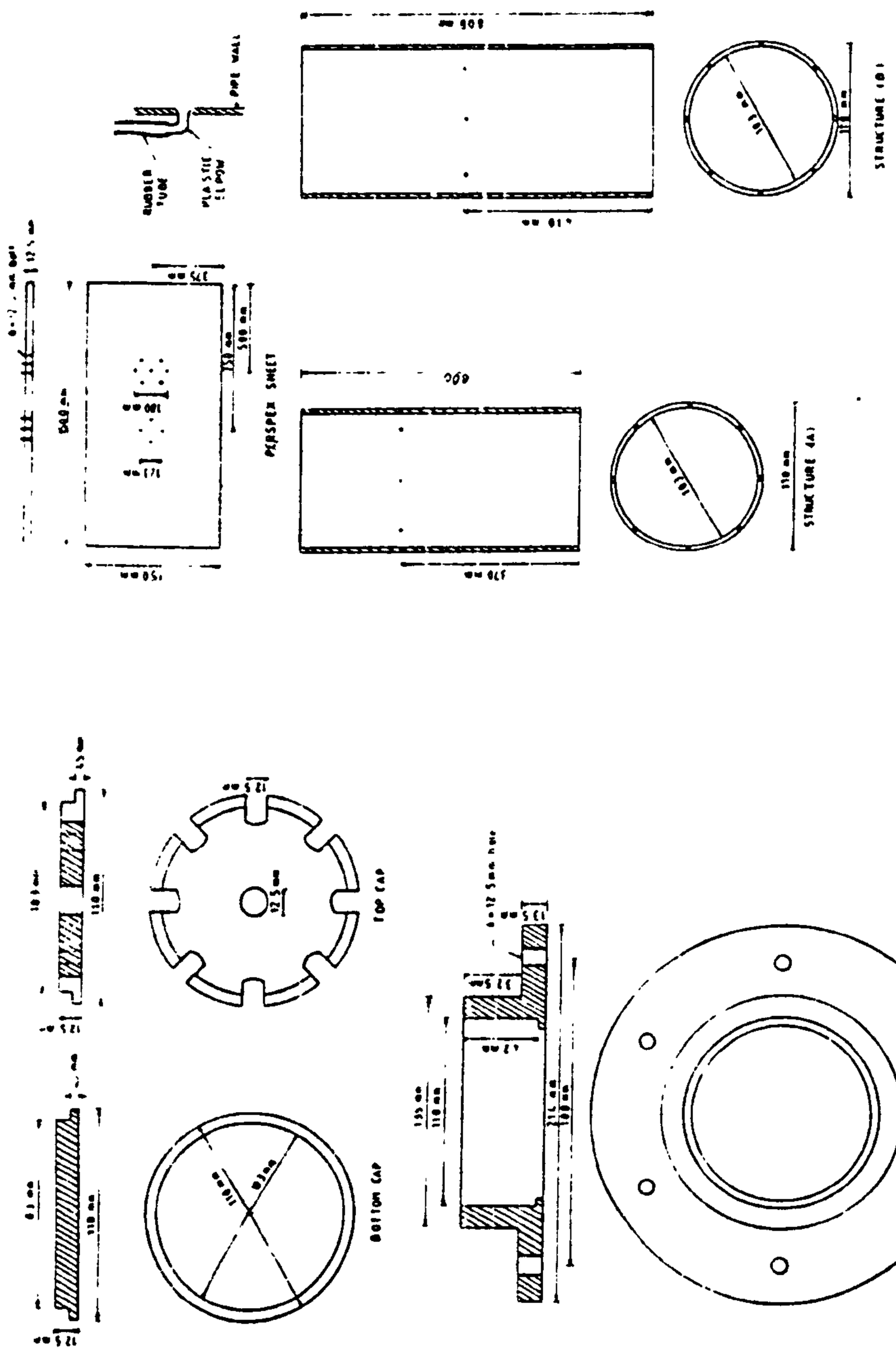


Figure (4.1.2) - The Parts of the Structures (A) and (B) Before Assembly.

Group two consists of two structures (C) and (D). The structures consist of 0.0605 metre outside diameter tube and similar parts as in group one with difference in dimension. See Figure (4.1.3).

At two levels of the group two structure were eight 0.004 metre diameter holes. A plastic PVC (soft) tube instead of the rubber tube was used as there is smaller space inside the tube structure than those of group one.

Figures(4.1.4) and (4.1.5) show the structures of groups one and two after they have been assembled.

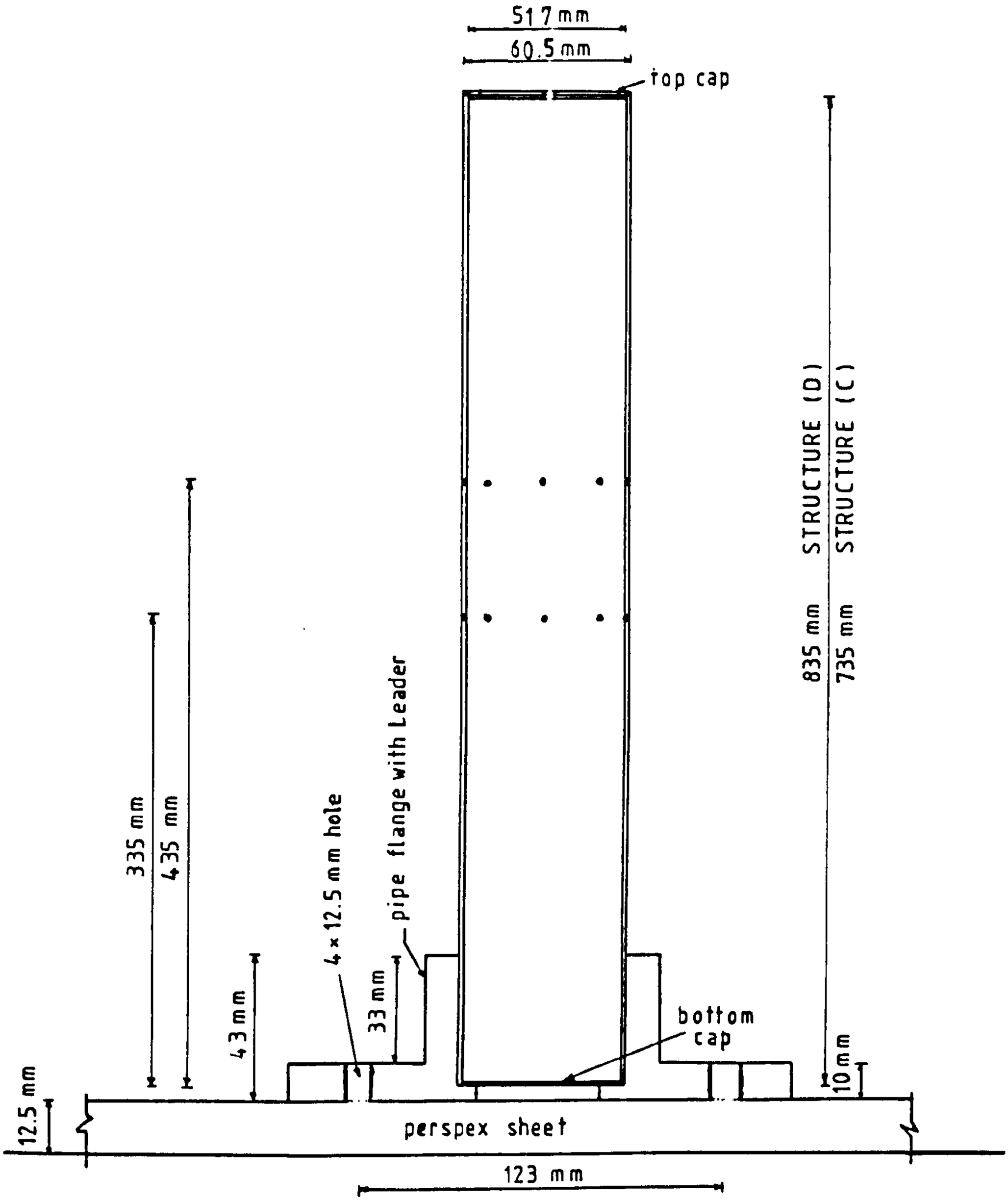


Figure (4.1.3) - Dimension of Structures (C) and (D).

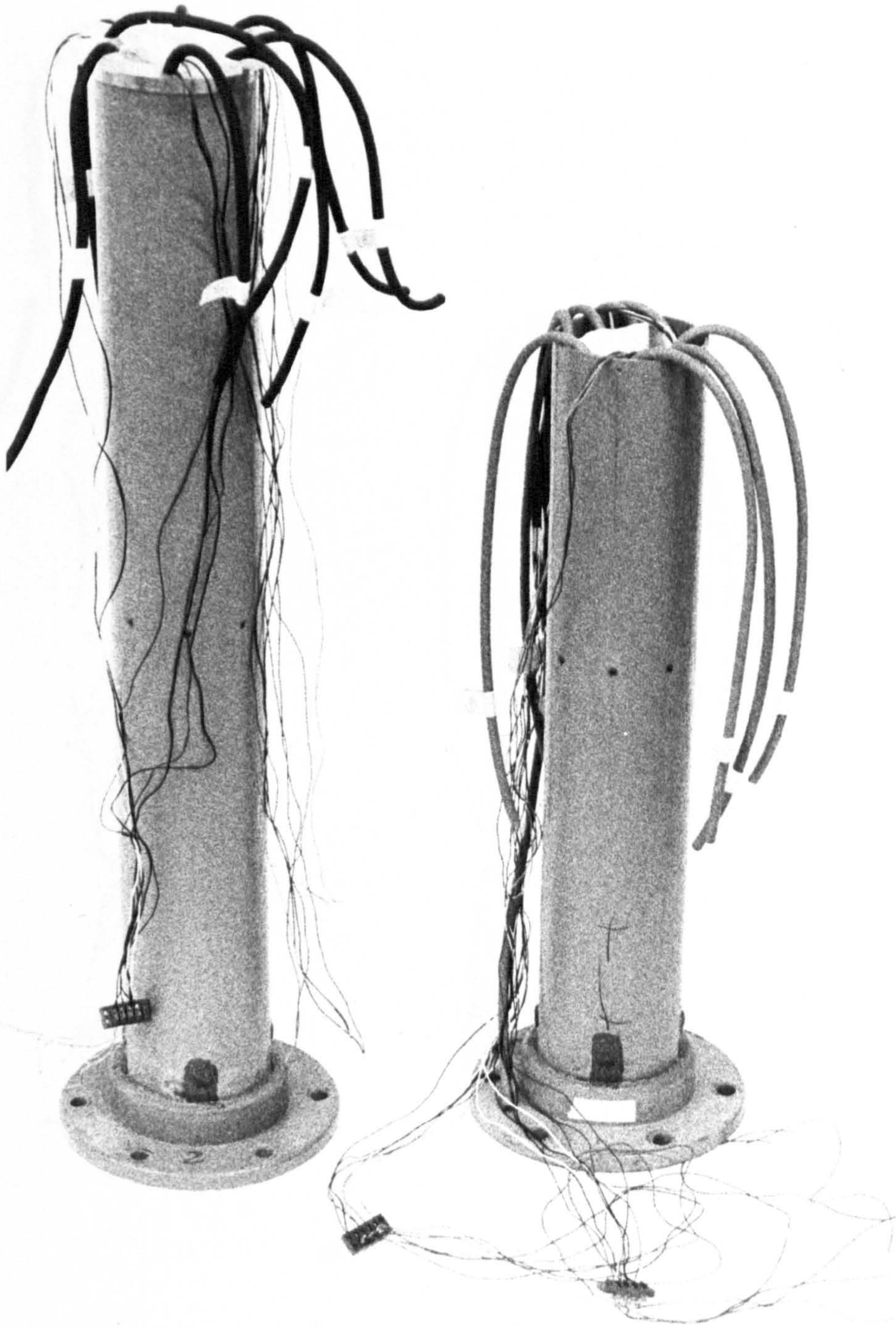


Figure (4.1.4) - Structures (A) and (B).

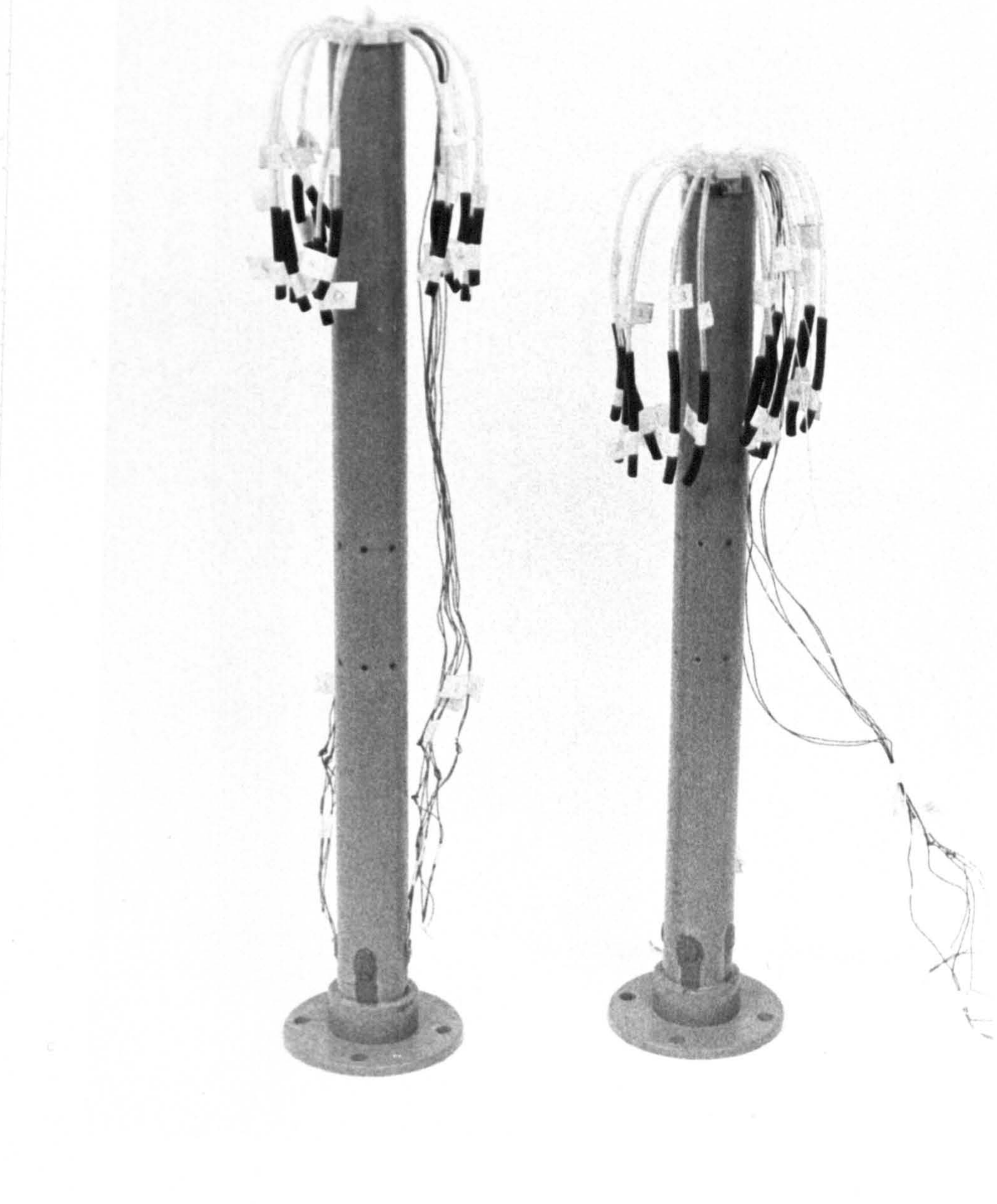


Figure (4.1.5) - Structures (C) and (D).

4.1.2.2 ALUMINIUM STRUCTURES

A third group consists of two aluminium structures. Structure (E) and structure (F) each has 0.0574 metre diameter, structure (E) was 0.735 metre high and structure (F) was 0.835 metre high. A circular PVC (hard) top cap having a mass of 520 grammes was mounted at the top of the structure. It was used as a concentrated mass above the water line so that the structure would behave like a one-degree-of-freedom system when subjected to the wave, thus representing the platform of the offshore structure.

Four masses were made of steel balls contained in a plastic bag each having a mass of 505 grammes. These four masses provided two possible combinations of masses in increments of 1010 grammes. The mass was put in the top cap, the top cap was rigidly fixed by four 0.0125 metre screws to the structure so that the top cap would not wobble with respect to the tube. The top of the cap was level with the structure top.

The eight 0.004 metre diameter holes were made at three different levels of the structure. Figure (4.1.6) shows the parts and dimensions of the structures of this group. The masses and also structure (E) after assembly are shown in Figure (4.1.7).

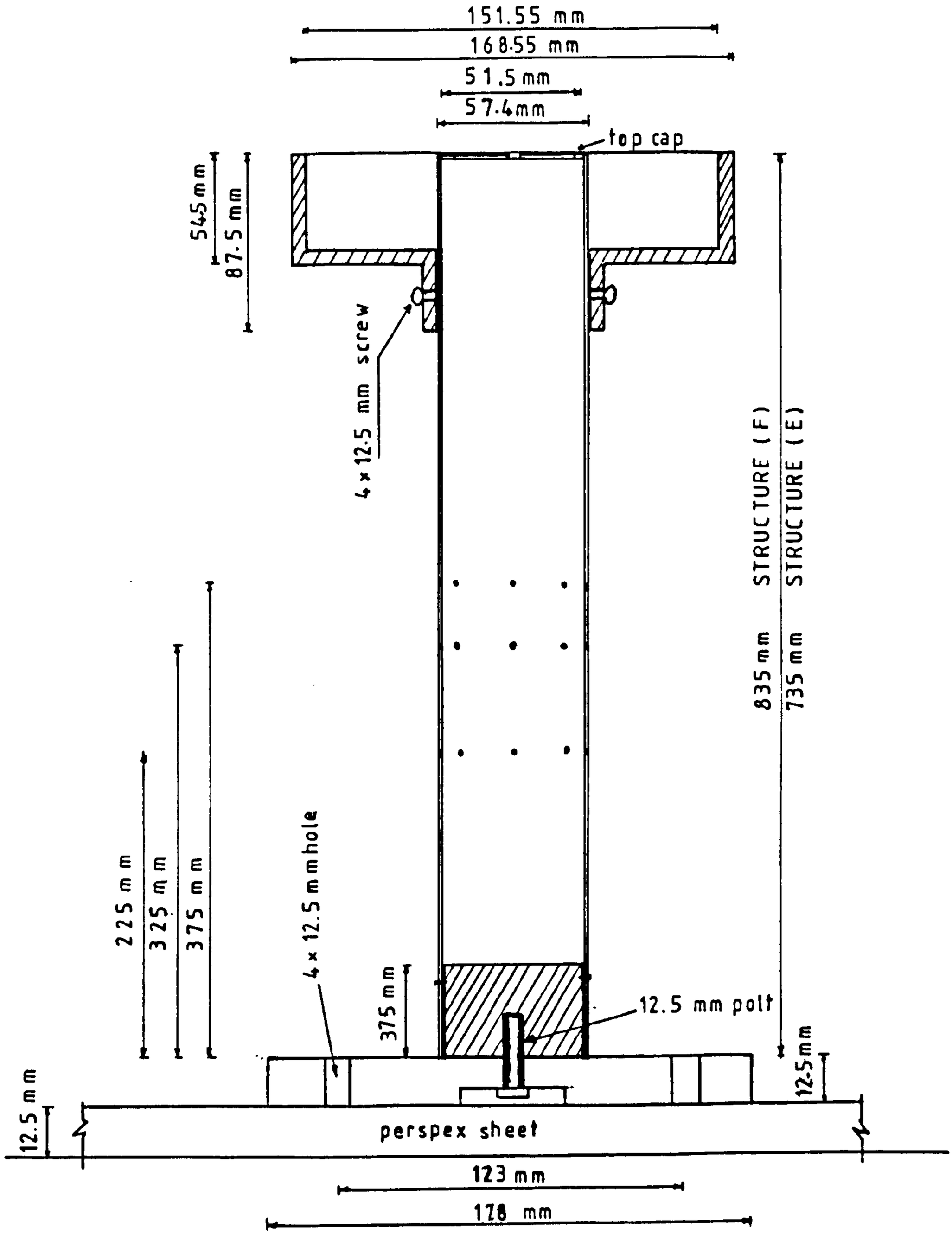


Figure (4.1.6) - Dimensions of Structures (E) and (F).

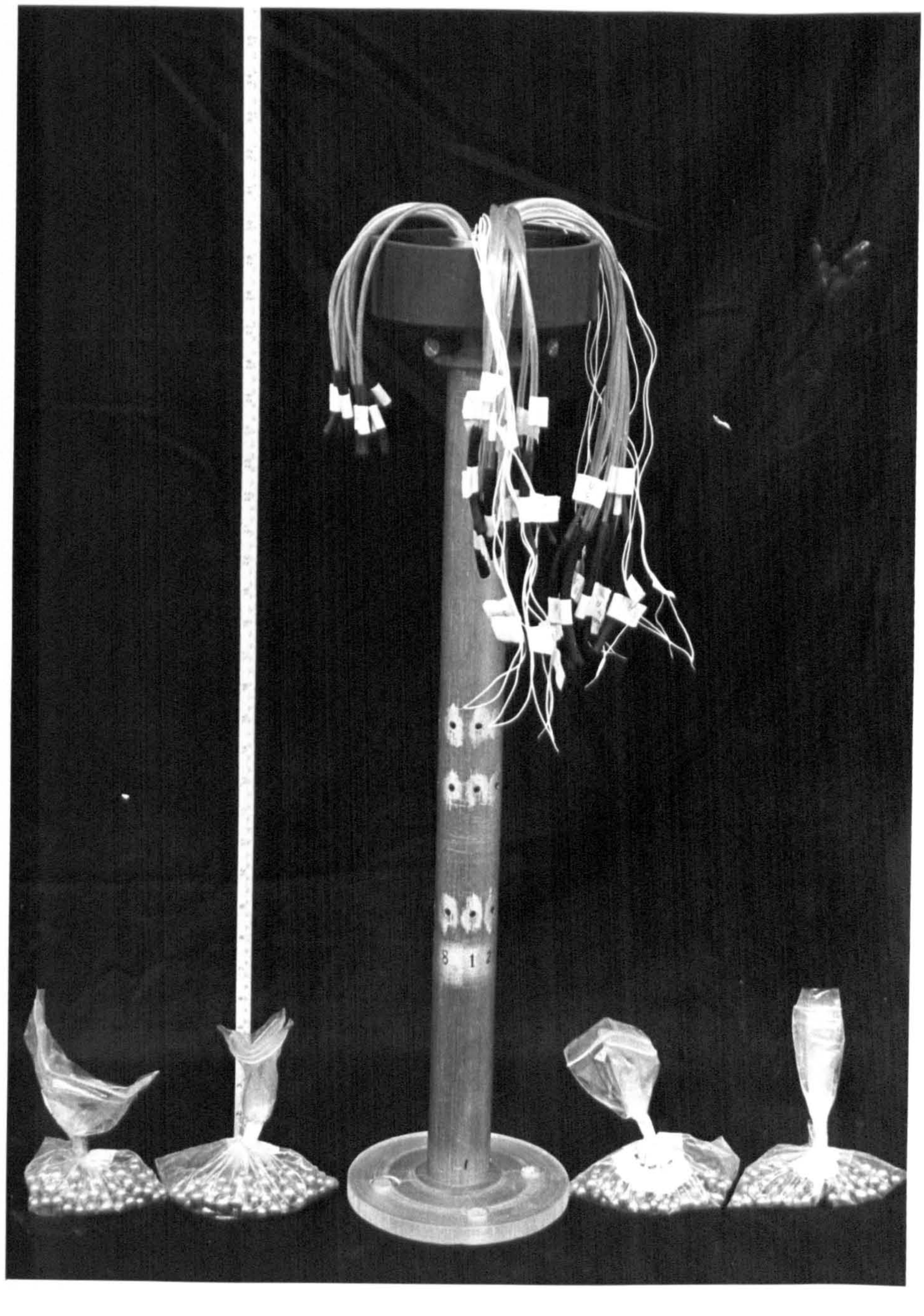


Figure (4.1.7) - Structure (E) and the Four Masses.

The mechanical properties of the structures tested are shown in Table (4.1).

Group	Structure designate	Material	Length meter	Diameter meter	Inside Diameter meter	Second Moment of area m^4	Modulus of Elasticity kN/m
One	A	Durapipe ABS	0.608	0.11	0.102	1.7048×10^{-6}	1509174.8
	B		0.808	0.11	0.102	1.7048×16^{-6}	1509174.8
Two	C	Durapipe ABS	0.735	0.0605	0.0517	0.2777×10^{-6}	1509174.8
	D		0.835	0.0605	0.0517	0.2777×10^{-6}	1509174.8
Three	E	Aluminium	0.735	0.0574	0.0515	0.0822×10^{-6}	58022298.0
	F		0.835	0.0574	0.0515	0.0822×10^{-6}	58022298.0

Table (4.1) - Mechanical Properties of the Structures.

4.2 DESCRIPTION OF INSTRUMENTATION

All the instrumentation used in this work was electrical, some of which was commercially available and the rest was designed and built especially for this work.

The instrumentation used may be divided into the measuring instrumentation and the reading instrumentation. The instrumentation used for measuring each parameter needed in the analysis and results was designated as the measuring instrumentation, while the reading instrumentation was connected to the measuring instrumentation to give the final measured data.

4.2.1 MEASUREMENT INSTRUMENTATION

The measuring instrumentation was designed and placed to give the most reliable measurement without seriously affecting the characteristics and the environment of the test.

4.2.1.1 WAVE PROBE

The wave probe was used to measure the wave profile. The measurement was done by detecting the change of voltage due to the change of resistance of the wire probe following the equation

$$\tilde{I} \times \tilde{R} = \tilde{V}$$

where $\tilde{I}$ = the constant alternative current

$\tilde{V}$ = the variable D.C Voltage

$\tilde{R}$ = the variable resistance of the probe wire due to the change of the wave profile

Alternating current was used at the tips of the probe to avoid the polarisation of waver.

Two stainless steel wires having a diameter of 0·002 metre and a length of 0·5 metre with a resistance of 0·006 ohms per metre were used as the probe. The two wires were stretched 0·025 metre apart and fixed to a perspex plate; the perspex plate was suspended over the tank by means of an adjustable boom. This mounting facilitated the daily calibration of the probe and also aided in maintaining the probe in the proper position relative to the still water level.

There are certain disadvantages in the use of the resistance type probe. The relationship between water surface position and the electrical signal output is non-linear. Furthermore, the calibration of the system is dependent upon the conductivity of the water between the two wires. These difficulties have been overcome satisfactorily by always adjusting the position of the probe to the still water level existing in the tank and also by adding a resistance to the electrical circuits before the electrical current reaches the probe.

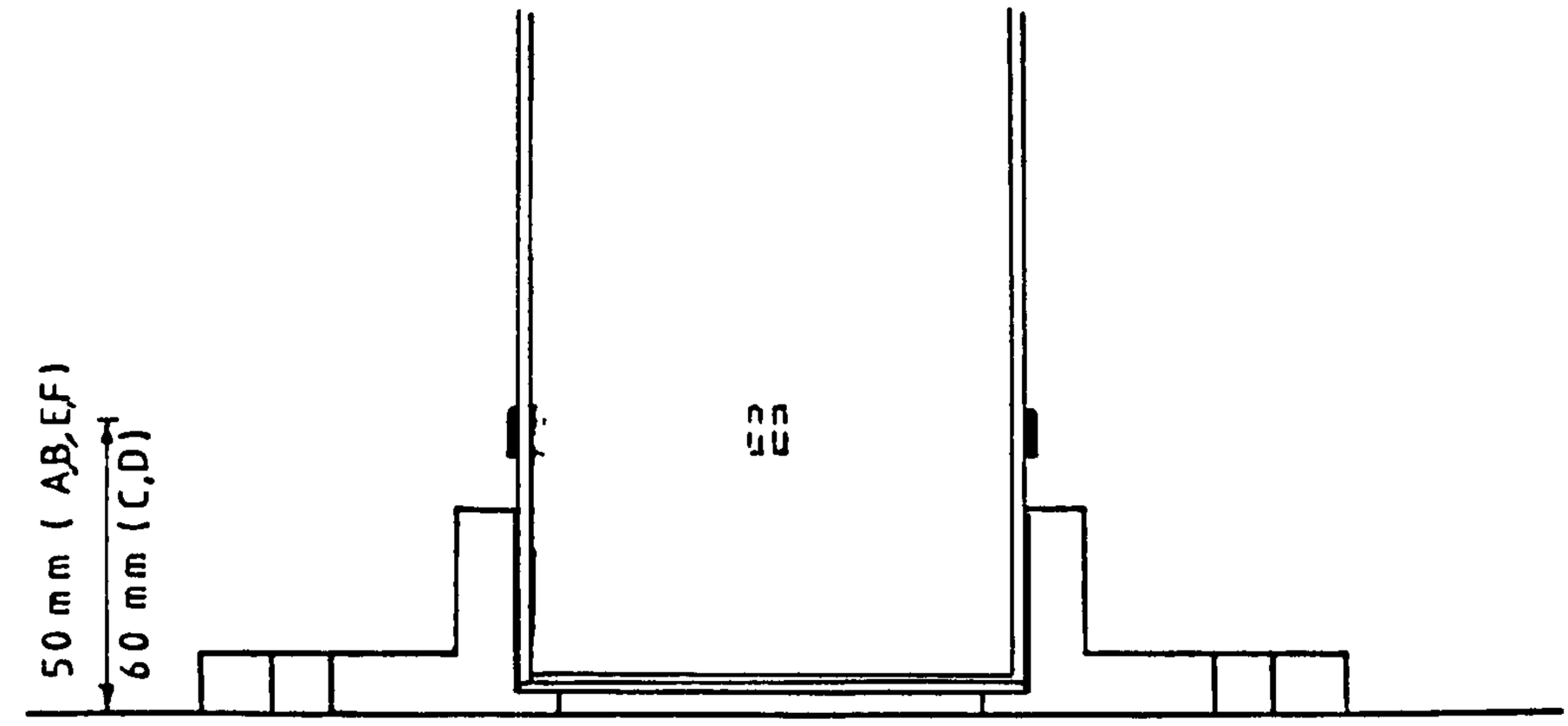
4.2.1.2 STRAIN GAUGES AT THE BASE OF THE STRUCTURES

The bending moment at the base of the structure was measured with strain gauges. They were Tokyo Sokki strain gauges, type FLA-6-11, the gauges lengths were 0·006 metre and resistance of 120 ohms. For each structure two sets of gauges were used with four gauges in each set. The gauges were placed and connected as shown in Figure (4.2.1), one set was fixed to measure the bending moment in-line with the wave direction and the other perpendicular to the wave direction.

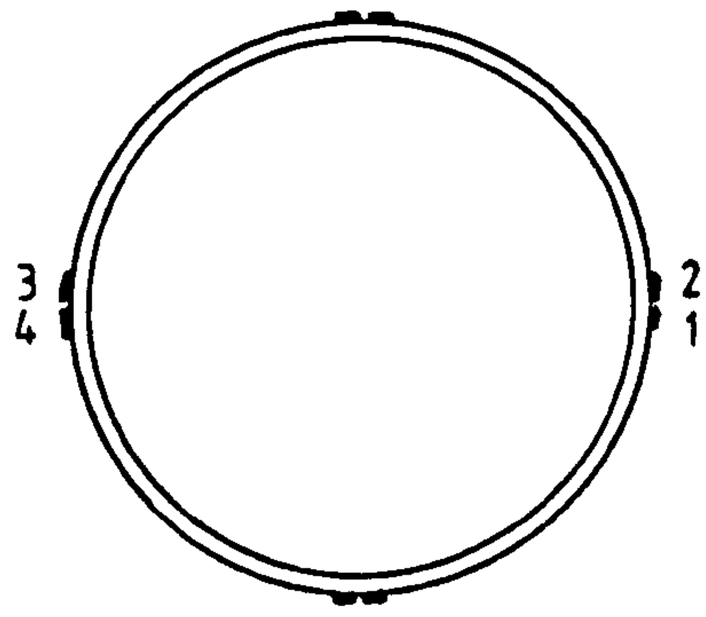
The strain gauges for the Durapipe (ABS) structures were placed outside of the cylinder and covered with watertight paint while the wires from the strain gauges were carried inside the cylinders. The reason for placing the strain gauges outside is to keep them at a constant temperature throughout the test (water temperature). For the aluminium structures the strain gauges were placed inside.

The strain gauges were connected so that only the bending moment was measured. The direct load, such as that form the concentrated mass, did not influence the strain gauge reading.

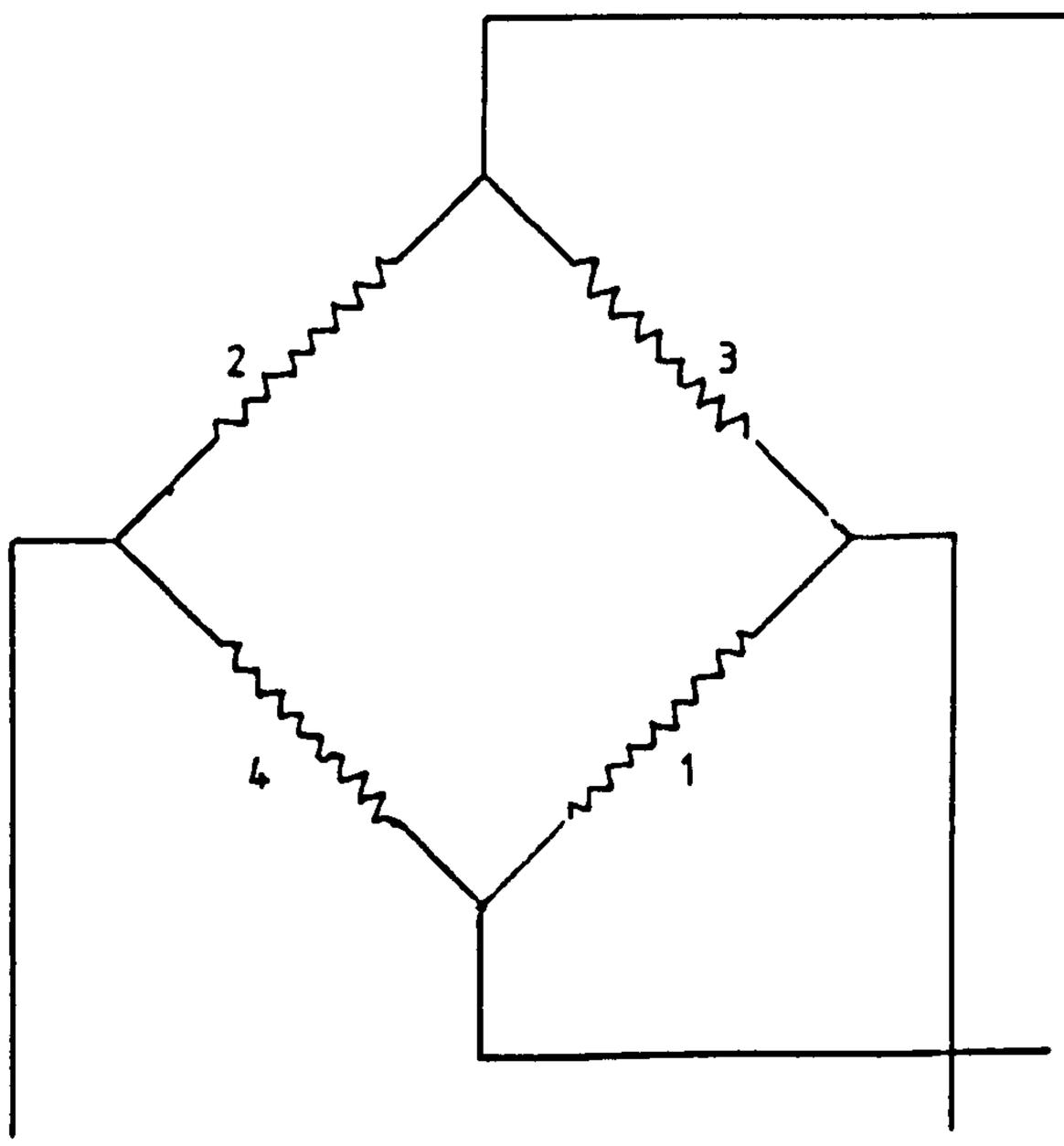
The strain gauges were aligned with respect to the tank by a vertical mark inscribed throughout the length of the cylinder and its base.



Sectional elevation of strain gauge installation



Plan view of strain gauge installation



Full Wheastone bridge of strain gauge

Figure (4.2.1) - Strain Gauge Installation.

4.2.1.3 PRESSURE TRANSDUCERS

A more accurate method for measuring the force of the waves at any level of the structure is one in which the pressure is measured at equal intervals around the circumference of the structure at that level by use of pressure transducers. The total pressure force on a unit height of the cylinder may then be obtained by integrating the pressure on the individual surface elements around the circumference.

The general description of design of the pressure transducers was desirable to incorporate the following features:

1. The pressure transducers should be capable of measuring pressures from 0·01 psi to 0·5 psi

2. The output of the transducers versus the applied pressure should be linear

3. The output traces should be clean "no noise", this could be obtained by the high natural frequency of the transducer sensing element. This feature reduces the error in reading the output

4. The pressure transducers should be easily attached and removed to read the pressure at any point along the pile

Because no known commercial transducer met the requirements listed above, a special transducer shown in Appendix (A) was developed.

During the development of the transducer, a number of different types of sensing element were considered, fabricated and tested. Based on the experience obtained during these tests, the final design of the sensing element was arrived at.

Eight pressure transducers were built so that eight individual pressure readings around the circumference of the structure at the level required can be obtained at the same time.

The pressure transducers were mounted on the top of the tank. To measure the pressure at any particular points on the structure several soft tubes were connected from the various points on the structure to the transducers.

Figure (4.2.2) shows the mounting and the connection of the pressure transducers to the structure.

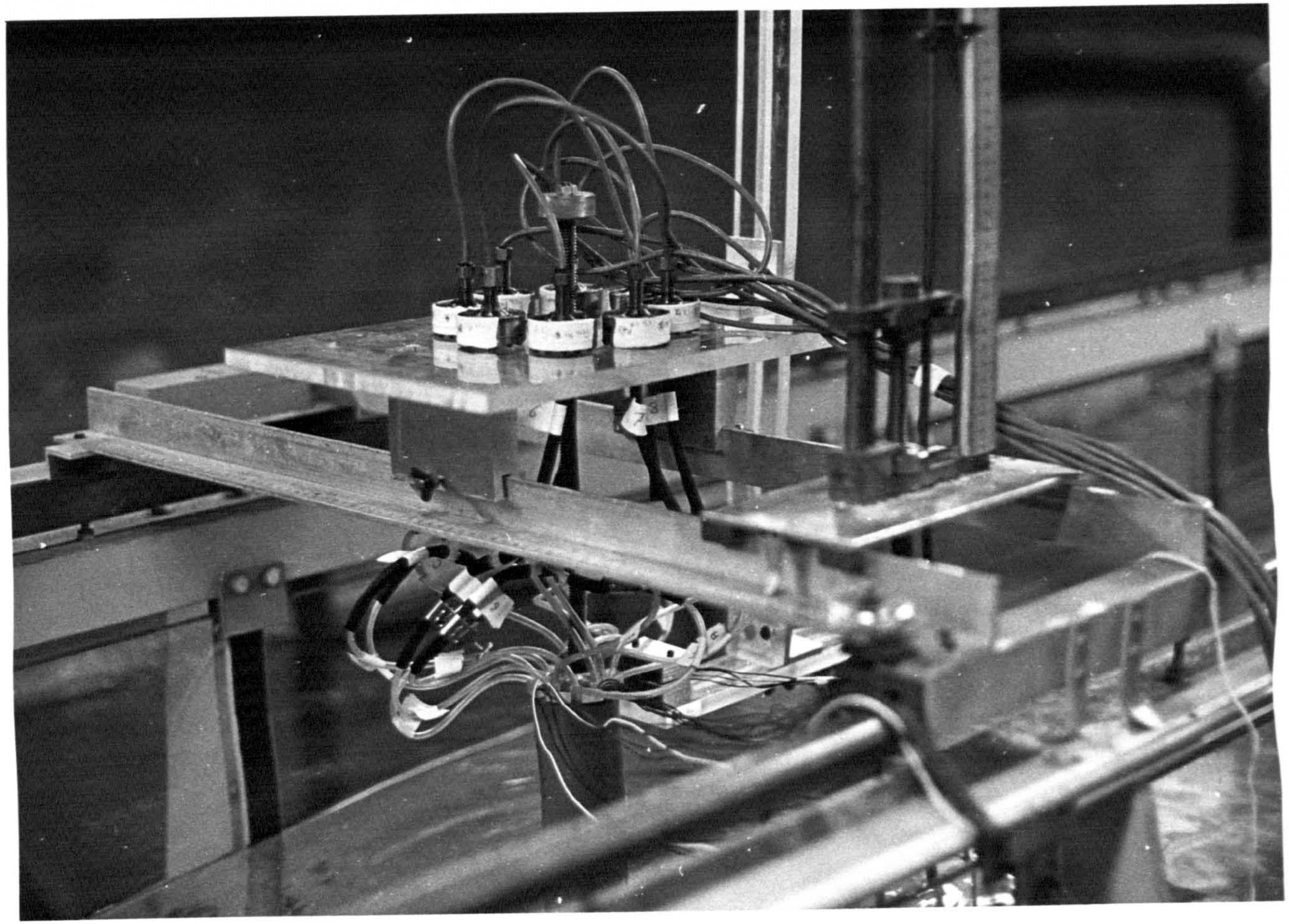


Figure (4.2.2) - The Mounting and Connection of the Pressure Transducer to the Structure.

4.2.1.4 DISPLACEMENT TRANSDUCER

An inductive type of displacement transducer was attached to the tip of the structure to measure the tip displacement of the structure caused by the wave action.

This displacement transducer utilises the well proven linear variable differential transformer principle. A transducer with a free unguided armature assembly was used so that the movement of the structure will not be affected by the stiffness of the spring loaded armature. This is especially so because the wave loading on the structure is small and consequently the movement of the structure was also small.

The working range of the transducer is ± 0·0125 metre.

4.2.2 READING INSTRUMENTS

The reading instruments were used in conjunction with the measuring instruments. The reading instruments were used to drive, amplify the magnitude of the electrical signal coming out from the sensing element and purify it from any noise interference in order to be able to record and digitalize it without any alteration of the actual electric signal measured.

Ten channels out of the fourteen strain gauge bridge amplifier units of the basic Sensonic amplifier were used to drive and pick up the electric signals from the two sets of the strain gauges which were fixed at the base of the structure, and also from the strain gauges of each of the pressure transducers.

Two channels of inductive amplifier units were used with the wave probe to measure the wave profile.

A RDP transducer amplifier type D7M was used to drive and amplify the electrical signal of the displacement transducer.

The low level signals "output" from the measurement instrumentation drive amplifier were converted into higher level signals and also filtered with low output impedence by using a preamplifier. The preamplifier was designed and manufactured to satisfy the necessary requirements of the signals level to be recorded. For the electric circuit of the preamplifier see Appendix (B).

The amplification and the phase shift due to filtration in the preamplifier were tested by using the output from a signal source type 471 which generates a sinusoidal wave-form of 100 MV at frequency ranges between 1 Hz to 10 Hz as an input to the preamplifier. Both the output from the preamplifier and the signal source were fed into a

two channel oscilloscope type D1011 which attenuated the signal from the amplifier so that both signals of similar amplitude were displayed on the screen, the attenuation inside the oscilloscope was equal to the gain of the preamplifier. The phase shift from the two signals was negligible.

All the outputs from the preamplifier channels were recorded simultaneously by one fourteen channel RACOL tape recorder store 14D5. The electric signals were recorded at a speed of $1\frac{7}{8}$ and at the window of recording $\pm$ 2 volt peak to peak.

Figure (4.2.3) shows the strain gauge Sensonic amplifier, preamplifier and the tape recorder.

The electric signals analogue recorded were fet to an Intercole Transmitter/Receiver model HS6202C which was connected to a PDP Compulog Intercole system in order to be digitalized at regular intervals to be saved on a floppy disc.

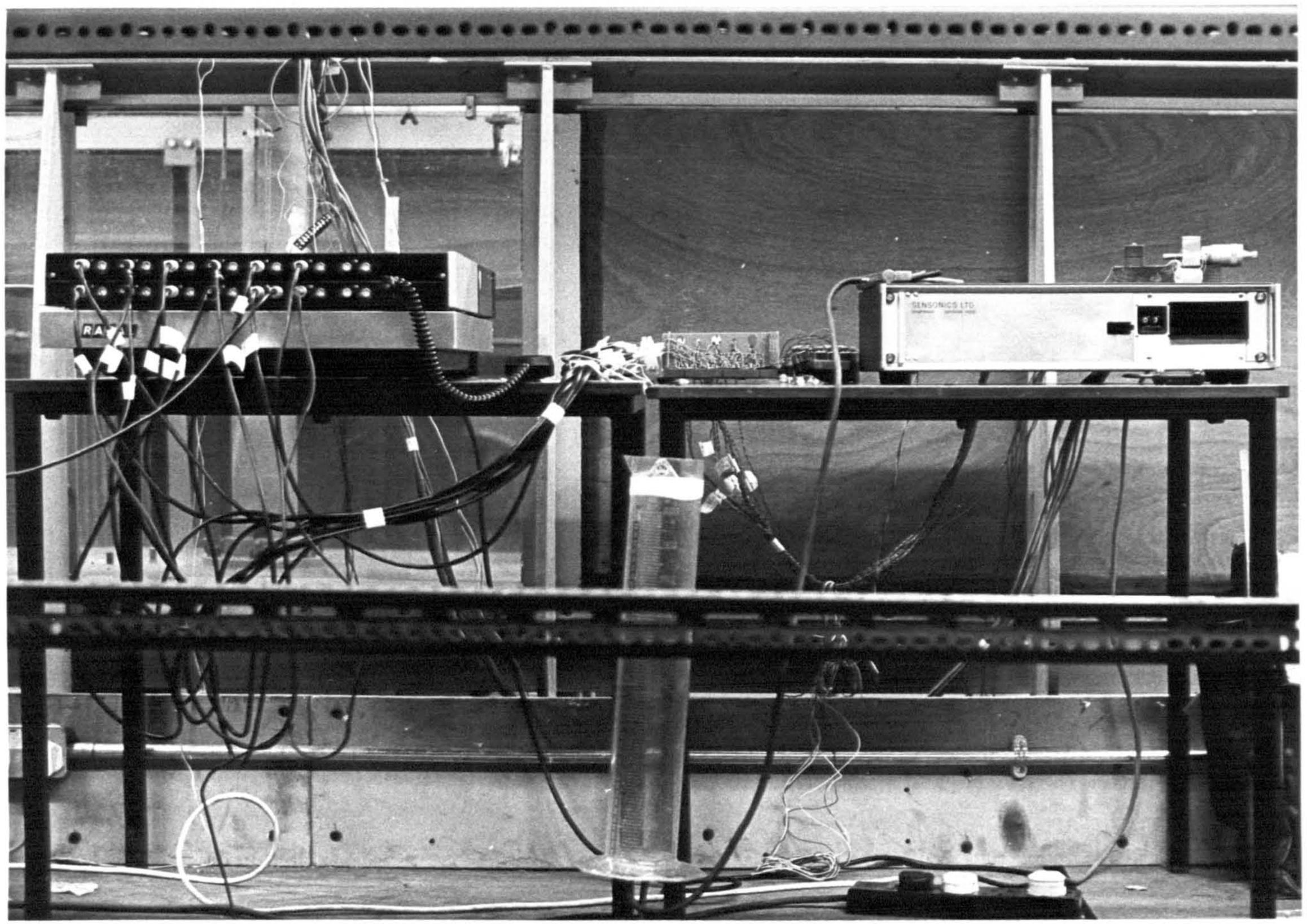


Figure (4.2.3) - Strain Gauge Amplifier, Preamplifier and the Tape Recorder.

4.3 CALIBRATION OF THE MEASURING INSTRUMENTATION

For the calibration tests of the measuring instrumentation the structure was first located and fixed in the tank. The displacement transducer was attached to the structure and the pressure transducers were mounted, so the structure was ready for the main test. The calibration tests were done twice each day; once before and once after the main tests.

All the electric equipment was switched on and left for half an hour before the first calibration tests start in order to achieve a stable electrical signal.

4.3.1 STRAIN GAUGES

The strain gauges at the base of the structure for measuring the bending moments were calibrated by applying a known horizontal load equal to 4·45 N at the top of the cylinder with a wire and pulley arrangement. This horizontal load had three increment of loading each equal to 4·45 N. For the in-line bending moment the tests were done once where the horizontal load was in the direction of the wave propagation and once in the opposite direction in order to obtain the positive and negative bending moments reading. Similar tests were carried out for the strain gauges in the transverse direction.

The bending moments at the base of the structure were calculated by multiplying the horizontal force by the distance from the top of the structure to the strain gauge. For the cases where added mass was applied at the top of the structure (third group of structures) the bending moment was calculated as the above plus the moment due to the concentrated mass at the top of the structure when the structure is in a deflection position. The deflection of the tip of the structure was determined in a previous test where a dial guage was attached to the top of the structure on theother side and on the same line with the wire so that the tip displacement due to the horizontal load were known. This test was useful also in the calibration of the displacement transducer as shown below.

Figure (4.3.1) shows the arrangement of the test determining the tip deflection.

4.3.2 DISPLACEMENT TRANSDUCER

The displacement transducer was calculated at the same time as the strain gauge. The horizontal displacements were known due to the horizontal applied loads. A mark on the transducer armature was made in order to correct the relative positions between the armature and the body of the transducer.

The relationship between the horizontal displacement and the electrical signal recorded was linear.

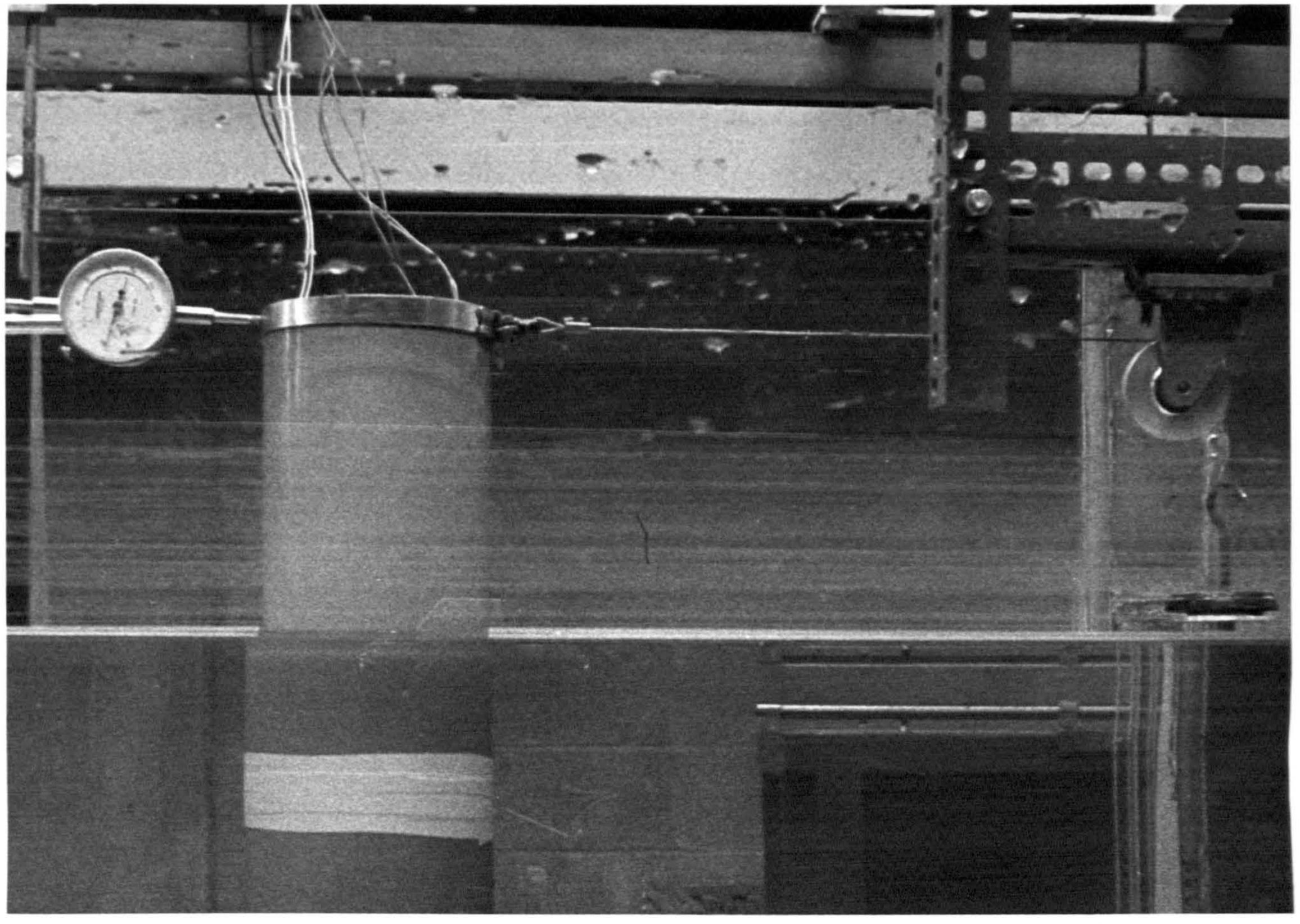


Figure (4.3.1) - The Test Determining the Tip Deflection due to the Horizontal Load.

4.3.3 PRESSURE TRANSDUCERS

The pressure transducers were calibrated for a positive gauge pressure by connecting the pressure transducers at the top of the tank with the eight holes around the circumference of the structure by the tube assembly in each structure. For that purpose the water level in the tank decreased to the level of these holes. This gave a zero gauge pressure reading. The electric signals from the pressure transducers were recorded. The water level in the tank increased by 0.025 metre which gave a gauge pressure equal to the head of water above the holes, the electric signals were recorded, the water level in the tank was again decreased to the level of these holes, again giving zero gauge pressure and the electric signals were recorded. By repeating this process with increments in the water level of 0.025 metre up to the maximum water level of 0.1 metre, the pressure transducers were calibrated for the positive pressure gauge readings.

The negative gauge pressures were obtained by disconnecting the pressure transducers from the tubes and connecting them again when the water level was 0.11 metre above the holes which give zero gauge pressure. The procedure of positive gauge pressure was reversed in order to calibrate the pressure transducers for the negative pressure gauge reading.

The calibration of the pressure transducers was done with the uppermost holes in the cases of structures with more than one holes level.

The following test was previously made in order to find the pressure losses along the length of the tube from the hole on the structure to the pressure transducer at the top of the water tank. The pressure transducer was connected to a single bore at the bottom of a graduated measuring cylinder with a single bore stopcock by 0.04 metre long rubber tube in the case of testing for 0.025 metre head of water. This was in order to keep the pressure transducer away from the water as it is not waterproof. The electric signals were recorded, water was poured down to the level 0.025 metre in the measuring cylinder then electric signal was recorded. The water was emptied from the measuring cylinder by the stop-cock. A 0.75 metre long rubber tube of the same length as those used in the tests, replaced the 0.04 metre tube and the test was repeated.

The above test was repeated with different head of water pressure and different lengths of the short tube. The difference between the electric signals recorded from the short tubes and from the long tube were negligible.

4.3.4 WAVE PROBE

After the position of the wires of the probe had been adjusted to the still water level the electric signal gave a zero reading. The probe was lowered by means of an adjustable boom by O·Ol metre which was equivalent to raising the water level by the same amplitude and the electric signal was recorded. This test was carried out by lowering the probe in stages of O·Ol metre until a distance of O·l metre was reached.

Similar tests were carried out by raising instead of lowering the wave probe to give the calibration in the case of lowering of the water level.

4.4 EXPERIMENT

To achieve the main objective of the investigation, it was decided to carry out experiments in order to determine all the parameters required in the analysis, which are the water wave criteria and the structure dynamic properties.

4.4.1 WATER WAVE CRITERIA

The first test was carried out to see how close small amplitude harmonic waves conformed to sinusoidal shape, and how their characteristics varied with time.

The test was done by adjusting the wedge arm and the speed of the driving motor of the wave generator. The wave profile was measured by the wave probe which was located at the position of the tested structure. The test was carried out for 45 minutes.

For conformation of sinusoidal shape a typical case is shown in Figure (4.4.1).

The variation of the wave characteristic during the test was insignificant.

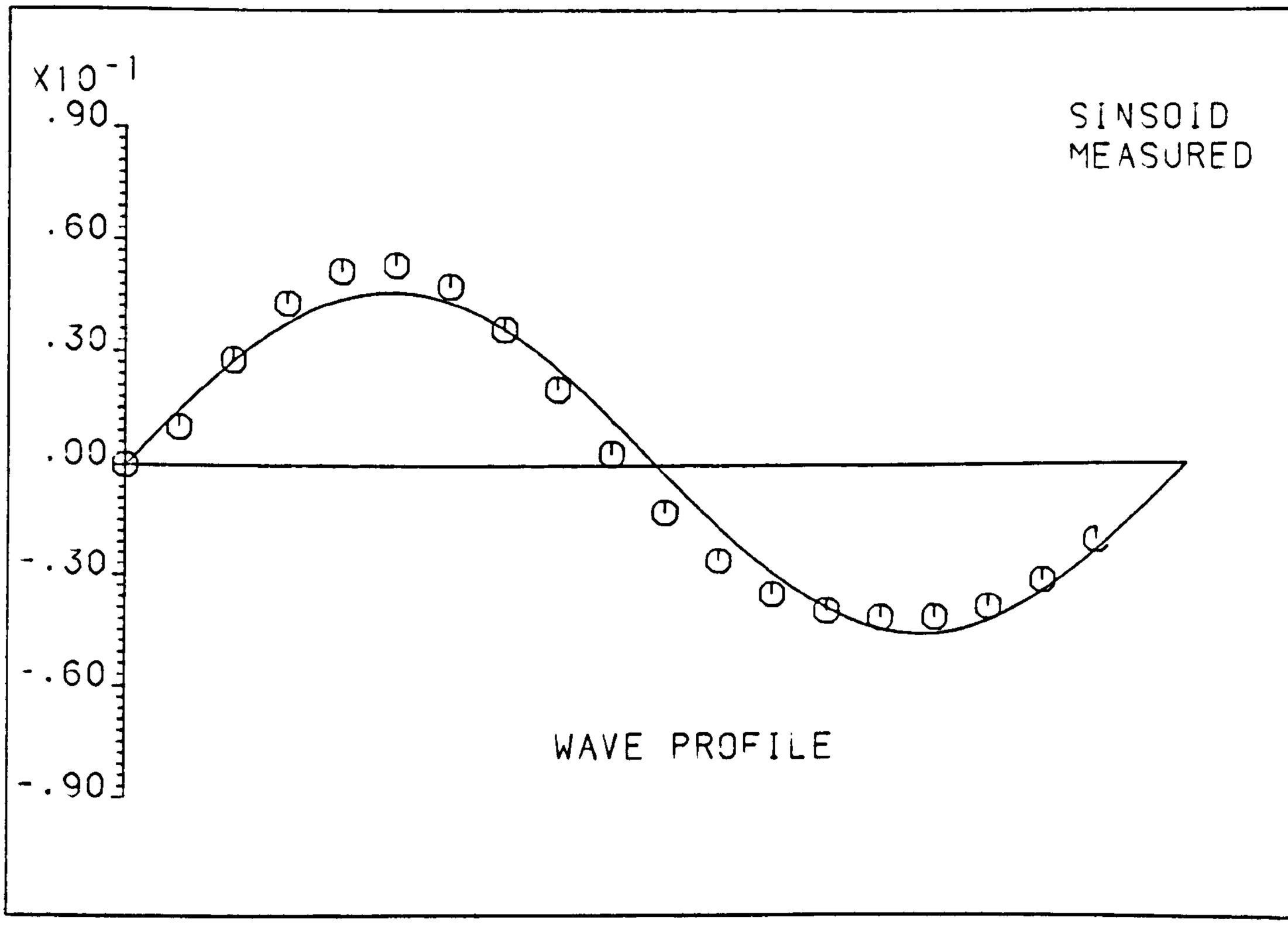


Figure (4.4.1) - Water Surface Profile.

It was necessary to check the degree of wave reflection from the beach. The amount proved to be quite small.

Eagleson and Dean 1966[110] show that a portion of a wave impinging on a sloping surface will reflect upon itself. The magnitude of the reflected wave can be expressed as a coefficient of reflection times the magnitude of the incident wave. The reflected wave will disturb the regularity of the wave surface. The crests and troughs of the resulting wave will describe an envelope as shown in Figure (4.4.2).

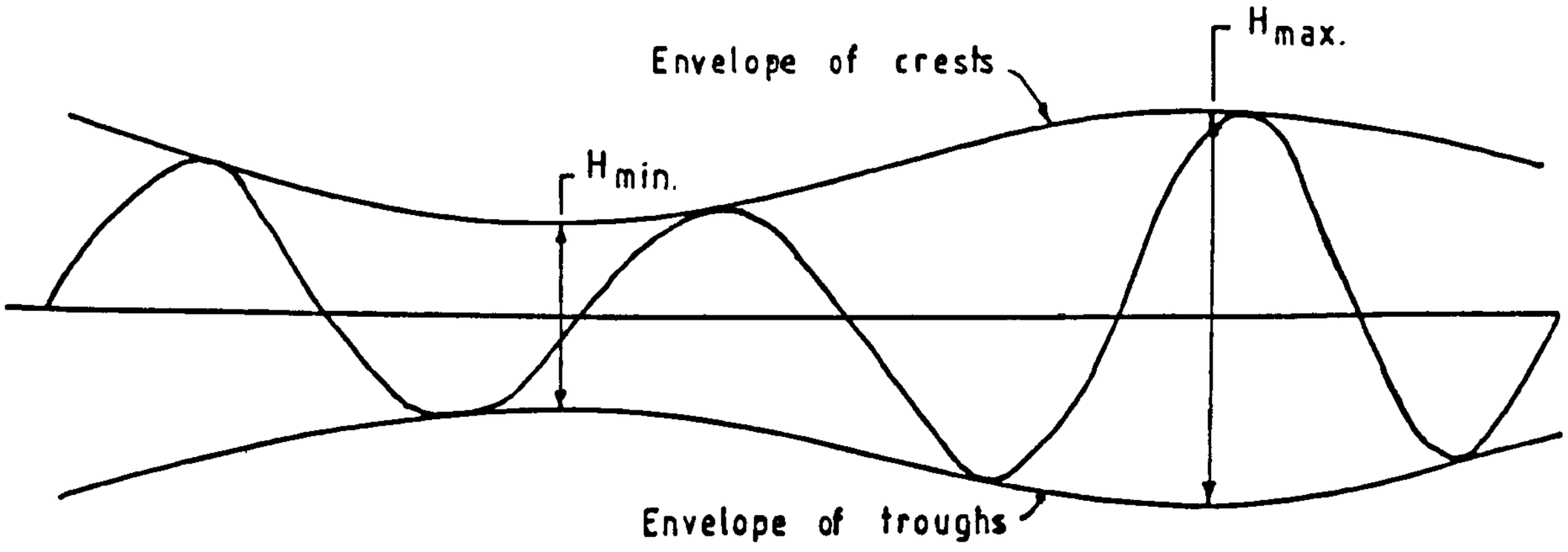


Figure (4.4.2) - Typical Reflected Wave.

It is shown in the above reference that the reflection coefficient can be obtained from:

$$K_r = \frac{H_{max} - H_{min}}{H_{max} + H_{min}} \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (4.4.1)$$

where K_r = Reflection coefficient

H_{max} = Maximum wave height

H_{min} = Minimum wave height

Waves of various characteristics were tested. From the results it can be concluded that reflection was minimal, and thus it was ignored in all analysis. Figure (4.4.3) shows the result of the reflection coefficient verse wave period.

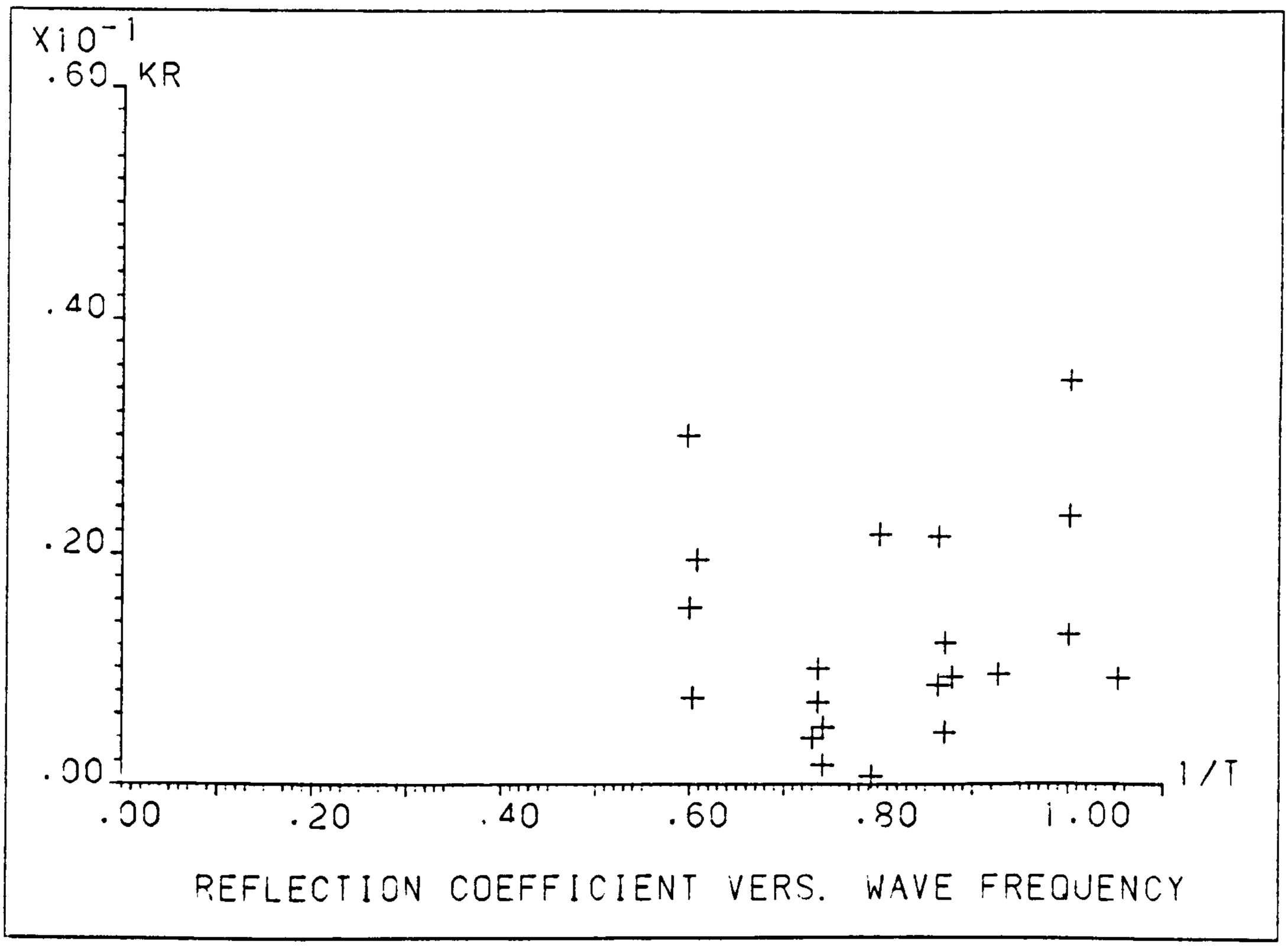


Figure (4.4.3) - Reflection Coefficient for the Wave Tank.

4.4.2 STRUCTURE DYNAMIC PROPERTIES

As the determination of certain structural parameter such as stiffness, natural frequencies and damping coefficients are important to be obtained, the following experiments were carried out.

4.4.2.1 STIFFNESS TEST

In the determination of the stiffness of the structure it is vital to know the rate of the deformation of the structure material, especially the Durapipe ABS structure. The Durapipe ABS is not a truly elastic material. Maxwell and Harrington 1952[111] show that the strain of this material depends on the time rate of deformation. That is, if the velocity of deformation is high enough the stress-strain curve will be linear. When the velocity of deformation becomes small or negligible, then a plastic flow takes place which is recoverable when the load is removed. There is a limit of strain however, beyond which plastic deformation is not recoverable. This can be called deformation types I, II and III and can be illustrated schematically as in Figure (4.4.4).

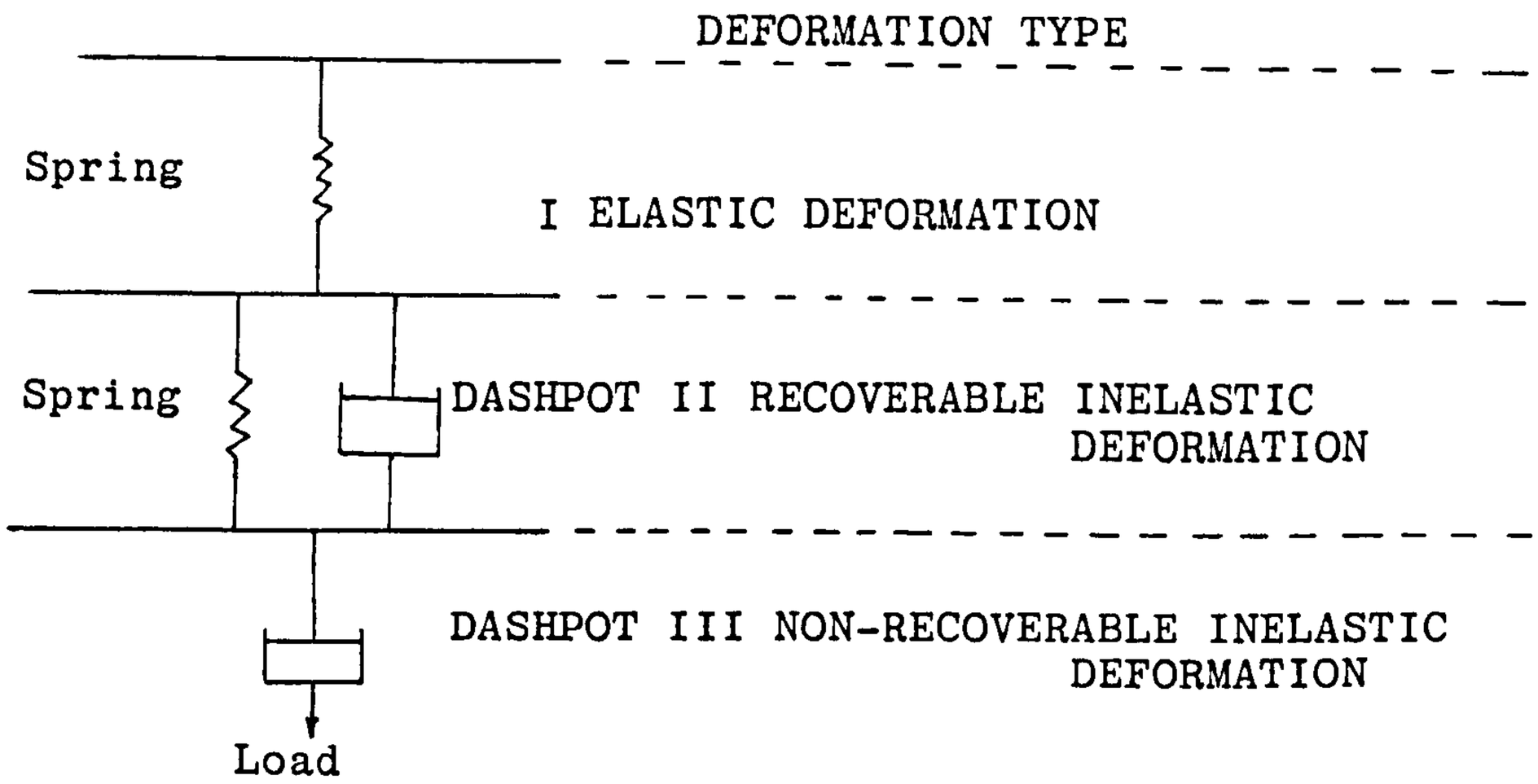


Figure (4.4.4) - Schematic Diagram of Deformation Types for Durapipe ABS.

For this work it will be shown that type I deformation occurred for the Durapipe ABS cylinders under dynamic loads.

The experimental procedure of this test was exactly as that of the calibration of the strain gauges at the base of the cylinder structure except that the output signal was taken to the Oscillograph instead of the tape recorder.

Typical records of the test procedure are shown in Figure (4.4.5), it was noted that the strain started at zero and ended at zero, indicating that type III deformation did not take place. The load was applied fairly suddenly so that a few oscillations occured in the record. It can be seen that strain did not decrease

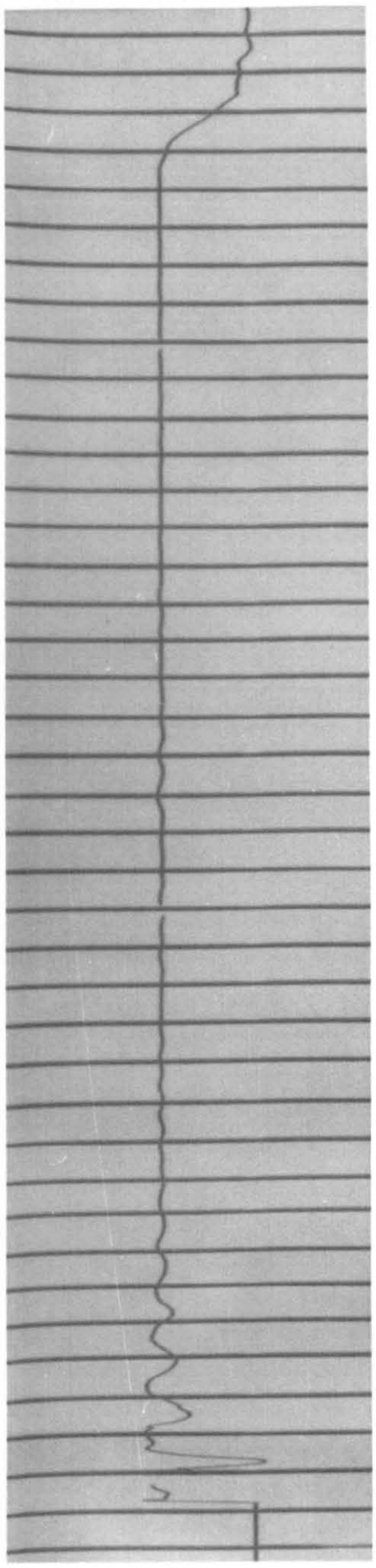


Figure (4.4.5) - Typical Recording of the Strain Gauge.

with time for constant load. If decrease showed in the strain under static load condition type II deformation would occur.

The stiffness of the cylinders was determined by the following test. The cylinder was mounted vertically as in the case of the main test and a series of weights were applied at the free end. The free end deflection corresponding to each weight was measured using a dial gauge. Graphs of load against the deflection were plotted, the slope yielding the equivalent stiffness.

4.4.2.2 NATURAL FREQUENCIES AND STRUCTURAL DAMPING

For convenience, measurements of the natural frequency and damping were made simultaneously. Essentially, the scheme was to force the oscillation of the structure by means of an impulse load and then to take the reading of the strain gauge at the base of the structure. The signals were recorded on an oscillograph from which the structural damping and natural frequency in the fundamental mode could be measured.

Care was taken to make the amplitude of vibration during an impulse load test about the same as the amplitude of vibration during the experiments with waves.

In the second approach, the structures were deflected with a steady load and then released.

The natural damped frequencies of the structures were obtained from the same test records as those from which the damping coefficient were obtained.

Typical records of the test results are shown in Figure (4.4.6).

The structures were tested in air and water.

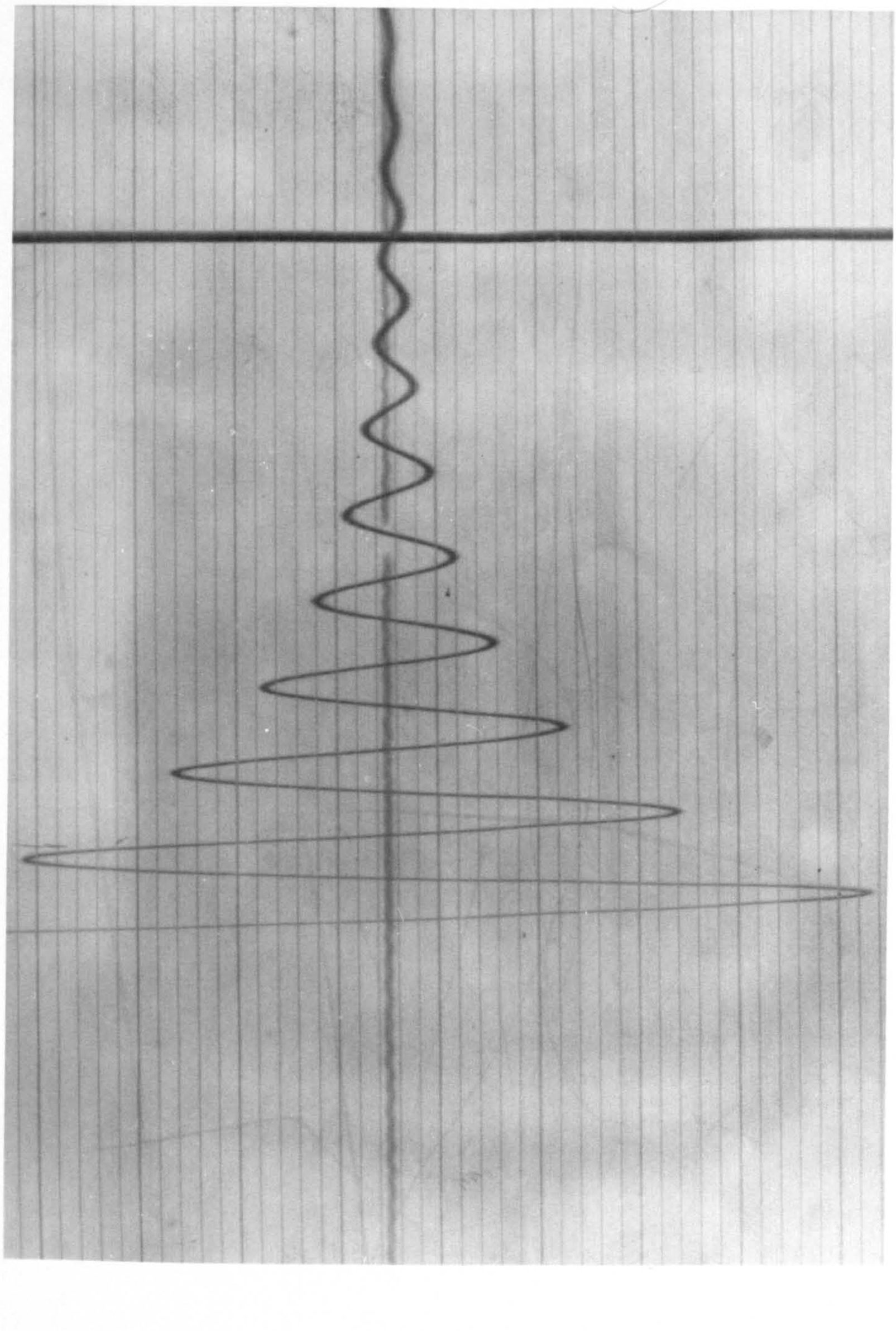


Figure (4.4.6) - Typical Record of the Natural Frequencies and Structural Damping.

4.4.3 MAIN TESTS

After the structure was installed in the water tank and all the measurement instrumentation was calibrated the tests were started by adjusting the water level in the tank. The wave generator arm was fixed at the position required and then switched on; the speed of the motor which controlled the frequency of the wave was set up by the digital counter. The records were taken after a few waves passed the structure in order to establish a steady state wave and structure motion. This was approximately five minutes after the operation of the wave generator.

First the structure was acting as a cantilever "free to vibrate" the signals from the eight pressure transducers connected to one level of the structure, from the wave probe located beside the structure, from the strain gauges at the base of the structure and from the displacement transducer at the top of the structure were recorded simultaneously for 15 minutes. Then the structure was clamped from the top "Fixed Structure" by screwing a stiff shaft fixed on the top of the water tank to the top cap of the structure. The effectiveness of the clamp at the top was tested by monitoring the signal of the displacement transducer. If the signal was zero throughout the wave cycle then the structure was fully clamped. After the structure was fully fixed the signals from the measurement instrumentation were recorded for 15 minutes. The stiff

shaft was released, returning the structure to the initial condition and the signals were recorded for 15 minutes.

This process was repeated for waves of other characteristics.

For the first group of the Durapipe ABS, structures A and B which had only one level in which the pressure could be measured; for this case the above test procedure was complete for testing this group.

Figure (4.4.7) and Figure (4.4.8) show structures A and B when they were under test, respectively.

The second group of Durapipe ABS, structures C and D which had two levels at which the pressures could be measured. For that reason after the above test procedures were carried out then the second level of pressure was recorded by locking the connection between the pressure transducers and first level and by opening the connection to the second level and repeating the procedure.

Figure (4.4.9) shows structure C when it was under test and also it shows clearly the position of the wave probe.

Figure (4.4.10) shows structure D under test and it also shows the way of the mounting of the displacement transducer and the stiff shaft used for clamping the top of the structure.

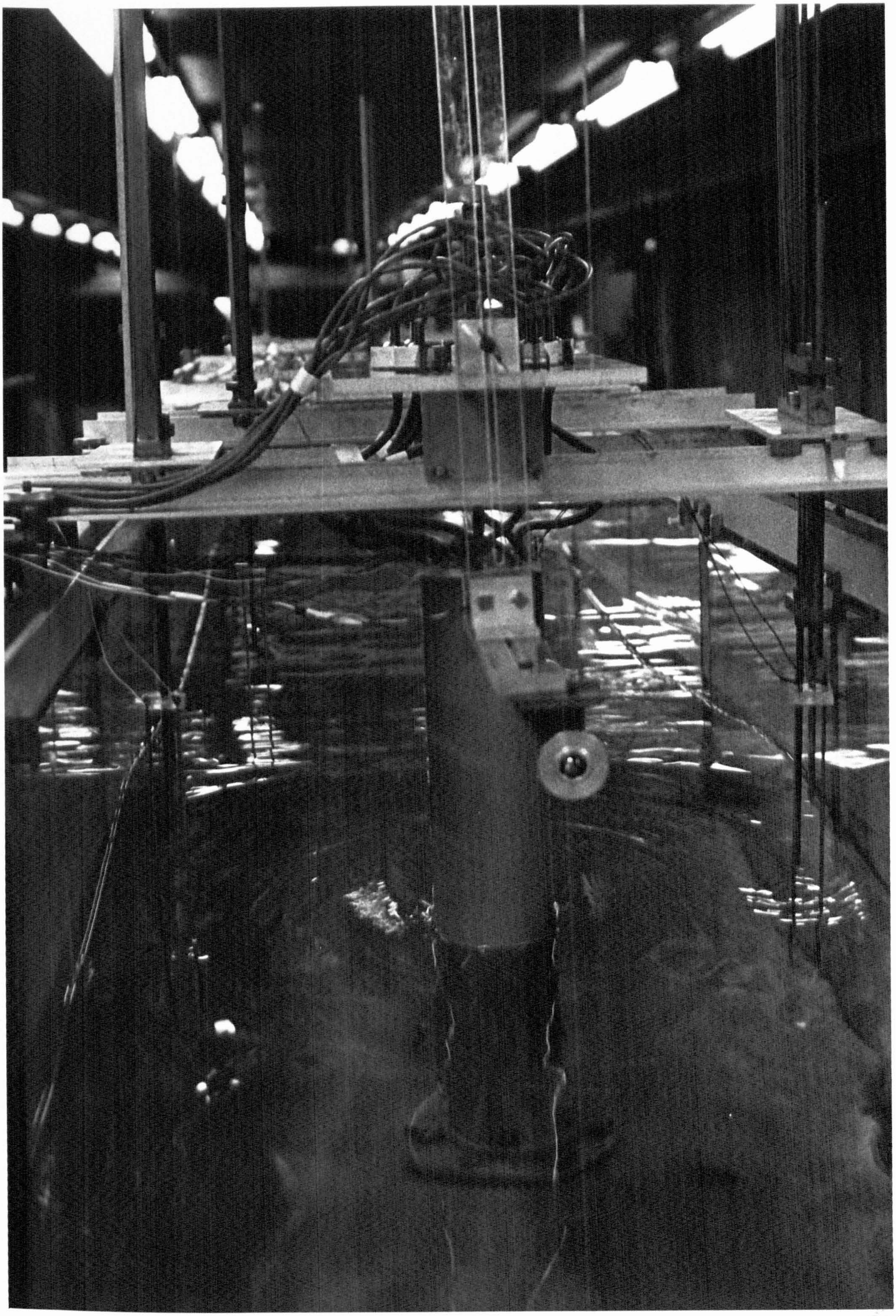


Figure (4.4.7) - Structure (A) under Test.

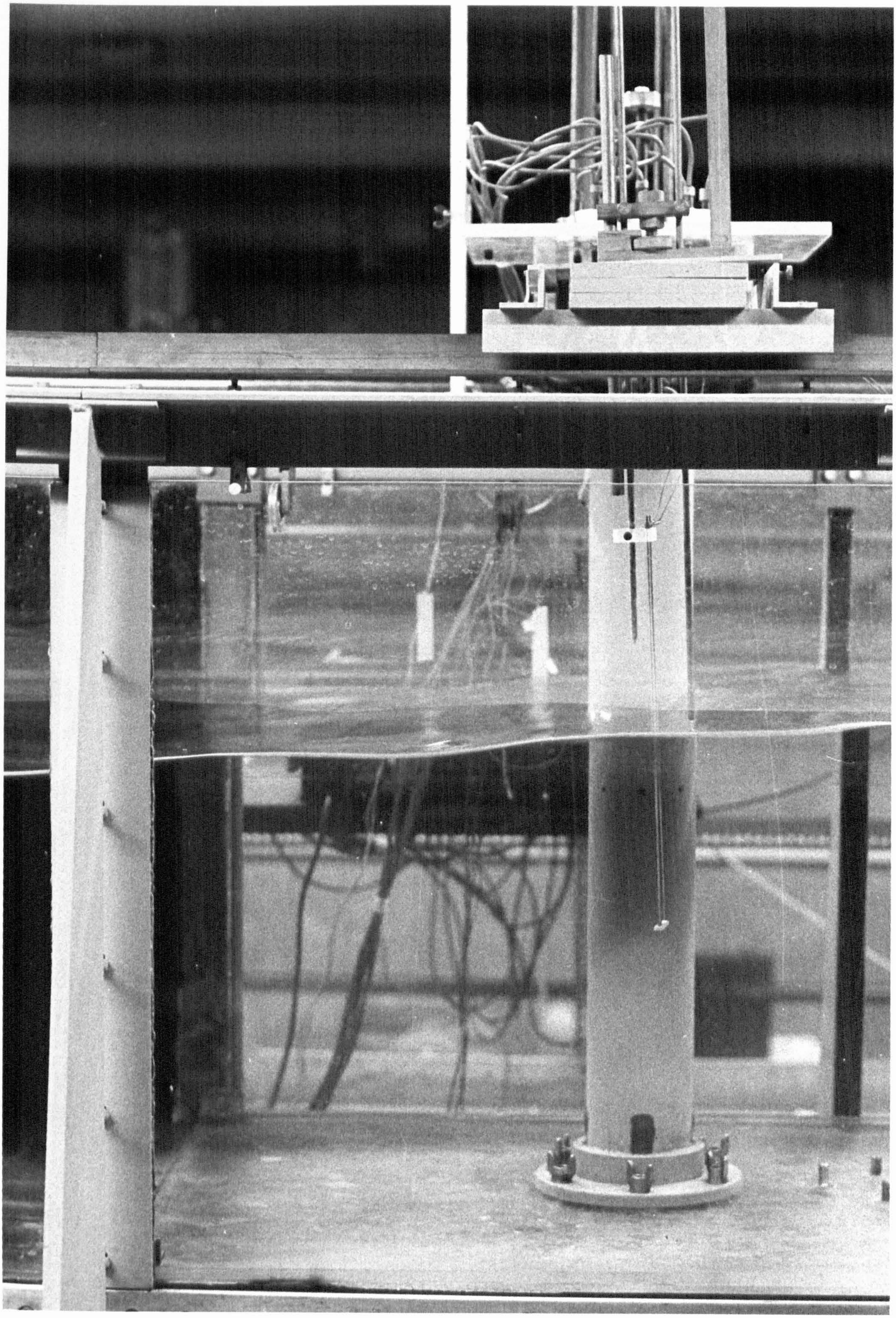


Figure (4.4.3) - Structure (B) under Test.

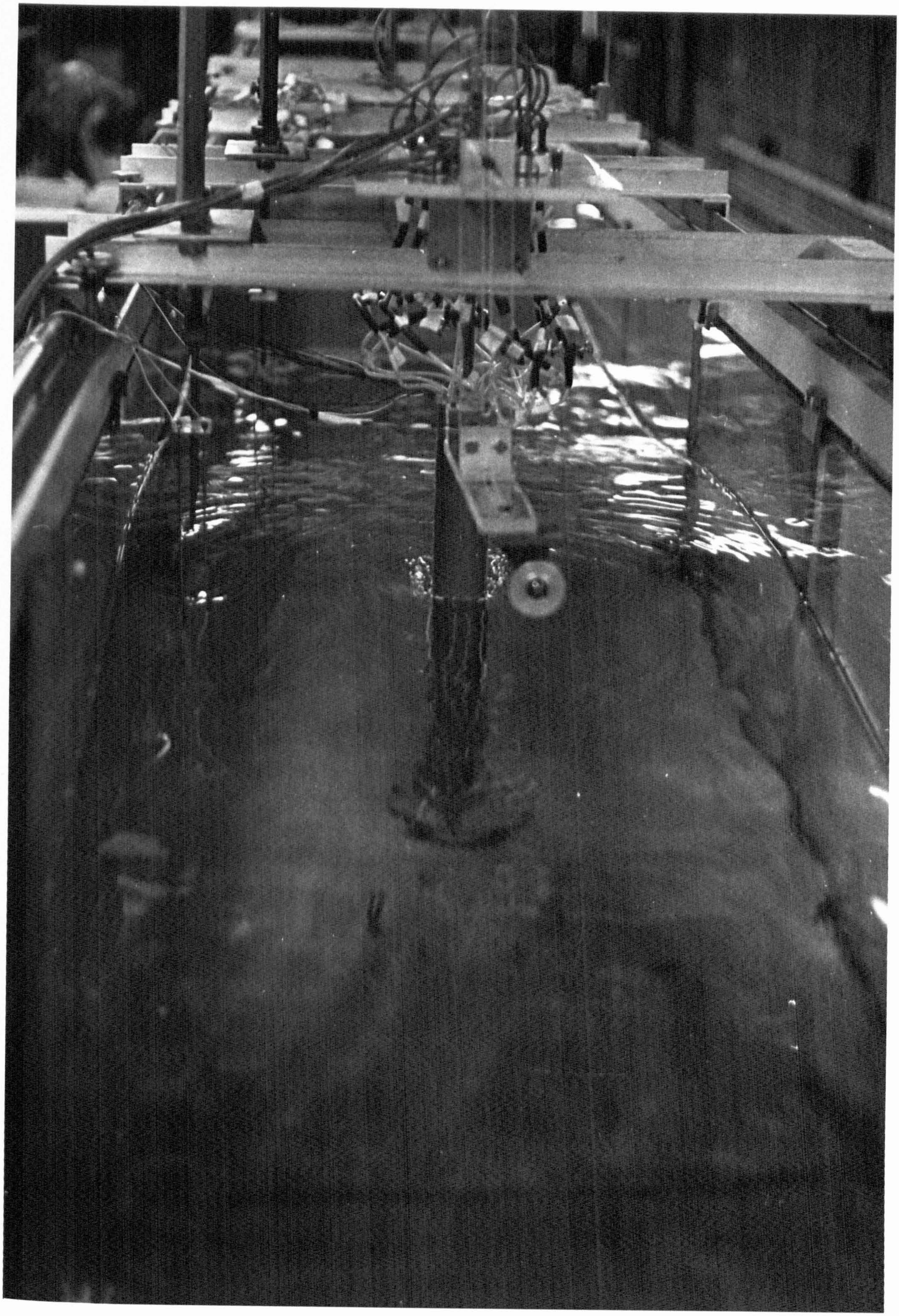


Figure (4.4.9) - Structure (C) under Test.

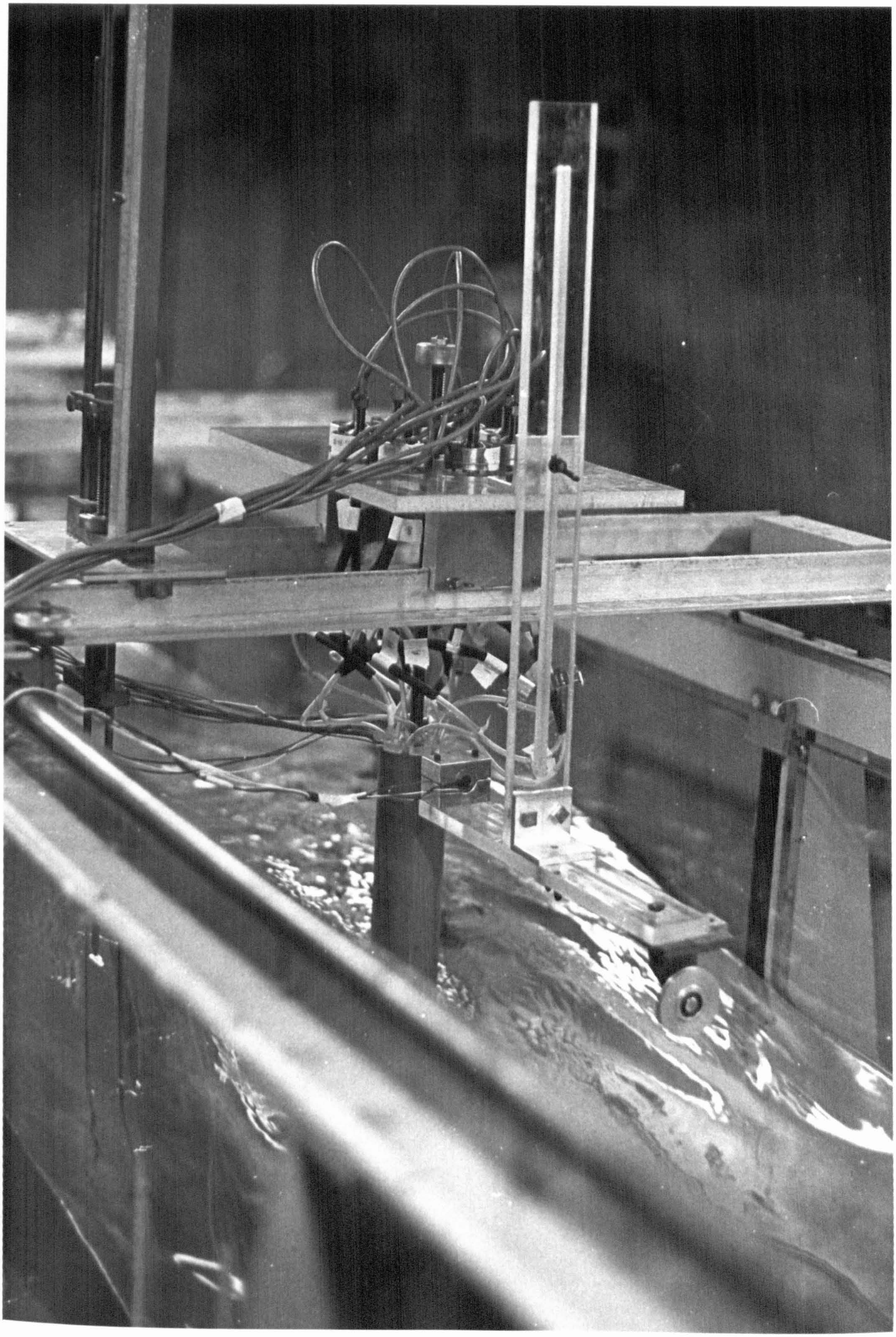


Figure (4.4.10) - Structure (D) under Test.

Figure (4.4.11) - Structure (E) under Test.

For the group of the Aluminium structures, structures E and F the test was carried out without any added mass for the highest level where the pressure can be recorded. The test was repeated for each increment of added mass for the same level of the pressure. For the other two levels which the pressure could be measured the same test procedure as used for the first level was repeated.

Figures (4.4.11) and (4.4.12) show the structure (E) and structure (F) while they were under test.

Table (4.2) shows the reading of the digitize electric signals of one of the calibration test, column one to eight are the eight pressure transducers, column nine is the strain gauge reading for the in-line bending moment while column ten is the reading of the transverse bending moment, column eleven is the reading of the displcaement transducer and column twelve is the reading of the wave probe.

Table (4.3) shows the digitize electric signals reading for one of the test.

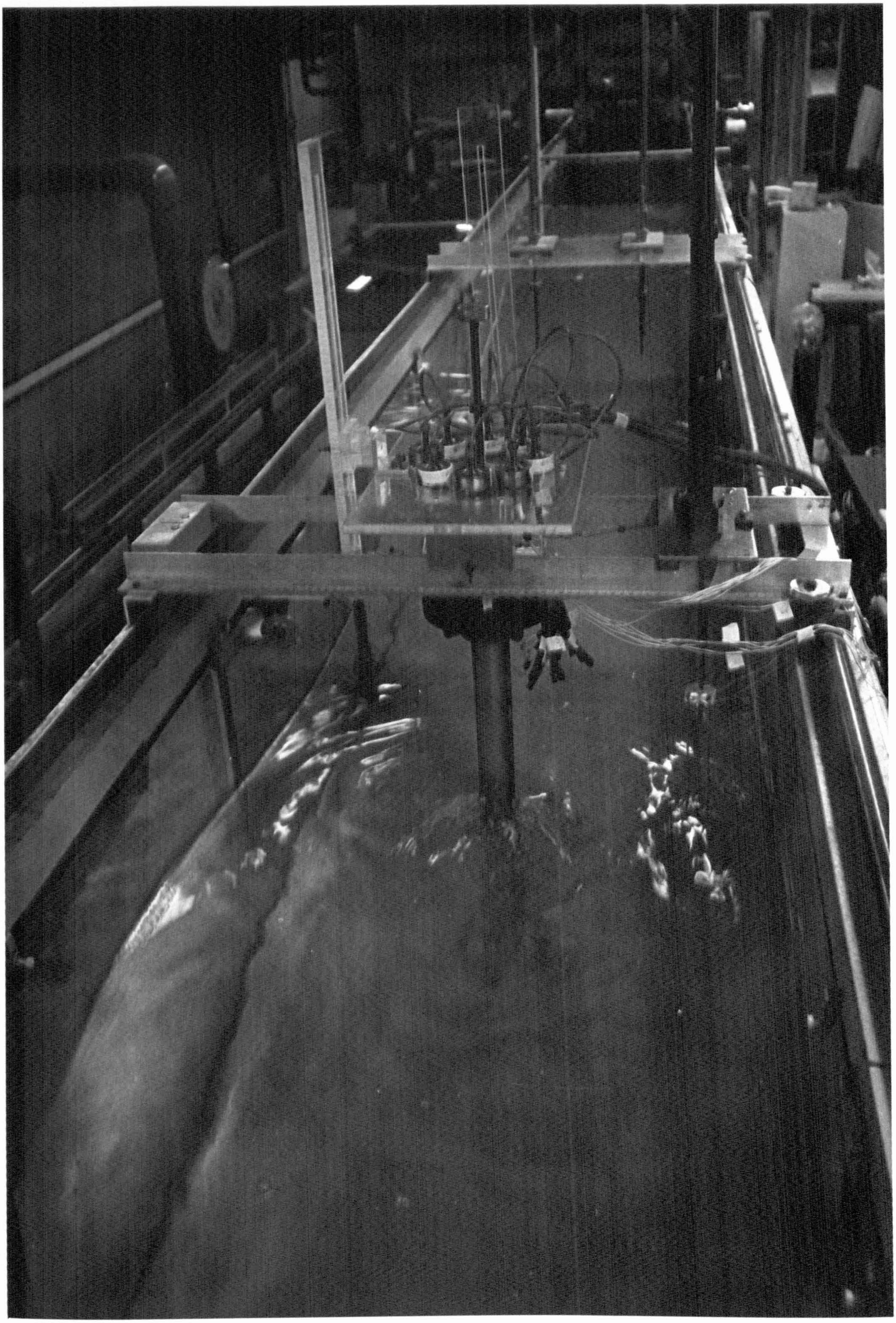


Figure (4.4.12) - Structure (F) under Test.

CALIBRATION FOR TEST NO=50 (C)

12	20	150									
-222	27	197	-174	-121	242	-177	3	-95	42	-45	-54
-220	29	194	-172	-119	240	-174	1	-97	40	-43	-51
-219	27	197	-170	-117	240	-177	-3	-97	36	-30	-62
-222	24	190	-165	-119	246	-169	3	-93	42	-35	-57
-237	26	186	-171	-128	234	-189	-38	-73	44	-105	-50
-240	27	188	-179	-119	223	-177	-53	-44	36	-45	-59
-236	32	182	-170	-117	226	-186	-59	-26	41	-36	-59
-231	31	182	-163	-121	231	-185	-50	-29	40	-35	-56
-229	25	192	-179	-112	230	-188	-56	-32	34	-84	-64
-237	24	184	-180	-112	219	-195	-53	-45	46	-86	-55
-234	29	174	-169	-116	225	-196	-48	-39	46	-94	-49
-710	-1340	555	-497	-407	663	-737	811	118	324	389	-202
-712	-1358	551	-483	-408	661	-741	824	122	323	381	-205
-708	-1367	554	-477	-405	664	-739	816	139	310	397	-211
-711	-1388	561	-487	-394	650	-731	806	136	311	393	-216
-712	-1392	557	-487	-393	651	-732	792	115	317	402	-216
-719	-1410	547	-489	-414	651	-750	820	123	327	392	-204
-712	-1411	562	-485	-395	652	-729	812	141	315	410	-215
-711	-1415	556	-479	-401	662	-737	828	132	326	400	-206
-709	-1421	557	-483	-391	654	-732	863	136	312	412	-214
-711	-1425	557	-475	-404	661	-739	828	126	328	402	-201
-531	-972	424	-352	-282	503	-565	501	-2	38	-76	-58
-537	-880	428	-363	-281	496	-561	475	5	32	-71	-62
-536	-967	418	-357	-280	493	-557	472	10	34	-76	-51
-536	-861	416	-352	-282	494	-562	480	9	35	-68	-53
-533	-967	421	-362	-281	491	-559	480	1	34	-68	-65
-533	-964	424	-357	-278	495	-561	484	6	45	-78	-52
-531	-969	425	-359	-276	490	-556	470	10	32	-61	-64
-525	-866	428	-357	-276	502	-566	477	-2	40	-69	-59
-530	-962	424	-349	-280	501	-565	470	-3	45	-70	-54
-530	-964	421	-353	-279	497	-566	487	-1	39	-69	-59
-534	-864	418	-348	-283	498	-565	453	6	763	942	-360
-360	-1025	628	-522	-416	721	-854	870	430	758	943	-363
-367	-1045	617	-518	-422	717	-863	872	423	765	917	-355
-360	-1064	625	-516	-416	713	-851	871	429	758	948	-363
-365	-1086	620	-514	-425	716	-863	874	417	760	931	-357
-365	-1100	626	-524	-413	708	-854	862	453	756	929	-363
-365	-1110	616	-519	-418	712	-858	869	427	760	941	-355
-365	-1121	615	-511	-424	718	-857	871	428	755	933	-358
-360	-1129	626	-516	-414	718	-866	877	413	755	938	-357
-359	-1134	623	-516	-409	719	-862	871	420	757	943	-359
-360	-1149	623	-517	-413	718	-859	874	410	756	931	-355
12	-85	374	-277	-190	441	-492	311	11	55	-225	-60
-10	-65	367	-284	-190	429	-449	314	9	44	-92	-67
-5	-38	369	-273	-197	439	-489	317	3	46	-95	-59
-11	-37	364	-291	-194	430	-497	321	11	44	-48	-59
-5	-19	368	-270	-199	442	-491	332	11	44	-88	-55
-2	-13	371	-271	-181	439	-491	313	9	45	-89	-56
-10	-14	363	-288	-183	423	-490	316	14	30	-75	-68
-3	-5	372	-279	-173	425	-482	293	7	27	-77	-65
-1	-1	368	-273	-190	433	-487	304	15	37	-78	-64
1	2	371	-270	-178	435	-494	299	14	41	-87	-59
3	4	374	-273	-170	427	-484	280	19	31	-80	-67
-489	-1532	747	-594	-474	854	-1037	876	443	954	989	-500
-489	-1544	746	-590	-477	854	-1039	876	445	856	993	-504
-493	-1541	739	-596	-468	839	-1029	865	453	848	1005	-515
-480	-1550	755	-589	-462	853	-1027	868	452	860	987	-501
-484	-1550	749	-600	-464	841	-1029	866	480	851	1001	-513
-485	-1549	749	-598	-466	842	-1024	868	452	849	1004	-513
-483	-1551	742	-585	-469	855	-1037	870	440	864	988	-503
-477	-1558	760	-599	-456	856	-1027	874	450	860	994	-504
-484	-1548	747	-587	-470	849	-1036	882	443	861	994	-498
-473	-1550	750	-584	-492	851	-1017	886	460	848	1001	-516
17	-96	374	-261	-148	427	-509	308	15	-280	999	-65
17	-86	369	-255	-150	433	-506	317	13	-266	-601	-58
19	-88	378	-262	-143	427	-502	309	11	-266	-596	-55
14	-76	367	-254	-151	429	-501	305	19	-277	-592	-63
17	-79	371	-258	-146	426	-499	300	16	-276	-593	-66
16	-79	371	-257	-156	428	-510	314	15	-270	-596	-63
8	-66	366	-271	-155	417	-524	274	-1	-268	-608	-53

Table (4.2) - Digitized Electric Signal for One of the Calibration Tests.

```
ZERO FOR TEST WD=50 S=45 (C
  12     19    150      0       3       0
 .85     24   -102    -20      28     -10     .82     38     -13      64      52     -54
 .84     24    -94    -25.     43      -1      .65     .38    -14      76      40     -45
 .87     36   -107    -11.     31      11     -108     -60    -10      65      61     -55
 .79     26   -102    -20      27       9      104     .45    -12      73      76     -60
 .81     26   -102    -19      25       8      106     -32     -9      75      77     -55
 .85     26    -99    -13.     23      15      102      69    -19      76      66     -57
 .83     24   -101    -27.     33       6     -106     -74    -21      79      50     -49
 .83     28   -100    -17.     23       9      111      46    -16      62      47     -53
 .85     25    -92    -26.     33       4     -119      52    -13      76      34     -45
 .82     27   -102    -19.     30       4      119      .46     -6      67      51     -48

TEST FOR STRU (C) WD=50 S=45 W=10 FREE(A)
  12    128      1       0       3       0
-201   -241    -26    -142    -150     135     -50     374     667    -441    1442     653
-192   -232    -30    -139    -156     179     -52     359     676    -244    1447     557
-175   -218    -36-   -132.   -157     178     -52     348     670    -121    1446     490
-156   -208    -27    -138.   -142     165     -51     316     658       4    1437     431
-137   -194    -32    -135.   -136     156     -45     308     655.    119    1433     359
-115   -178    -45    -126.   -127     143     -32     275     646     222    1440     292
 -84   -154    -53.   -115.   -116     133     -23     238     639     331    1436     219
 -53   -123    -61    -101.   -103     113       3     188     639     425    1443     136
 -12    -96    -69    -86.    -99      102     -13     144     617     517    1441.     50
  22    -51    -77.   -78     -79       79      38      81.    610     580.   1439     -25
  64     -6    -104.  -54     -76       77     -51      22     586     626.   1448    -120
 103     38    -118.  -42     -66       62     -74     -37     575     649.   1450    -198
 144.    78    -132   -35     -52       49     105     -90     546     657.   1442    -265
 171    117    -156   -22     -30       18     146    -149.    542     624    1438    -334
 185    152    -176    -8     -13      -10     128    -188     522     582    1442    -408
 192    190    -212    23      -5      -30     207    -228     505     539    1447    -484
 202    224    -229    45       9      -62     242    -275     423     482    1444    -542
 219    263    -260    60      31     -104     275    -329     479     411    1448    -599
 241    276    -267    75      48     -127     294    -352     458     362.   1436    -639
 263    294    -276    83      72     -160     314    -383     445     300    1435    -678
 278    307    -287   105      88     -176     320    -390     429     235.   1440    -720
 286    317    -296   109     113     -198     335    -410     423     178    1445    -739
 295    320    -296   114     125     -205     344    -424     405     123    1443    -759
 301    325    -302   124     121     -200     338    -425.    386      73    1442    -771
 307    329    -311   130     125     -203     341    -437     372      14    1448    -789
 308    333    -306   130     133     -208     346    -437     351     -35    1439    -780
 308    334    -309.  131     138     -219     356    -452     345.   -108    1446    -785
 310    341    -312   131.    139     -219     359    -453.    321    -173    1445.   -774
 312    334    -305.  131.    139     -214     353    -439     289    -230    1437.   -757
 309    340    -316   141     134     -211     353    -438     278    -307    1444    -748
 311    343    -308.  133     142     -212     355    -445     262.   -380    1438    -714
 311    341    -309   135     143     -219     366    -453     237    -467    1445    -685
 308    345    -310.  143     140     -211     360    -445     242    -542.   1446    -658
 296    342    -304   133.    145     -215     362    -445     230    -624.   1438    -603
 275    348    -314   141.    145     -216     366    -449     240    -715    1449.   -564
 253    339    -311   136     151     -223     373    -424     246    -796    1450.   -507
 229    326    -304   138.    145     -211     361    -383     255    -863    1446.   -450
 204    302    -289   123.    149     -210     353    -353     272    -942    1443.   -376
 191    221    -276   108.    136     -196     324    -308     280    -999    1437.   -305
 161    243    -264    99.    120     -171     293    -262     300   -1057    1438.   -235
 130    208    -245    77.    102     -151     222    -227     338   -1119    1442.   -158
  94.   120    -225    65.     75     -113     231    -163.    356   -1159    1438     -84
  51    136    -214    54.     50      -86     203    -112     390   -1205    1452.    -10
  12     86    -130    24.     41      -68     179     -56     408   -1228    1439      77
 -27     47    -162.    8.     19      -40     147      -3     440   -1254    1439.    164
 -61     10    -149   -14.      1      -22     126      56     472   -1264    1439     254
 -93    -25    -126   -23.    -29       17      94     128     492   -1262    1440.    342
-123    -63    -117   -42.    -50       36      77     155     528   -1260    1451     431
-143    -113    -90.   -64.    -63       57      53     200     541   -1217    1436.    521
-164    -153    -89.   -68.    -94       92      36     215     571   -1177    1449     582
-182    -183    -70    -93.   -102       99      20     238     627   -1105    1442.    649
-193    -212    -56.   -107   -122      124      -9     220.    601   -1016.   1439     691
-208    -214    -52.   -116   -139      146     -25     291     624    -932.   1454.    712
-219    -229    -44.   -129   -133      149     -30     300     645    -838.   1449.    731
-225    -232    -34.   -136   -149      162     -40     317     653    -727.   1444.    735
-226    -243    -34    -147   -148      168     -46     329     682.   -619    1445     722
-224    -244    -33.   -139   -167      186     -61     344     659    -502    1444     683
-212    -244    -27    -147   -160      179     -53     332     673    -371    1440     646
-199    -236    -26.   -144.  -162      185     -55     338     671    -238    1439     591
-183    -225    -26.   -144.  -155      178     -53     316     671    -130    1440     531
-162    -219    -32.   -136.  -159      183     -63     307     658     -17    1443     456
-143    -204.   -36.   -129.  -151      169     -54     287.    662      97    1447     380
-120    -190    -39.   -130.  -137      156     -46     263     649     214    1439     316
 -93    -165    -47.   -121.  -123      134     -28     238     644     313    1439     240
 -62.   -136    -64.    -99.  -117      122     -14     198     639     410    1450     155
 -27    -102    -68.    -96.  -107      118       3     148     621     498    1446      68
```

Continue

11	-62	-77	-31	-63	95	28	92	618	575	1440	-1
57	-26	-83	-60	-64	37	45	38	600	630	1440	-84
97	33	-112	-42	-63	76	67	-33	504	655	1449	-173
135	76	-129	-32	-55	52	94	-73	464	661	1444	-247
169	114	-143	-19	-35	30	123	-118	549	644	1442	-313
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215	252	-235	58	33	-82	220	-305	495	445	1441	-583
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298	333	-290	131	116	-184	349	-420	425	154	1450	-767
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17	100	-194	35	33	-71	188	-89	422	-1230	1448	59
-19	47	-172	5	25	-58	169	-28	449	-1253	1442	150
-53	-10	-151	-10	-2	-28	141	45	466	-1265	1439	237
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-116	-84	-120	-40	-53	-28	85	155	519	-1262	1446	411
-141	-119	-104	-58	-74	-54	52	196	540	-1234	1448	491
-164	-150	-83	-76	-93	-71	27	225	556	-1187	1441	568
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-197	-200	-57	-110	-120	-102	-3	284	594	-1043	1438	676
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-219	-240	-25	-158	-154	179	-52	340	671	-408	1441	662
-210	-238	-34	-155	-157	177	-56	330	688	-287	1438	609
-193	-226	-33	-147	-156	175	-53	315	677	-187	1450	541
-176	-216	-42	-139	-161	175	-65	312	662	-62	1449	469
-158	-210	-36	-148	-145	154	-56	272	665	49	1438	406
-137	-191	-51	-130	-144	153	-56	264	648	160	1440	323
-108	-170	-51	-129	-127	138	-34	229	604	260	1446	258
-79	-142	-63	-107	-123	130	-20	192	646	365	1449	170
-43	-111	-72	-92	-112	118	-2	158	632	460	1449	85
-7	-76	-76	-88	-99	90	23	98	619	535	1445	9
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-152	-120	-102	-67	-72	59	46	195	535	-1263	241	526
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-240	-238	-43	-143	-145	145	-38	332	625	-896	284	744
-247	-244	-39	-141	-160	168	-42	351	639	-785	306	733
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-244	-241	-40	-146	-167	187	-58	354	657	-559	332	689
-232	-246	-30	-154	-165	192	-62	360	663	-427	335	657
-223	-242	-25	-157	-160	192	-56	348	671	-310	354	602
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-85	-147	-57	-113	-115	126	-14	209	641	359	374	178
-48	-113	-69	-97	-105	112	4	161	635	455	373	92
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-235	-246	-42	-141	-153	172	-49	362	649	-706	318	711
-232	-252	-34	-147	-161	186	-56	368	656	-596	320	688
-225	-251	-28	-150	-157	183	-52	355	670	-485	322	646
-214	-242	-29	-148	-161	187	-59	352	678	-354	326	598
-201	-239	-25	-150	-152	175	-55	337	671	-230	313	552
-184	-226	-23	-149	-150	169	-54	324	669	-114	314	487
-166	-212	-38	-139	-151	171	-54	305	699	-3	333	400
-146	-198	-38	-141	-138	151	-39	276	665	112	330	342
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-88	-155	-57	-111	-125	138	-18	212	650	333	351	178
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292	322	-295	125	126	-203	355	-488	410	92	214	-744
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239	340	-321	135	142	-222	365	-431	245	-811	51	-510
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-215	-227	-34	-138	-153	188	-40	373	641	-607	293	719
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-87	-152	-55	-106	-125	157	-19	248	635	313	340	198
-53	-126	-55	-96	-104	137	-3	202	632	406	335	124
-17	-94	-65	-84	-93	129	10	168	608	494	340	31
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295	330	-286	129	110	-121	346	-385	417	188	259	-714
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-196	-200	-67	-109	-115	119	-7	306	570	-1045	271	697
-212	-215	-63	-116	-134	141	-19	318	597	-976	278	711
-223	-226	-49	-130	-139	157	-30	332	612	-874	272	733
-228	-243	-39	-150	-143	169	-41	350	619	-746	256	734
-227	-239	-45	-149	-152	180	-45	349	640	-644	259	713
-224	-244	-31	-157	-151	185	-51	362	635	-518	238	696
-212	-242	-35	-152	-159	185	-54	349	641	-399	249	645
-210	-237	-34	-158	-154	178	-58	353	643	-272	246	601
-192	-231	-31	-157	-152	179	-58	344	641	-163	250	535
-174	-221	-35	-150	-148	168	-44	322	650	-52	255	473
-151	-205	-33	-154	-140	158	-41	305	640	69	263	396
-125	-187	-48	-133	-139	163	-39	283	639	165	290	310
-99	-169	-55	-133	-120	142	-26	254	621	280	291	244
-70	-150	-56	-123	-107	127	-8	212	613	379	300	154
-34	-112	-68	-108	-91	115	8	185	604	455	314	70
2	-75	-87	-85	-89	106	20	122	590	522	338	-30
42	-32	-99	-79	-69	91	49	58	562	574	336	-98
84	17	-119	-54	-63	62	74	-1	564	602	344	-184
122	60	-141	-28	-59	57	85	-45	538	610	356	-262
151	100	-161	-11	-44	37	119	-93	528	585	370	-333
169	137	-179	1	-26	14	150	-137	506	555	369	-404
175	174	-207	12	-7	-22	196	-192	489	522	355	-459
190	211	-230	29	14	-49	235	-243	473	483	347	-519
206	247	-253	47	21	-70	268	-270	458	426	337	-576
235	278	-264	62	49	-101	298	-348	456	366	326	-616
257	295	-280	91	55	-121	313	-367	435	315	310	-663
276	310	-286	98	88	-157	333	-388	417	262	278	-683
288	312	-285	109	109	-175	342	-399	406	211	256	-709
298	332	-289	121	120	-191	345	-405	389	157	239	-737
304	335	-296	127	130	-187	350	-412	383	107	211	-740
308	336	-299	135	129	-191	354	-427	361	46	193	-761
308	344	-307	138	132	-196	361	-433	344	-10	158	-765
315	345	-297	129	149	-208	371	-454	326	-68	122	-752
312	347	-303	140	140	-197	362	-436	300	-131	106	-756
311	349	-304	144	138	-201	367	-442	288	-210	92	-757
312	350	-304	135	148	-208	372	-440	267	-275	63	-734
313	348	-299	132	152	-214	378	-434	245	-346	51	-716
310	352	-303	136	153	-219	383	-440	238	-440	44	-698
306	353	-308	141	145	-210	382	-446	228	-520	41	-677
295	355	-305	143	143	-210	380	-438	203	-606	37	-640
276	350	-307	141	147	-210	372	-433	218	-681	29	-597
253	336	-303	133	154	-217	389	-426	225	-767	22	-536
233	314	-300	134	152	-218	385	-398	235	-845	25	-479
214	302	-293	128	144	-209	363	-336	241	-921	36	-417
193	283	-285	119	129	-189	334	-284	257	-988	53	-349
167	259	-267	109	113	-167	308	-241	280	-1050	69	-275
133	234	-249	90	101	-140	226	-205	313	-1121	98	-199
93	190	-225	66	92	-124	257	-171	336	-1162	103	-103
56	157	-208	56	67	-89	225	-127	359	-1202	127	-27
17	119	-192	41	44	-63	196	-72	388	-1239	149	52
-15	74	-168	12	35	-43	176	-7	420	-1271	163	147
-46	29	-154	-2	8	-18	156	67	455	-1294	192	239
-75	-24	-135	-19	-17	14	122	138	471	-1285	206	333
-104	-67	-114	-46	-27	29	105	187	495	-1270	212	434
-130	-107	-100	-52	-60	66	77	227	553	-1250	235	513

FIXEDFA

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272	346	-313	142	139	-195	368	-497	144	732	-772	-783
278	345	-311	137	150	-201	376	-504	157	636	-820	-773
281	348	-311	139	151	-204	382	-507	160	536	-853	-762
280	349	-311	139	150	-204	377	-510	161	431	-893	-739

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279	330	-322	147	139	-197	372	-509	162	305	-908	-724
280	350	-316	138	151	-207	376	-514	166	166	-922	-692
278	351	-317	144	145	-197	320	-506	165	19	-918	-664
266	352	-318	143	144	-199	321	-512	174	-138	-909	-627
253	355	-319	145	142	-200	320	-512	174	-285	-891	-582
230	349	-317	144	147	-209	383	-516	185	-443	-863	-526
203	328	-323	135	146	-211	369	-498	179	-592	-833	-470
174	297	-312	128	147	-209	349	-451	178	-754	-790	-403
142	275	-303	124	126	-188	317	-406	179	-905	-733	-336
110	236	-285	103	113	-169	290	-360	190	-1053	-667	-261
76	204	-255	80	94	-139	256	-303	180	-1172	-614	-178
42	168	-240	65	72	-112	226	-259	183	-1298	-543	-96
5	130	-224	44	51	-84	206	-299	191	-1419	-469	-5
-33	86	-209	27	25	-54	172	-139	186	-1482	-393	75
-62	38	-189	10	5	-26	141	-69	183	-1508	-318	161
-101	-6	-160	-22	-3	-11	120	-5	182	-1517	-249	265
-135	-40	-150	-31	-35	22	92	53	110	-1524	-170	353
-165	-73	-126	-59	-50	41	74	96	176	-1524	-96	451
-182	-105	-116	-70	-77	69	46	126	175	-1526	-25	533
-204	-148	-90	-91	-90	91	24	149	168	-1522	42	614
-220	-182	-90	-92	-117	118	-6	180	158	-1518	107	663
-233	-206	-73	-104	-131	137	-19	202	150	-1517	165	712
-244	-225	-65	-115	-140	152	-39	231	156	-1519	224	729
-252	-232	-61	-131	-147	166	-39	244	162	-1522	284	737
-259	-241	-57	-136	-156	188	-49	269	157	-1496	328	728
-256	-248	-42	-152	-151	191	-54	274	152	-1431	364	713
-251	-249	-43	-156	-152	193	-55	220	152	-1356	402	684
-245	-247	-42	-153	-155	169	-58	269	162	-1250	428	642
-234	-240	-42	-143	-164	195	-67	222	152	-1130	446	586
-220	-234	-42	-142	-162	185	-70	268	154	-1010	457	523
-202	-223	-42	-146	-153	175	-68	247	154	-875	457	465
-182	-209	-44	-142	-140	162	-54	226	158	-739	458	397
-159	-191	-52	-131	-140	163	-55	209	151	-588	455	315
-134	-175	-60	-122	-127	148	-44	180	148	-438	425	251
-110	-152	-73	-105	-123	146	-39	206	142	-293	404	171
-79	-125	-78	-101	-108	121	-10	97	159	-157	381	85
-43	-89	-85	-80	-104	116	3	60	130	-8	345	5
-8	-49	-100	-78	-84	95	28	9	152	136	292	-62
33	-3	-116	-62	-69	73	52	-50	149	272	245	-151
68	35	-132	-50	-54	62	81	-115	150	391	190	-228
102	80	-161	-25	-45	47	105	-161	144	498	133	-304
125	117	-177	-9	-31	26	128	-203	141	608	53	-369
146	161	-202	14	-14	-1	167	-256	132	698	-7	-440
157	209	-230	32	8	-41	219	-321	141	760	-73	-506
174	243	-244	44	27	-70	254	-413	130	803	-159	-554
192	276	-274	76	36	-89	281	-413	128	817	-232	-611
220	299	-281	87	64	-124	312	-449	128	825	-316	-647
240	307	-293	107	77	-147	328	-469	132	818	-384	-694
252	319	-296	113	104	-173	343	-481	129	822	-461	-730
262	329	-309	132	109	-175	341	-484	126	833	-535	-756
268	335	-306	129	126	-189	352	-495	137	837	-609	-764
273	340	-308	126	135	-196	361	-507	138	815	-661	-782
274	344	-317	140	131	-194	358	-509	132	794	-725	-790
274	345	-313	139	135	-194	360	-506	138	736	-774	-788
275	343	-320	142	135	-200	362	-516	143	658	-815	-791
275	345	-315	138	137	-204	371	-525	149	570	-859	-771
277	349	-314	143	136	-198	368	-518	154	467	-888	-761
273	349	-316	141	140	-202	368	-518	160	340	-910	-742
274	343	-315	133	146	-212	376	-527	165	207	-921	-716
276	351	-319	142	137	-202	370	-517	165	68	-924	-689
276	353	-328	145	137	-202	379	-528	173	-82	-912	-660
267	354	-319	135	145	-208	376	-528	175	-230	-904	-610
249	345	-315	137	151	-213	381	-529	184	-392	-866	-563
222	337	-312	138	150	-210	378	-505	185	-543	-835	-504
192	324	-301	132	149	-211	372	-471	195	-701	-792	-437
164	300	-293	131	131	-190	340	-406	191	-848	-734	-376
135	275	-273	118	113	-167	309	-340	190	-985	-686	-283
104	237	-255	92	114	-152	289	-295	198	-1124	-624	-199
71	198	-241	80	84	-116	253	-248	195	-1245	-556	-112
32	157	-218	53	73	-89	222	-206	201	-1370	-481	-36
-9	115	-199	34	52	-65	196	-153	191	-1452	-423	57
-49	77	-177	26	22	-31	161	-82	188	-1489	-347	132
-85	40	-159	7	-1	-2	128	-11	186	-1511	-265	220
-116	-16	-134	-25	-15	19	107	35	187	-1519	-189	320
-142	-64	-125	-34	-39	46	82	59	189	-1526	-116	411
-164	-112	-109	-59	-53	65	61	106	181	-1518	-55	505
-185	-149	-97	-66	-83	101	29	146	158	-1516	14	586
-205	-124	-84	-76	-103	120	10	177	170	-1519	90	652
-221	-198	-62	-98	-108	123	5	200	175	-1525	158	693
-236	-210	-55	-106	-133	142	-18	246	163	-1518	204	733
-247	-225	-56	-119	-145	157	-32	249	159	-1516	256	756
-254	-235	-47	-133	-147	166	-41	258	128	-1519	315	747
-256	-244	-45	-139	-157	182	-53	265	160	-1466	351	734

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-253	-254	-38	-154	-150	133	55	259	155	-1388	380	719
-249	-256	-39	-147	-160	190	58	262	168	-1304	422	663
-243	-245	-36	-154	-157	187	58	253	164	-1189	439	613
-229	-232	-45	-142	-161	194	69	255	160	-1068	451	555
-212	-226	-40	-148	-152	179	-56	233	166	-951	467	489
-196	-212	-41	-144	-146	171	52	219	162	-814	469	418
-175	-202	-47	-134	-144	163	-49	196	157	-662	448	355
-150	-180	-62	-115	-141	166	48	173	155	-520	439	275
-124	-161	-60	-109	-126	153	-36	144	151	-366	411	205
-95	-136	-73	-94	-122	144	-25	103	148	-220	390	117
-61	-108	-78	-82	-99	115	8	45	142	-80	352	29
-23	-70	-97	-81	-97	104	22	-1	155	67	299	-43
15	-29	-103	-71	-74	79	47	-54	155	197	254	-131
56	18	-126	-51	-69	69	62	-102	148	327	205	-211
86	62	-144	-41	-51	44	122	-159	148	446	147	-285
113	112	-176	-17	-42	22	122	-207	138	550	86	-360
130	144	-189	-14	-15	-15	160	-318	139	641	12	-421
146	184	-216	5	-1	-41	197	-324	135	723	-58	-482
164	223	-246	30	16	-65	230	-372	125	794	-133	-540
185	256	-268	46	36	-99	275	-438	133	807	-205	-603
208	280	-284	67	56	-125	301	-423	129	816	-279	-649
231	299	-300	91	73	-146	314	-484	120	832	-363	-683
244	310	-304	101	93	-175	335	-508	132	824	-426	-722
253	320	-311	121	97	-178	332	-512	120	830	-504	-749
258	322	-313	114	119	-196	346	-523	125	835	-578	-762
262	333	-323	129	121	-197	345	-518	126	827	-636	-783
266	332	-323	131	121	-202	353	-532	133	791	-689	-800
266	335	-327	135	123	-202	353	-534	132	743	-741	-798
268	338	-322	133	132	-209	367	-546	148	678	-787	-794
269	339	-321	134	131	-203	355	-540	143	597	-841	-776
269	342	-329	138	128	-207	365	-549	155	485	-865	-776
270	341	-322	132	134	-210	361	-541	142	383	-905	-744
272	345	-316	131	143	-215	373	-552	166	246	-912	-726
271	346	-325	134	142	-217	375	-542	121	109	-910	-709
270	354	-333	144	131	-206	365	-538	170	-32	-907	-677
265	358	-326	145	139	-209	378	-541	184	-190	-889	-633
245	350	-330	147	137	-207	376	-535	182	-332	-875	-584
222	336	-325	142	142	-211	380	-536	193	-496	-842	-530
195	316	-318	143	140	-204	364	-498	184	-639	-813	-463
169	280	-297	130	144	-204	355	-442	194	-797	-766	-383
139	254	-262	113	131	-197	320	-390	188	-936	-719	-308

FIXED(A)

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-217	-164	-103	-90	-105	69	26	117	166	-1322	53	618
-232	-199	-72	-102	-129	118	55	52	167	-1519	121	676
-246	-219	-62	-116	-142	131	16	174	152	-1515	171	716
-257	-227	-61	-127	-148	143	-22	191	145	-1531	240	732
-262	-243	-53	-130	-151	159	-36	200	150	-1514	266	741
-269		-58	-139	-164	173	-46	220	153	-1508	320	717
-266	-254	-47	-143		180	-30	220	159	-1445	345	707
-264	-255	-45	-149	-172	190	-59	226	142	-1362	373	673
-258	-243	-41	-155	-164	135	-55	219	143	-1251	397	628
-242	-240	-36	-154	-164	130	-58	214	155	-1142	416	575
-232	-231	-41	-148	-161	177	-49	194	156	-1018	421	511
-214	-212	-46	-135	-167	182	52	186	136	-880	433	439
-191	-202	-46	-133	-154	172	49	170	144	-729	422	375
-162	-191	-54	-124	-149	165	46	148	150	-587	411	310
-140	-163	-67	-111	-144	151	30	124	148	-433	394	227
-111	-147	-77	-101	-132	136	-22	86	142	-297	363	152
-79	-116	-83	-98	-113	106	4	29	136	-157	336	64
-46	-82	-96	-76	-113	101	11	-14	142	-1	286	-13
-2	-37	-114	-70	-96	77	44	-92	150	130	245	-105
38	5	-124	-63	-80	58	68	-156	142	263	187	-177
76	51	-156	-38	-71	45	86	-206	127	390	123	-252
109	93	-172	-20	-63	27	113	-259	119	498	68	-326
123	134	-199	-13	-34	-18	154	-316	126	595	9	-398
139	181	-227	14	-21	-39	191	-367	131	688	-53	-473
154	213	-247	35	-5	-70	236	-421	122	779	-133	-522
171	247	-268	45	15	-99	273	-478	124	802	-202	-583
194	273	-286	64	33	-119	302	-512	124	817	-275	-626
215	281	-302	88	43	-136	315	-522	126	817	-337	-668
230	295	-308	99	67	-167	328	-528	121	828	-431	-718
240	310	-314	109	85	-181	332	-526	119	835	-504	-753
244	319	-325	122	90	192	341	-536	125	839	-571	-783
244	319	-328	119	100	-206	342	-544	123	812	-642	-792
239	322	-329	116	115	-218	356	-552	128	769	-701	-810
248	327	-329	121	111	-215	351	-552	136	718	-754	-817
250	322	-325	119	123	-226	362	-559	135	649	-813	-806
253	331	-330	124	115	-217	355	-554	140	570	-859	-793
254	333	-325	123	129	-227	368	-555	150	466	-893	-788
256	336	-329	127	126	-228	372	-565	159	354	-904	-766
253	338	-329	134	124	-216	365	-548	151	241	-923	-735

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-267	-251	-44	-154	-163	181	-49	248	155	-1435	347	691
-261	-250	-40	-154	-165	182	-50	244	161	-1353	384	647
-252	-251	-36	-150	-169	186	-52	248	153	-1233	395	614
-241	-239	-35	-145	-166	187	-54	240	164	-1128	420	554
-227	-232	-40	-147	-164	176	-50	225	162	-1006	424	496
-207	-219	-44	-146	-165	166	-45	212	159	-971	426	470
-184	-200	-55	-129	-162	172	-52	200	148	-721	412	364

FIXEQ(A)

12	128	1	0	0	0						
258	337	-323	135	122	-209	321	-584	144	576	-866	-776
260	337	-323	139	113	-206	370	-581	152	456	-898	-771
258	336	-327	137	120	-212	324	-597	162	327	-913	-755
258	341	-327	135	123	-216	378	-591	166	202	-923	-776
257	342	-325	136	121	-211	371	-578	154	76	-935	-704
254	341	-333	140	116	-210	366	-584	157	-69	-929	-676
-248	-244	-331	138	119	-211	371	-588	164	-217	-916	-642
231	341	-328	130	129	-223	384	-601	180	-386	-894	-599
204	332	-328	127	123	-228	329	-586	179	-541	-868	-547
181	321	-325	125	126	-229	371	-558	180	-695	-830	-490
154	298	-310	119	116	-211	346	-508	179	-845	-784	-421
124	266	-296	105	113	-203	325	-457	180	-991	-737	-340
91	229	-272	92	92	-170	284	-396	179	-1128	-688	-250
60	192	-256	75	71	-145	261	-356	185	-1264	-619	-182
26	164	-242	53	52	-119	239	-315	181	-1377	-565	-86
-9	123	-227	42	27	-85	208	-273	202	-1465	-494	-11
-46	77	-208	21	9	-66	186	-230	183	-1504	-417	66
-36	27	-186	2	-14	-40	154	-164	179	-1520	-344	156
-121	-31	-162	-20	-33	-18	126	-88	176	-1523	-274	244
-150	-74	-149	-35	-57	13	92	-18	163	-1513	-217	345
-127	-108	-128	-56	-72	36	27	38	168	-1512	-145	432
-198	-139	-114	-79	-87	52	62	68	170	-1523	-68	504
-218	-168	-102	-83	-117	82	34	92	164	-1528	3	573
-232	-191	-91	-94	-134	106	12	156	157	-1517	59	639
-248	-212	-82	-108	-142	117	-6	140	150	-1514	107	693
-260	-230	-69	-117	-161	142	-24	164	150	-1516	174	725
-268	-236	-64	-128	-169	158	-32	198	153	-1523	235	736
-272	-250	-50	-146	-161	155	-33	189	160	-1522	295	743
-270	-254	-50	-152	-167	169	-43	209	153	-1472	312	747
-265	-254	-47	-158	-174	174	-46	196	141	-1408	356	720
-257	-252	-43	-145	-182	189	-56	204	150	-1305	373	694
-252	-247	-44	-155	-176	172	-52	187	161	-1200	409	639
-240	-240	-39	-156	-170	169	-47	192	158	-1072	422	583
-224	-229	-54	-141	-180	175	-55	181	148	-935	416	523
-204	-216	-49	-144	-166	165	-47	162	150	-796	426	448
-184	-201	-57	-132	-164	160	-46	145	147	-651	409	376
-155	-179	-61	-124	-152	151	-32	115	153	-509	402	282
-129	-163	-74	-110	-142	136	-22	78	150	-359	374	145
-99	-136	-83	-98	-131	123	-4	41	149	-206	330	19
-67	-102	-94	-87	-120	108	12	-5	152	-63	293	-28
-35	-72	-102	-79	-104	90	30	-61	146	79	258	-69
-7	-31	-116	-67	-86	66	58	-127	153	208	211	-157
48	19	-139	-49	-75	55	27	-164	138	343	141	-226
85	67	-146	-37	-59	39	113	-237	143	450	89	-305
-115	-102	-170	-21	-45	24	138	-286	136	559	22	-374
134	142	-193	0	-32	1	172	-341	139	638	-41	-446
151	198	-229	25	19	-26	209	-392	131	721	-112	-522
163	228	-254	41	3	-62	248	-499	135	778	-182	-580
182	260	-279	63	21	-96	280	-492	131	809	-260	-629
202	285	-290	87	39	-127	303	-517	125	828	-338	-672
220	299	-300	101	66	-163	332	-545	130	827	-407	-713
236	313	-312	115	87	-190	337	-548	126	834	-496	-743
246	312	-306	115	104	-192	356	-560	139	827	-551	-770
248	325	-322	134	101	-191	348	-558	130	827	-627	-785
251	329	-320	126	116	-207	367	-576	159	795	-682	-804
254	329	-322	129	115	-203	361	-571	144	753	-745	-808
256	331	-321	130	123	-214	373	-580	150	685	-805	-803
257	337	-334	141	116	-205	364	-567	141	615	-857	-790
252	337	-320	134	125	-215	372	-584	159	507	-884	-780
257	340	-328	144	115	-205	367	-573	154	402	-914	-759
255	335	-321	133	129	-219	380	-581	168	266	-920	-744
255	341	-329	136	121	-213	372	-576	166	138	-927	-712
253	342	-336	143	116	-207	368	-577	165	-1	-924	-688
250	344	-332	133	128	-220	382	-595	179	-160	-910	-656
244	345	-332	140	120	-209	370	-588	170	-299	-907	-604
223	346	-341	143	117	-207	371	-589	176	-463	-877	-557
199	325	-332	129	126	-222	376	-589	185	-616	-848	-496
173	303	-329	136	115	-207	352	-519	178	-764	-812	-436
140	277	-308	115	120	-209	340	-472	183	-921	-754	-378
105	247	-288	99	106	-187	309	-410	190	-1063	-699	-303
70	222	-269	89	80	-148	265	-362	185	-1196	-635	-222
36	183	-245	65	64	-128	243	-318	181	-1312	-578	-128
9	153	-234	50	38	-95	212	-269	182	-1419	-511	-43

Continue

-29	116	-203	27	23	-73	186	-212	185	-1480	-440	39
-61	73	-193	9	5	-48	157	-141	175	-1498	-374	126
-97	30	-178	-6	-20	-16	124	-68	172	-1508	-290	213
-131	-19	-153	-32	-31	5	100	-15	182	-1522	-210	295
-162	-63	-146	-46	-60	35	70	26	176	-1514	-150	394
-187	-99	-126	-62	-78	54	58	49	172	-1524	-72	464
-207	-136	-106	-76	-99	80	33	88	162	-1514	-14	549
-225	-167	-89	-90	-113	97	21	108	172	-1522	54	611
-241	-194	-82	-102	-141	120	-4	146	163	-1522	115	652
-251	-217	-72	-113	-145	133	-11	160	166	-1526	177	694
-262	-235	-64	-123	-161	156	-33	184	164	-1516	218	728
-268	-249	-57	-136	-169	165	-42	196	152	-1513	259	741
-270	-252	-52	-147	-173	178	-50	199	161	-1502	318	726
-268	-263	-43	-157	-179	174	-45	192	161	-1440	348	719
-264	-258	-51	-147	-183	188	-61	202	153	-1351	372	697
-253	-255	-44	-148	-180	186	-63	199	159	-1239	390	659
-243	-249	-45	-158	-169	170	-54	183	161	-1134	417	603
-232	-236	-49	-146	-175	176	-63	181	155	-997	420	549
-212	-226	-50	-141	-164	165	-54	155	153	-858	417	485
-193	-212	-48	-142	-157	159	-45	132	163	-730	426	405
-172	-194	-61	-125	-160	154	-42	107	144	-576	406	335
-143	-176	-61	-126	-140	137	-27	73	163	-433	394	244
-112	-142	-79	-106	-142	136	-19	45	151	-280	361	159
-80	-120	-83	-98	-128	122	-5	-8	149	-131	326	77
-46	-84	-94	-83	-112	102	20	-41	149	12	288	-12
-8	-43	-109	-60	-105	89	44	-96	150	148	237	-96
31	4	-124	-52	-80	61	76	-161	157	266	195	-188
72	55	-140	-33	-62	46	101	-217	148	401	124	-257
108	97	-158	-23	-46	20	131	-254	149	500	66	-342
133	141	-189	5	-41	13	152	-317	140	597	-5	-412
150	177	-209	8	-13	-28	198	-381	107	671	-72	-469
164	218	-233	31	-2	-54	238	-434	146	748	-146	-556
181	254	-260	61	7	-68	269	-472	133	906	-234	-603
201	276	-269	66	36	-107	307	-517	139	811	-298	-658
221	294	-293	95	45	-127	320	-533	135	823	-368	-707
233	302	-297	104	66	-157	332	-544	129	933	-452	-728
243	322	-308	120	81	-180	344	-559	135	829	-520	-765
248	318	-301	-117	104	-199	359	-573	143	819	-589	-784
254	322	-311	129	105	-195	356	-568	144	805	-650	-803
257	333	-306	124	118	-210	363	-576	150	766	-711	-808
259	336	-317	130	123	-214	320	-581	151	711	-769	-802
258	341	-322	137	119	-210	367	-582	148	635	-820	-799
252	332	-315	130	123	-213	371	-592	150	546	-871	-780
260	343	-313	130	128	-216	375	-597	164	434	-888	-776
259	341	-321	130	127	-222	378	-601	179	316	-915	-757
258	340	-324	125	128	-226	380	-603	170	186	-920	-740
258	336	-320	130	129	-225	380	-604	173	47	-927	-706
253	343	-325	137	121	-214	371	-570	162	-84	-930	-665
242	348	-326	138	122	-218	374	-593	171	-234	-918	-627
222	346	-325	133	131	-224	386	-606	169	-395	-890	-589
198	337	-332	143	123	-218	380	-589	137	-554	-856	-536
178	322	-327	139	126	-219	-359	-553	185	-702	-823	-472
155	297	-306	122	126	-215	-340	-458	192	-852	-777	-407
133	224	-288	111	113	-191	-314	-445	191	-996	-720	-331
102	243	-274	93	97	-172	-291	-396	204	-1133	-664	-248

Table (4.3)- Digitized Electric Signal for One of the Tests.

CHAPTER FIVE

RESULTS

In this chapter the results of the laboratory experiments and the theoretical work required to determine certain structural parameters(Dynamic Properties) will be covered. The experimental data analysis for the six structures will be presented, and the results obtained from these tests will be shown.

5.1 THE DYNAMIC PROPERTIES OF THE STRUCTURES

5.1.1 DETERMINATION OF DAMPING COEFFICIENTS AND NATURAL FREQUENCIES

Damping coefficients were determined from the logarithmic decrement method as set forth by Den Hartog 1956[112]
For a single degree of freedom system, if an impulse load is applied to the mass centre, the sequential amplitudes of the oscillation are reduced by the damping in the system. For a system with linear, or viscous damping the relationship between succeeding peaks is expressed by:

$$\frac{X_{n+1}}{X_n} = e^{-2\pi\xi\ FN/FD} \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (5.1.1)$$

where X_n = magnitude of the n^{th} oscillation

X_{n+1} = magnitude of the next oscillation

ξ = relative damping coefficient

FN = natural undamped frequency of vibration

FD = natural damped frequency of vibration

and

$$FD = FN\sqrt{1 - \xi^2} \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (5.1.2)$$

For low damping as in the case of this work

$\xi^2 << 1$ so that

$$\ell n\, X_{n+1} - \ell n\, X_n = -2\pi\xi$$

or

$$\xi = \frac{1}{2\pi}\, \ell n\left(\frac{X_n}{X_{n+1}}\right) = \frac{1}{2\pi} \times \frac{1}{n}\, \ell n\left(\frac{X_o}{X_n}\right) \quad \ldots \quad \ldots \quad (5.1.3)$$

In the experiments the magnitudes of succeeding oscillations were obtained by constructing verticals at the upcrossings. The damping coefficients were determined from the smooth part of the record, after the higher modal vibration had disappeared. The decay trace was obtained by displacing the cantilever tip and releasing it. The first few cycles of motion from release were therefore occupied by the cantilever changing the imposed deflection shape to a curve compatible with a freely vibrating uniform loaded cantilever. For this

reason, the first three cycles from the start of the trace were ignored in the calculation. It was observed that over the remaining cycles, the amplitude ratio did not significantly vary with time.

Very little difference in the damping coefficient was observed from the two tests (impulse load test and tip displacement released test). The damping coefficients calculated therefrom were averaged. Each of these tests were repeated several times.

From the results, it was observed that the damping coefficients in air were less than in water and also it was noticed that the damping coefficients in water increased with the increase of water depth. This could have been due to the added hydrodynamic damping.

The natural damped frequencies of the structures were obtained although in the analysis the undamped natural frequencies were required but since the relative damping coefficients were always very small the damped frequencies were essentially equal to the undamped frequencies as can be seen from Equation (5.1.2).

Table (5.1) gives the average of the results of testing for the relative damping coefficients and the natural frequencies for the structures in both air and water.

Struct-ture	Water Depth Meter	Relative Damping Coefficient		Natural Frequency		Remarks
		In Air	In Water	In Air	In Water	
(A)	0.45	0.043	0.045	58.9	55.4	
	0.47		0.046		55.1	
	0.50		0.048		54.8	
	0.525		0.049		54.5	
(B)	0.5	0.040	0.044	36.94	34.25	
	0.525		0.045		33.76	
	0.55		0.047		33.4	
(C)	0.5	0.030	0.034	25.8	23.35	
	0.525		0.036		22.92	
	0.55		0.037		22.49	
(D)	0.5	0.028	0.030	21.2	18.93	
	0.525		0.033		18.54	
(E)	0.5	0.064	0.067	20.6	16.07	Zero load
		0.068	0.070	15.13	12.33	1010 grams
		0.072	0.074	12.69	9.8	2020 grams
	0.525		0.069		15.64	Zero load
			0.071		12.27	1010 grams
			0.076		9.45	2020 grams
(F)	0.5	0.0625	0.065	17.65	14.65	Zero load
		0.065	0.068	13.46	10.76	1010 grams
		0.070	0.073	11.52	8.6	2020 grams
	0.525		0.067		13.8	Zero load
			0.069		10.65	1010 grams
			0.075		8.3	2020 grams

TABLE (5.1) - Relative Damping Coefficients and Natural Frequencies.

5.1.2 DETERMINATION OF STIFFNESS CONSTANTS

The values of stiffness constant obtained from the experiments were substituted into the following equation

$$W_n^2 = K/M \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad \ldots \quad (5.1.4)$$

$$FN = W_n/2\pi$$

From the above equation, the analytical values of the natural frequencies were obtained. These analytical natural frequencies were compared with those obtained by experimentation during the tests for damping coefficients in order to provide an experimental check on the analysis.

The results are shown in Table (5.1).

Structure		FN Experimental	FN Analytical
(A)		58.9	59.08
(B)		36.94	33.9
(C)		25.8	25.0
(D)		21.2	20.1
	Mass Added at the top		
(E)	Zero	20.6	18.88
	1010 grams	15.13	13.41
	2020 grams	12.69	10.97
(F)	Zero	17.65	15.09
	1010 grams	13.46	10.89
	2020 grams	11.52	8.95

TABLE (5.1) The Experimental and Analytical Natural Frequencies

Also the values of the stiffness constant were substituted into the following equation

$$K = \frac{3EI}{L_s^3}(1-b) \quad \ldots \quad (5.1.5)$$

wehre L_s = structure height

$$b = \frac{4WL_s^2}{\pi^2 EI}$$

From the above equation the modulus of elasticity of each structure was calculated. Small discrepancies were found between the calculated and experimental values of the modulus of elasticity and also between the calculated and experimental values of the natural frequencies. These could be due to the possibility that the wall thickness away from the end may be different from that at the end which will affect the value calculated for the sectional moment of inertia I as well as the value of mass.

5.2 DATA ANALYSIS

The obtained data for each setting of the wave generator, the water depth in the tank and each level on the structure where the pressure can be measured were

(1) The wave profile

(2) The pressure round the structure at that level

(3) The tip displacement

(4) the bending moment at the base of the structure.

The above data were taken, first when the structure was free to vibrate (Free) and second when the structure was prevented from vibrating (Fixed).

The wave profiles obtained from free structure and fixed structure for each test were compared. The difference between the wave profiles was negligible, which shows that for the waves used in the tests and for the tested structures there were no radiation waves due to the movement of the structure.

For the structures with more than one level of pressure which can be measured, the data of the wave profile, tip displacement and the bending moment, obtained from the first level, were compared with the other levels for the same condition of test. There were no differences between these data which indicates that the wave properties did not change throughout the whole time of tests and the structures' vibrations were in steady state conditions.

Figures (5.2.1) and (5.2.2) show a typical data obtained for free and fixed structure respectively. Each figure shows the wave profile, the tip displacement, the

pressures at each of the eight points due to the wave and the in-line and transverse bending moments. They also show the calculated in-line and transverse local forces at that level.

The calculated static pressure ($\gamma\eta$) is plotted with the measured value of the pressure at point (1) (the static pressure is the curve without symbols). The static pressure is higher than the pressure due to the wave throughout the wave cycle as expected.

Figures (5.2.3) and (5.2.4) show the corresponding pressure distribution across the structure a time increment of the cycle equal to ($t = T/8$). From the figures it can clearly be seen that the distribution of pressure around the periphery of the cylinder cross-section is not symmetrical about the axis of wave propagation. This asymmetry is attributable to the asymmetric formation and shedding of eddies which will give rise to a transverse force.

Figures (5.2.5) and (5.2.6) show the corresponding Fast Fourier's Transformation to the wave profile, in-line force and tip displacement. From the figures it can be noticed that there is a magnitude of the second harmonic in the wave profile, and they show also that the wave profile did not change for the two sets of

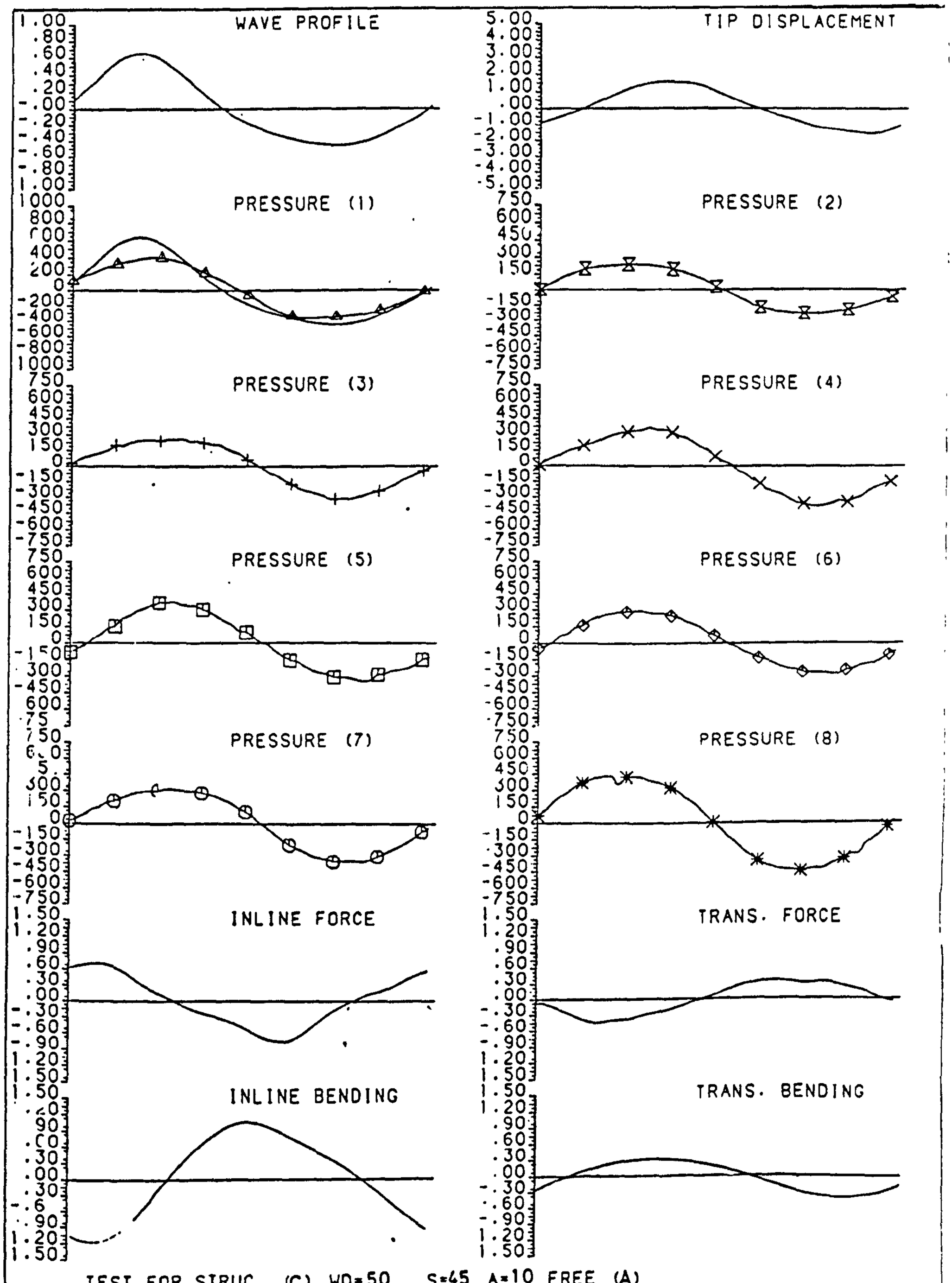

Figure (5.2.1) - Measured Data and the Calculated In-Line and Transverse Local Force for Free Structure.

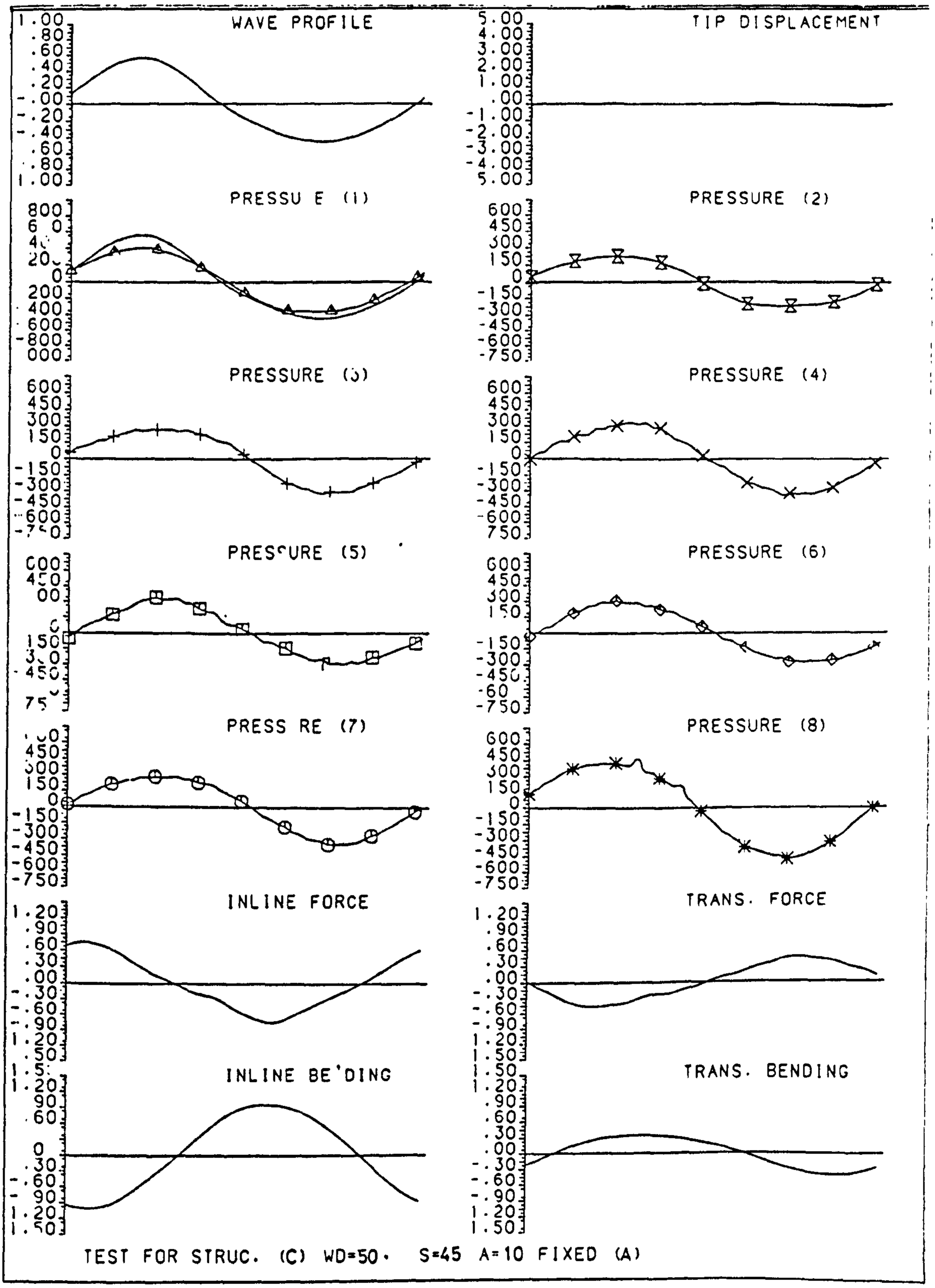

Figure (5.2.2) - Measured Data and the Calculated In-Line and Transverse Local Force for Fixed Structure.

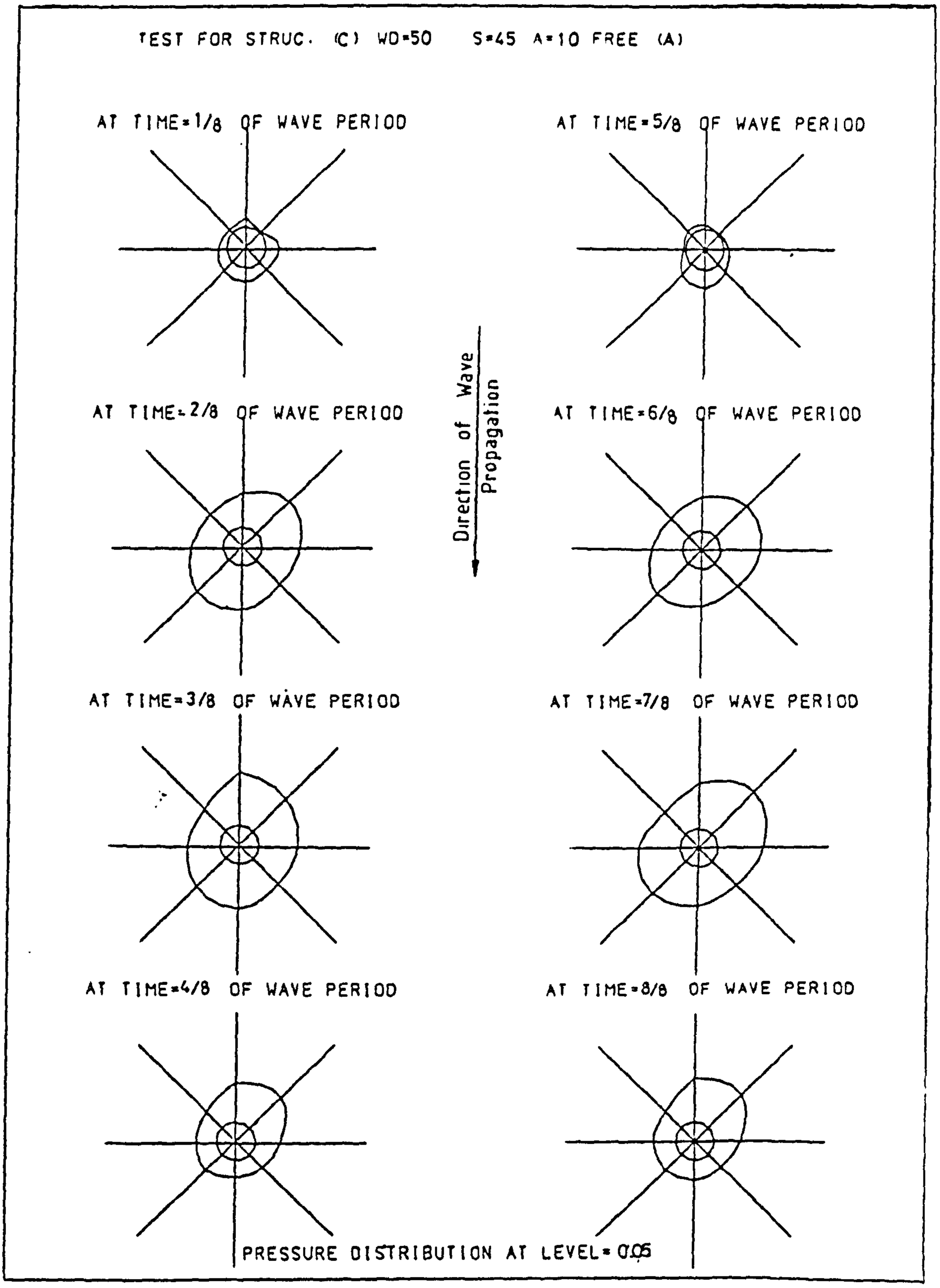

Figure (5.2.3) - Pressure Distribution across the Structure during the Wave Cycle for Free Structure.

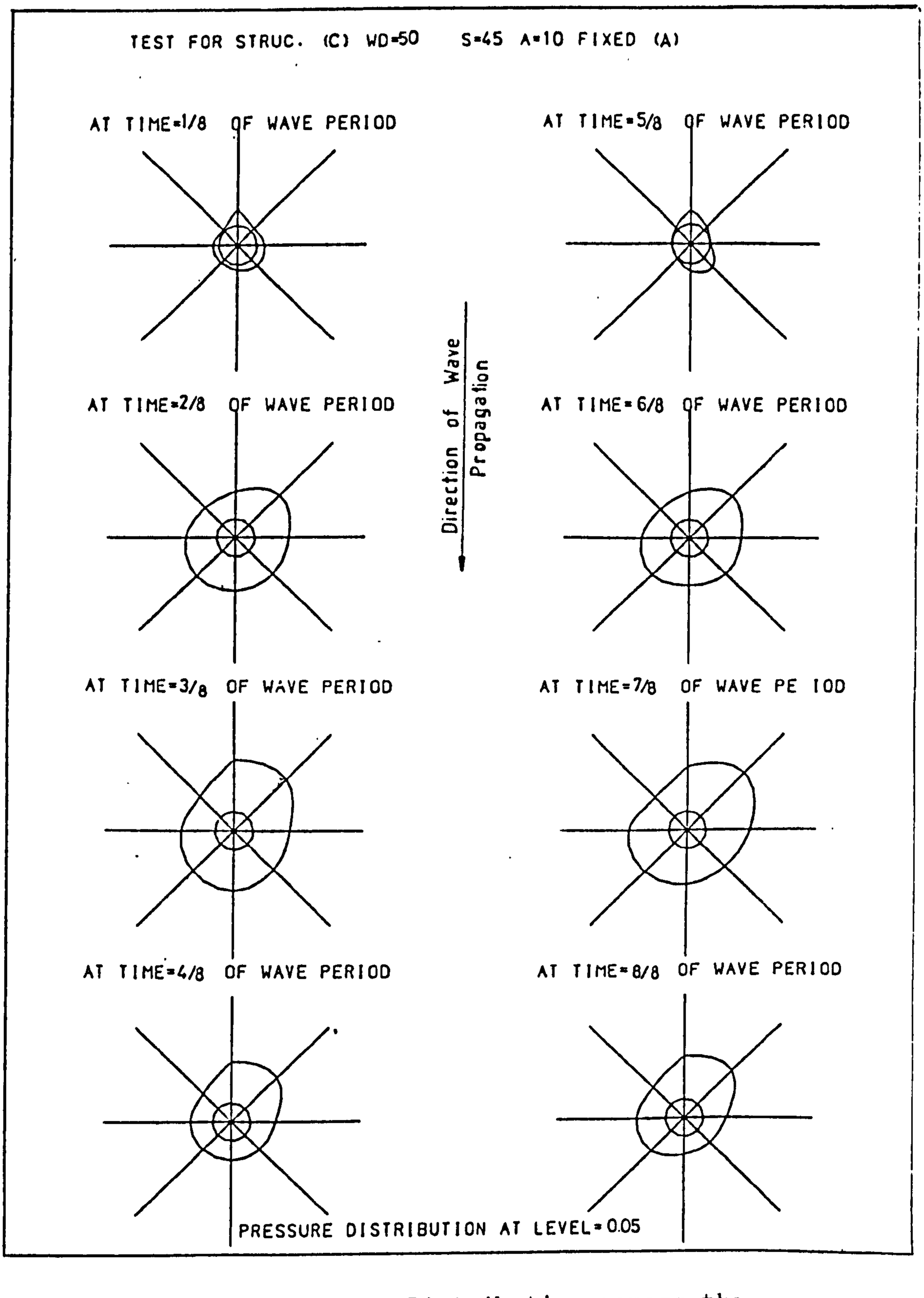

Figure (5.2.4) - Pressure Distribution across the Structure during the Wave Cycle for Fixed Structure.

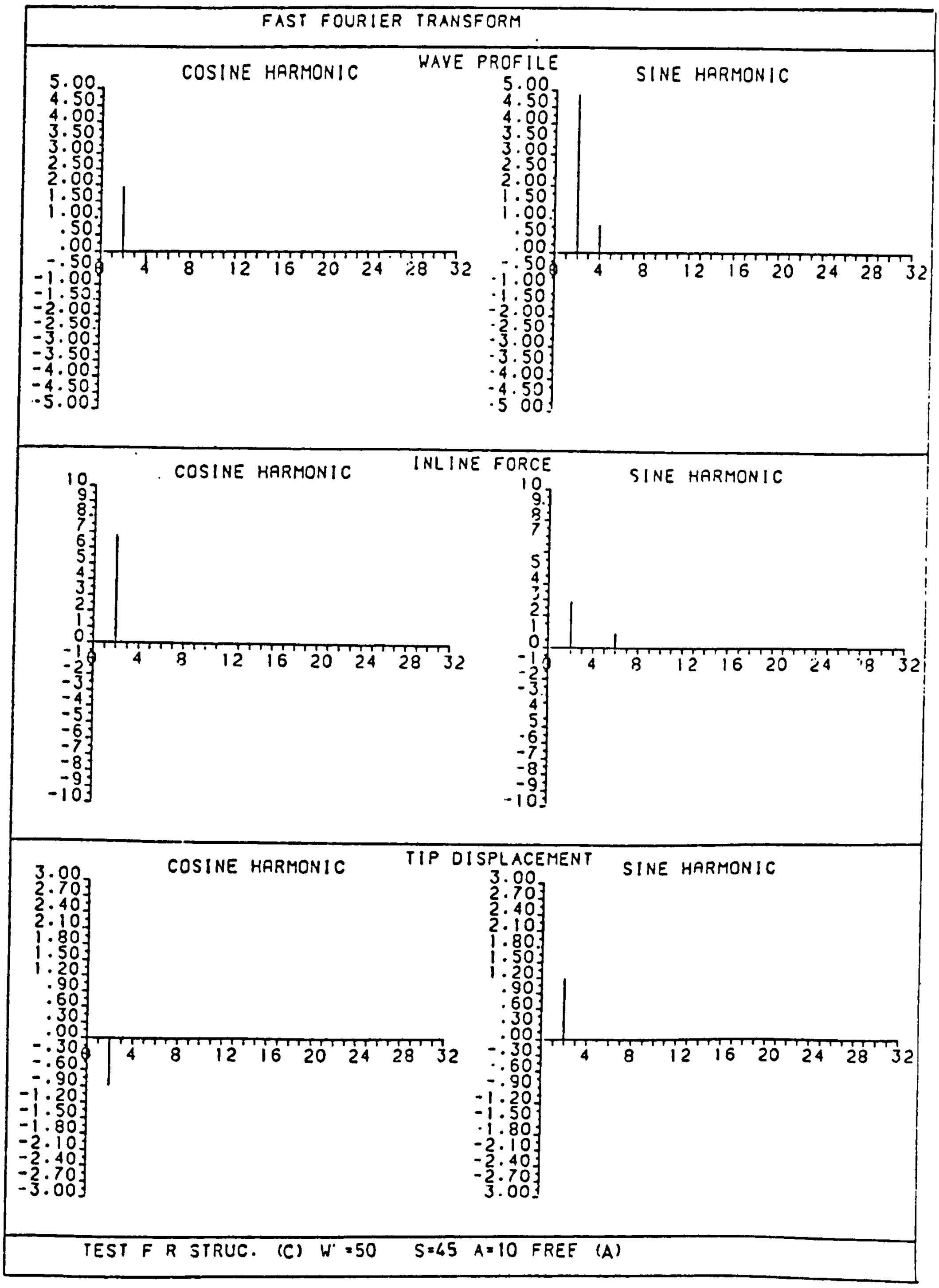

Figure (5.2.5) - Fast Fourier Transform for the Wave Profile, the In-Line Force and the Tip Displacement for Free Structure.

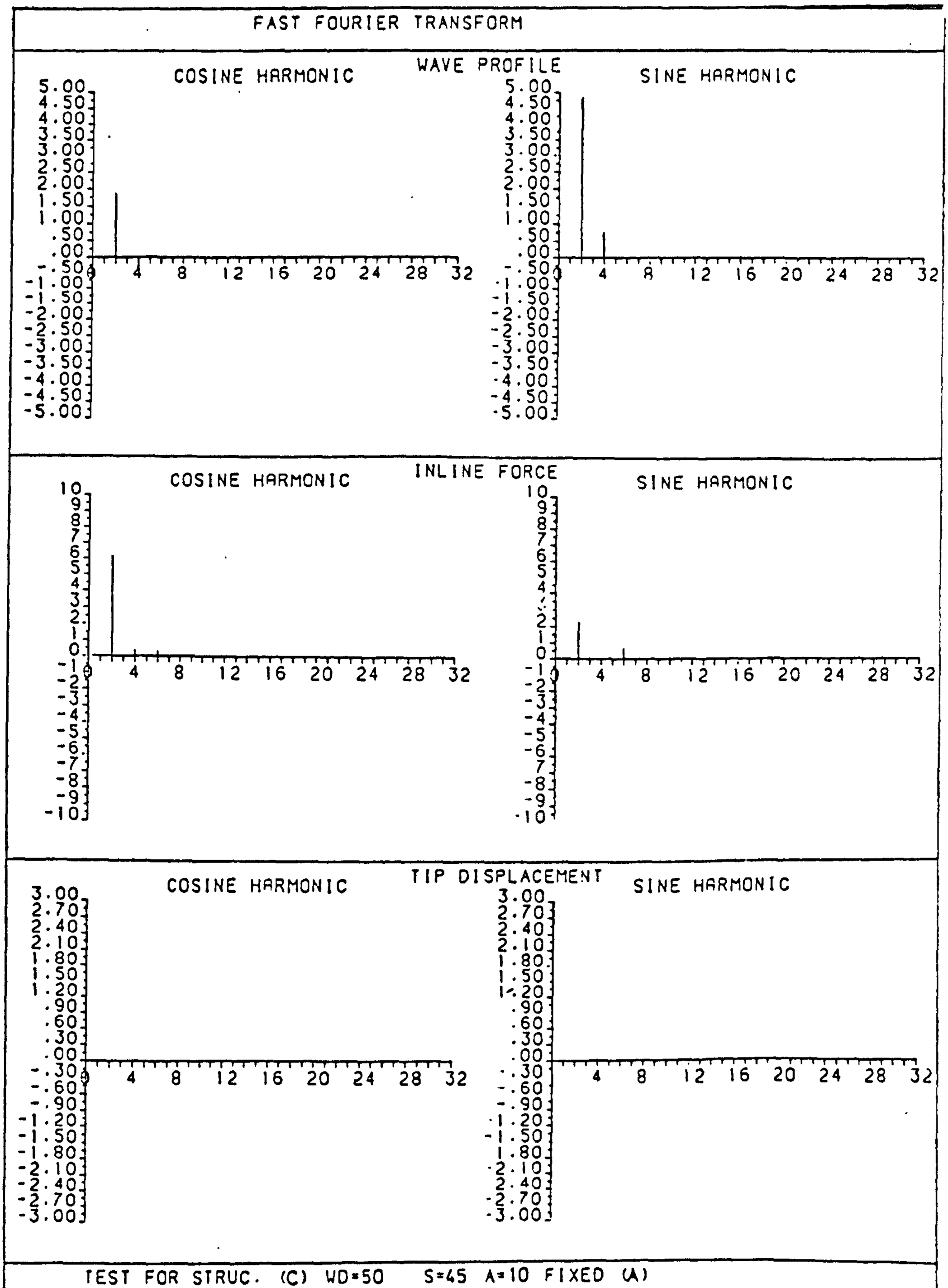

Figure (5.2.6) - Fast Fourier Transfor for the Wave Profile, the In-Line Force and the Tip Displacement for Fixed Structure.

tests. Also they show the difference in the force between the two cases, and the existence of the third sine harmonic in the force as it would be expected (see Keulegan and Carpenter[27]).

Detailed analyses were conducted for all tests. The analysis of the data obtained from each test may be divided into three sections. These three sections were wave properties, forces on structure and dynamic behaviour.

5.2.1 WAVE PROPERTIES

The measured wave profile was processed and wave properties were defined. The processing of the wave profile was carried out by digitizing the electric signal from the wave probe at 0.02 of a second. As the waves generated used in the tests were not purely sinusoidal waves as it is shown in Figure (4.4.1), sinusoidal interpolation was used to generate, from the number of data points obtained from one cycle of the wave, an exact 64 data points in wave cycle, in order to use the Fast Fourier Transformation. See Appendix (C) for sinusoidal interpolation.

From the F.F.T. the amplitudes of each sine and cosine harmonics contained in the wave profile were determined. Tne wave profile was represented mathematically by Equation (3.2.10).

The wave period was obtained directly from the time increment between the wave cycle. Wave height was defined as the difference between the crest elevation and the trough elevation.

Figure (5.2.7) shows a typical case of the calculated wave profile using Equation (3.2.10) and substituting for the first three harmonics only and the measured one.

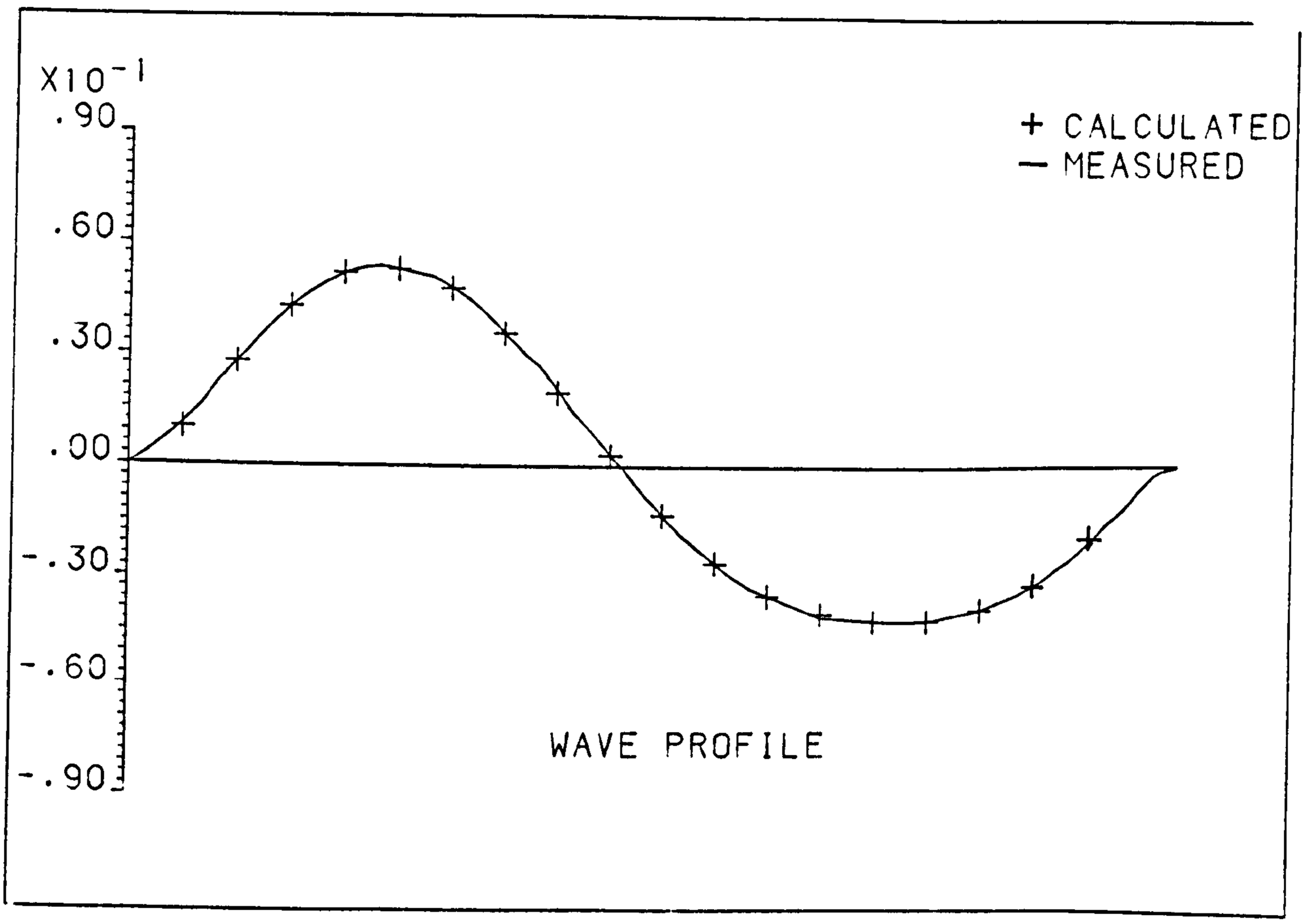

FIGURE (5.2.7) - The Measured and calculated Wave Profile.

Checks were made for the measured and calculated wave lengths using linear wave theory, from the Equation (3.2.13).

Equation (3.2.13) was solved by Newton's Raphson's iteration method getting the wave length of the first harmonic and comparing with the measured, the maximum difference was about 6%.

5.2.2 THE FORCES ON STRUCTURE

The signals from the eight pressure transducers located at one level of the structure were digitized at 0.02 of the second simultaneously with the wave profile, bending moment and tip displacement signals.

The force acting on the measured level of the structure at any time can be obtained from Equation (3.4.3). The numerical integration of this equation was carried out by using Simpson's rule:

$$dF(t) = \frac{1}{3}\,\Delta(2 \times \text{Odd value} + 4 \times \text{Even value})ds \qquad (5.2.4)$$

∴ the dF(t) in the direction of the wave propagation (in-line force) will be

$$dF(t) = \frac{1}{3} \frac{\pi D}{8} \left[(2 \times \delta P_I \cos \frac{\pi}{4}(I-1) + 4\ \delta P_J \cos \frac{\pi}{4}(J-1)\right] ds$$

... ... (5.2.5)

where I = 1, 3, 5, 7

J = 2, 4, 6, 8

The values of the $dF_x(t)$ for the wave cycle had been represented mathematically in the same way as of the wave profile.

The measured bending moments were digitized. The digitized bending moments were used to compare the measured and calculated. The bending moments were calculated by using Equation (3.2.20). The values of C_D and C_M appearing in the equation were calculated from the highest level and taken as constant along the length of the structure. The difference between the calculated and measured bending moments varied from one test to another with maximum difference of 17.5%. Some tests gave good agreement between the measured and calculated bending moments as shown in Figure (5.2.8).

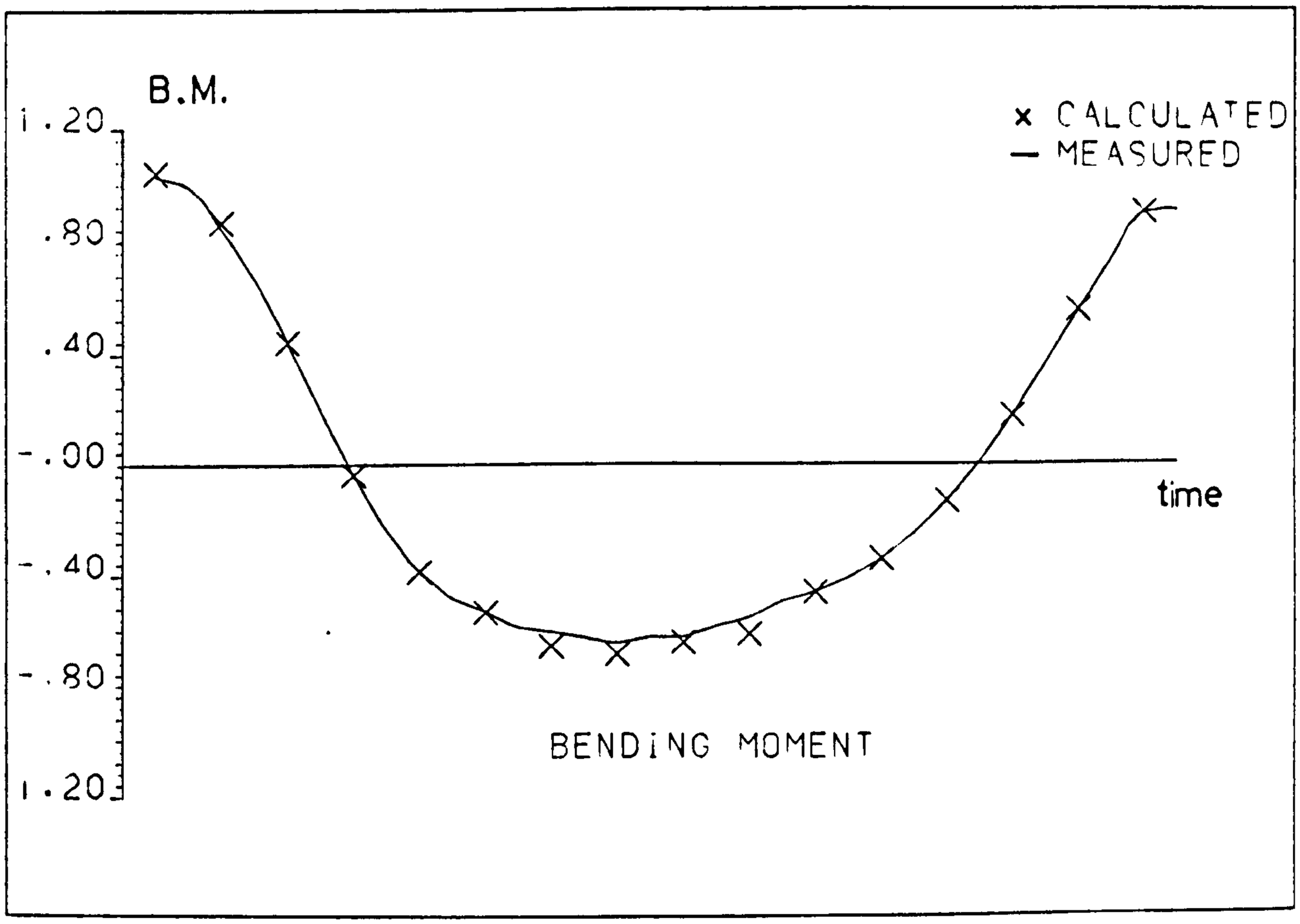

FIGURE (5.2.8) - The Calculated and Measured Bending Moments.

5.2.3 STRUCTURE DYNAMIC RESPONSE

The tip displacement values of x(t) had been digitized and represented mathematically in the same way as those of the wave profile and in-line force.

The mathematical representation of the tip displacement was used to calculate the structure velocity and acceleration at the measured level of pressure by Equations (3.4.7) and (3.4.8).

5.3 THE RESULTS FROM THE EXPERIMENTS

After the data had been analysed it was subjected to the appropriate theoretical approach to obtain the final results. The results obtained can be divided into two parts, the first part is the waves and the corresponding forces on the structures and the second part is the structural response due to the acting forces and the surrounding environment.

5.3.1 WAVES AND THE CORRESPONDING FORCES ON THE STRUCTURES

In the present study, the experimental waves were nearly sinusoidal, the position parameter (h/T^2) was varied between (0.247 - 0.564 m/sec^2) and the wave parameter (H/T^2) was varied between (0.079 - 0.201 m/sec^2). These values, according to the validity of the wave theory by Dean[8], lie within the region of the Airy theory. Also referring to Howell[5] the condition required to satisfy the linearization of the boundary condition is $H/2k^2h^3 << 1$, which allows for using linear wave theory without significant difference from the higher order. Linear wave theory was applicable as $H/2k^2h^3$ for this study was varied between (0.03 - 0.074). Finally according to the classification of the water wave length to the relative depth (h/L) the waves used

in this work can be classified as intermediate waves as h/L varies between (1/7 - 1/2.8).

From the above it can be concluded that the linear wave theory using the intermediate equation as shown in Chapter Three gives the most accurate theory to calculate the kinematic properties of the wave for the waves used.

As the structure used in this investigation had a small diameter relative to the wave length, the hydrodynamic forces were calculated by Morison's Equation (3.2.18) and the modified Morison's Equation (3.2.19).

The Keulegan-Carpenter number of the water wave was calculated on the basis of the maximum water velocity at the still water level. The forces corresponding to the wave acting on the structure were used to evaluate the values of the forces coefficients C_D and C_M. The average values of the drag and inertia coefficients for the cycle had been calculated according to a least squares fit between the experimental force history and that predicted by Morison's equation using Equation (3.2.18) and (3.2.19). Although the forces were measured at some distance from the still water level, the forces coefficients were related to the Keulegan-Carpenter number at the still water level.

For each structure the forces coefficients corresponding to the waves were represented as follows:

1. The forces coefficients C_D and C_M for fixed structures due to wave

2. The percentage ratio of the value of C_D for a free structure to that for a fixed structure (C_{DP}) and the same for C_M(C_{MP})

3. The percentage ratio of the values of C_D and C_M calculated by the modified Morison's equation to those calculated by Morison's equation (C_{DD} and C_{MD}) when the structures were free to vibrate.

The characteristics of the waves used in the tests and the corresponding forces as represented above for each structure are shown in Tables (5.2 through 5.11).

For the case of more than one level of pressure measured along the structure the tables show:

1. For the fixed structures the drag coefficients decrease with depth while the inertia coefficients increase with depth for the tested waves. These are due to the reduction of the water particle

velocity along the structure and consequently the reduction of the K-C number, as it shows later that the C_D decreases with the decrease of the K-C number while C_M increases with the decrease of K-C number, and due to the reduction of the vorticity with depth. The variation of C_D and C_M with depth should be taken into consideration for the determination of bending moment at the base of the structure especially for the case of deep water.

2. The percentage ratio of the drag coefficients for free structures to those for fixed structures slightly decreases with depth. The corresponding percentage ratio of the inertia coefficients behaves similarly. This can be attributed to the decrease of the effect of the vortices with depth, the forward movement of the separation points of the flow round the structure with depth and the reduced structural response with depth.

3. The percentage ratio of the drag coefficients calculated by modified Morison's equation to those calculated by Morison's equation increases with depth, similarly the corresponding ratio of inertia coefficients. These are due to the increase of the ratio of the relative velocity of the water particale to that of the structures' velocity, as the water particle velocity decreases exponentially with

STR. DIAMETER=0.1100 M STR. LENGTH=0.6080 M
NATURAL FREQ. IN AIR=58.000 DAMPING COEFF. IN AIR=0.043

TEST NO.	WATER DEPTH METER	WAVE HEIGHT METER	WAVE LENGTH METER	WAVE PERIOD SEC.	KEULEGAN-CARPENTER NUMBER	REYNOLDS NUMBER	LEVEL FROM SWL	CD FIXED	CM FIXED	CD FR % CD FX	CD FR(R) % CD FR	CM FR % CM FX	CM FR(R) % CM FR
1	0.500	0.0800	3.2500	1.6700	3.040	22026.0	0.110	0.784	2.037	110.10	100.1500	133.32	99.9500
2	0.500	0.0740	2.4400	1.3500	2.473	22167.0	0.110	0.716	2.239	107.37	99.9200	131.02	99.9400
3	0.500	0.1100	1.9150	1.1500	3.635	38244.0	0.110	0.826	1.963	110.22	100.6400	132.16	99.9800
4	0.525	0.0840	3.2000	1.6500	3.449	25289.0	0.135	0.810	1.967	111.25	100.6600	133.65	100.0600
5	0.525	0.0770	2.4930	1.3700	2.454	25798.0	0.135	0.709	2.252	106.86	99.8590	130.81	99.8000
6	0.525	0.1150	1.8880	1.1400	3.728	39569.0	0.135	0.839	1.942	108.22	100.7900	131.39	99.9500
7	0.470	0.0551	3.1545	1.6600	2.008	14638.1	0.080	0.662	2.411	100.00	100.0000	130.00	100.0000
8	0.470	0.0730	2.4220	1.3600	2.493	22176.9	0.080	0.719	2.214	107.77	100.1000	131.85	99.8600
9	0.470	0.1000	1.9165	1.1600	3.321	34642.9	0.080	0.800	1.980	111.02	100.4300	134.44	99.7000
10	0.450	0.0543	3.1502	1.6800	1.874	13498.7	0.060	0.654	2.548	100.00	100.0000	130.00	100.0000
11	0.450	0.0663	2.3899	1.3600	2.231	19850.1	0.060	0.686	2.314	104.42	99.7700	130.75	99.7200
12	0.450	0.0970	1.8977	1.1600	3.269	34094.4	0.060	0.799	1.993	110.82	100.0300	133.95	99.7410

Table (5.2) - Wave Characteristics and the Wave Force Coefficients for Structure (A).

STR. DIAMETER=0.1100 M STR. LENGTH=0.8080 M
NATURAL FREQ. IN AIR=36.940 DAMPING COEFF. IN AIR=0.040

TEST NO.	WATER DEPTH METER	WAVE HEIGHT METER	WAVE LENGTH METER	WAVE PERIOD SEC.	KEULEGAN-CARPENTER NUMBER	REYNOLDS NUMBER	LEVEL FROM SWL	CD FIXED	CM FIXED	CD FR % CD FX	CD FR(R) % CD FR	CM FR % CM FX	CM FR(R) % CM FR
1	0.500	0.0790	2.4400	1.3500	2.682	24039.0	0.068	0.742	2.106	113.75	100.0200	106.91	100.0200
2	0.500	0.1200	1.9510	1.1500	3.757	39534.0	0.068	0.850	1.922	110.79	100.5000	103.25	99.8400
3	0.500	0.0990	2.2070	1.2600	3.177	30510.0	0.068	0.792	1.988	112.74	100.2400	105.36	99.9300
4	0.525	0.0660	2.5020	1.3600	2.182	19417.0	0.093	0.668	2.291	114.15	99.7300	109.56	99.9000
5	0.525	0.1140	1.9610	1.1600	3.748	39100.0	0.093	0.838	1.934	111.07	100.5300	103.23	99.8000
6	0.525	0.1330	1.5200	1.0000	3.926	47506.0	0.093	0.856	1.900	108.37	100.7900	102.99	99.7600
7	0.550	0.0749	2.5336	1.3600	2.594	23082.0	0.119	0.726	2.226	113.90	100.2000	107.51	99.7780
8	0.550	0.1057	1.9771	1.1600	3.439	35876.1	0.119	0.821	1.977	112.05	100.1000	104.59	99.8400
9	0.550	0.1290	1.5274	1.0000	3.775	45677.5	0.119	0.852	1.917	110.58	101.3000	103.10	99.8800

Table (5.3) - Wave Characteristics and the Wave Force Coefficient for Structure (B).

STR. DIAMETER=0.0605 M STR. LENGTH=0.7400 M
NATURAL FREQ. IN AIR=25.800 DAMPING COEFF. IN AIR=0.030

TEST NO.	WATER DEPTH METER	WAVE HEIGHT METER	WAVE LENGTH METER	WAVE PERIOD SEC.	KEULEGAN-CARPENTER NUMBER	REYNOLDS NUMBER	LEVEL FROM SWL	CD FIXED	CM FIXED	CD FR %----- CD FX	CD FR(R) %-------- CD FR	CM FR %----- CM FX	CM FR(R) %-------- CM FR
1	0.500	0.0670	2.4670	1.3600	4.076	10969.4	0.050	0.881	1.806	115.71	100.8000	109.45	98.7500
							0.150	0.858	1.844	115.53	100.9000	109.37	99.3700
2	0.500	0.1050	1.9151	1.1500	6.127	19502.6	0.050	1.047	1.734	114.50	101.2000	104.55	99.2500
							0.150	1.026	1.768	114.20	101.4000	104.39	99.4800
3	0.500	0.1250	1.5130	1.0000	7.863	28782.3	0.050	1.118	1.699	111.30	101.7000	102.20	99.0800
							0.150	1.078	1.725	111.20	101.9000	102.22	99.4500
4	0.525	0.1120	2.4750	1.3500	6.697	18157.4	0.075	1.066	1.746	113.87	101.3000	103.68	99.0200
							0.175	1.046	1.748	113.74	101.4700	103.60	99.4000
5	0.525	0.1389	1.9330	1.1500	7.975	25381.7	0.075	1.110	1.679	111.13	101.8000	102.09	99.5400
							0.175	1.107	1.681	111.00	102.0000	102.17	99.5700
6	0.525	0.1801	1.5206	1.0000	10.604	38812.0	0.075	1.261	1.544	100.60	102.7000	100.20	99.5900
							0.175	1.207	1.559	100.52	102.6000	100.19	99.5200
7	0.550	0.1748	1.7516	1.0800	9.821	33286.4	0.100	1.154	1.579	104.00	102.5400	100.31	100.0500
							0.200	1.120	1.657	104.12	102.5400	100.29	100.0300
8	0.550	0.1920	1.5274	1.0000	10.391	38033.1	0.100	1.168	1.568	101.51	102.5120	100.22	100.3300
							0.200	1.162	1.638	101.31	102.6330	100.36	100.3000

Table (5.4) - Wave Characteristics and Wave Force Coefficients for Structure (C).

STR. DIAMETER=0.0605 M STR. LENGTH=0.8350 M
NATURAL FREQ. IN AIR=21.200 DAMPING COEFF. IN AIR=0.028

TEST NO.	WATER DEPTH METER	WAVE HEIGHT METER	WAVE LENGTH METER	WAVE PERIOD SEC.	KEULEGAN-CARPENTER NUMBER	REYNOLDS NUMBER	LEVEL FROM SWL	CD FIXED	CM FIXED	CD FR %----- CD FX	CD FR(R) %-------- CD FR	CM FR %----- CM FX	CM FR(R) %-------- CM FR
1	0.500	0.0930	2.1800	1.2500	5.524	16176.0	0.050	1.019	1.758	118.34	101.0000	109.28	99.9200
							0.150	0.978	1.837	118.27	101.2000	109.11	99.9700
2	0.500	0.1040	1.9150	1.1500	6.105	19431.0	0.050	1.044	1.740	117.76	101.1500	107.71	99.9900
							0.150	1.023	1.770	117.71	101.3200	107.67	100.1000
3	0.500	0.1480	1.3300	0.9500	8.399	32359.0	0.050	1.115	1.675	113.68	102.1400	104.68	100.3400
							0.150	1.109	1.679	113.62	102.2600	104.54	100.2700
4	0.525	0.0980	2.4800	1.3500	5.466	14820.0	0.075	0.982	1.813	118.45	100.9000	109.56	99.4680
							0.175	0.964	1.842	118.38	101.0000	109.46	99.6220
5	0.525	0.1190	1.9300	1.1500	7.576	24113.0	0.075	1.072	1.723	115.65	101.6000	105.34	99.6900
							0.175	1.058	1.733	115.44	101.7300	105.27	99.7600
6	0.525	0.1440	1.5200	1.0000	7.904	28929.0	0.075	1.136	1.688	114.93	101.6100	104.98	100.3700
							0.175	1.096	1.705	114.67	101.9000	104.82	100.3300

Table (5.5) - Wave Characteristics and Wave Force Coefficients for Structure (D).

STR. DIAMETER=0.0574 M STR. LENGTH=0.7350 M
NATURAL FREQ. IN AIR=20.000 DAMPING COEFF. IN AIR=0.064

TEST NO.	WATER DEPTH METER	WAVE HEIGHT METER	WAVE LENGTH METER	WAVE PERIOD SEC.	KEULEGAN-CARPENTER NUMBER	REYNOLDS NUMBER	LEVEL FROM SWL	CD FIXED	CM FIXED	CD FR % CD FX	CD FR(R) % CD FR	CM FR % CM FX	CM FR(R) % CM FR
1	0.525	0.1080	1.9600	1.1500	6.742	19148.2	0.120	1.054	1.719	111.18	101.3300	101.60	99.9000
							0.170	1.052	1.720	111.14	101.3500	101.56	99.9800
							0.270	1.050	1.731	111.12	101.3900	101.54	99.8000
2	0.525	0.1360	1.5200	1.0000	8.130	26787.7	0.120	1.111	1.672	107.68	102.1000	100.19	100.4000
							0.170	1.110	1.680	107.65	102.1200	100.17	100.3000
							0.270	1.104	1.690	107.64	102.2000	100.16	100.1000
3	0.525	0.1200	1.7400	1.0300	7.308	22295.7	0.120	1.062	1.725	109.75	101.4500	100.87	100.1000
							0.170	1.061	1.728	109.72	101.4000	100.84	100.0000
							0.270	1.057	1.730	109.70	101.4500	100.82	100.2000
4	0.500	0.0900	1.9200	1.1500	5.879	16698.6	0.095	1.010	1.790	112.27	101.0500	103.08	99.9000
							0.145	1.008	1.790	112.25	101.0000	103.04	99.8500
							0.245	1.002	1.801	112.24	101.0000	103.02	99.8600
5	0.500	0.1250	1.5100	1.0000	7.271	25955.3	0.095	1.062	1.700	109.95	101.3000	100.90	100.0700
							0.145	1.056	1.710	109.92	101.2000	100.88	100.2000
							0.245	1.054	1.730	109.93	101.3000	100.86	100.3000
6	0.500	0.1050	1.2800	1.0300	6.520	19893.4	0.095	1.048	1.738	111.45	101.2500	102.00	99.9500
							0.145	1.046	1.740	111.41	101.1000	101.95	99.8100
							0.245	1.044	1.750	111.41	101.1000	101.95	99.8000

Table (5.6) - Characteristics and the Wave Force Coefficients for Structure (E) with no Mass at the Top.

STR. DIAMETER=0.0574 M STR. LENGTH=0.7350 M
NATURAL FREQ. IN AIR=15.130 DAMPING COEFF. IN AIR=0.068

TEST NO.	WATER DEPTH METER	WAVE HEIGHT METER	WAVE LENGTH METER	WAVE PERIOD SEC.	KEULEGAN-CARPENTER NUMBER	REYNOLDS NUMBER	LEVEL FROM SWL	CD FIXED	CM FIXED	CD FR %----- CD FX	CD FR(R) %-------- CD FR	CM FR %----- CM FX	CM FR(R) %-------- CM FR
1	0.525	0.1080	1.9600	1.1500	6.742	19148.2	0.120	1.054	1.721	111.76	101.3300	102.17	99.9500
							0.170	1.054	1.723	111.73	101.3200	102.15	99.9500
							0.270	1.050	1.731	111.72	101.3300	102.12	99.9700
2	0.525	0.1360	1.5200	1.0000	8.130	26787.7	0.120	1.112	1.680	108.25	102.0500	100.67	100.4500
							0.170	1.111	1.682	108.22	102.0800	100.64	100.5000
							0.270	1.104	1.690	108.21	102.0700	100.63	100.5200
3	0.525	0.1200	1.7400	1.0800	7.308	22295.7	0.120	1.064	1.728	110.37	101.3800	101.42	100.0500
							0.170	1.063	1.729	110.34	101.4000	101.39	100.0700
							0.270	1.057	1.730	110.32	101.3800	101.38	100.0600
4	0.500	0.0900	1.9200	1.1500	5.879	16698.6	0.095	1.012	1.790	112.82	101.0000	103.67	99.9100
							0.145	1.009	1.790	112.79	101.0300	103.64	99.9200
							0.245	1.002	1.801	112.77	101.0400	103.62	99.9400
5	0.500	0.1250	1.5100	1.0000	7.271	25955.3	0.095	1.062	1.700	110.50	101.2800	101.50	99.9900
							0.145	1.057	1.715	110.46	101.3000	101.45	99.9900
							0.245	1.054	1.730	110.44	101.3000	101.46	99.9700
6	0.500	0.1050	1.2800	1.0800	6.520	19893.4	0.095	1.048	1.738	112.01	101.1000	102.57	99.9200
							0.145	1.046	1.740	111.97	101.1500	102.54	99.9200
							0.245	1.044	1.750	111.95	101.1400	102.53	99.9400

Table (5.7) - Wave Characteristics and the Wave Force Coefficients for Structure (E) with 1010 Grs at the Top.

STR. DIAMETER=0.0574 M STR. LENGTH=0.7350 M
NATURAL FREQ. IN AIR=12.690 DAMPING COEFF. IN AIR=0.072

TEST NO.	WATER DEPTH METER	WAVE HEIGHT METER	WAVE LENGTH METER	WAVE PERIOD SEC.	KEULEGAN-CARPENTER NUMBER	REYNOLDS NUMBER	LEVEL FROM SWL	CD FIXED	CM FIXED	CD FR % CD FX	CD FR(R) % CD FR	CM FR % CM FX	CM FR(R) % CM FR
1	0.525	0.1080	1.9600	1.1500	6.742	19148.2	0.120	1.055	1.721	112.14	101.3100	102.79	99.9000
							0.170	1.051	1.721	112.10	101.3100	102.76	99.9200
							0.270	1.050	1.731	112.09	101.3000	102.75	99.9300
2	0.525	0.1360	1.5200	1.0000	8.130	26787.7	0.120	1.112	1.681	108.83	101.9000	101.10	100.3000
							0.170	1.111	1.683	108.80	101.9100	101.06	100.3000
							0.270	1.104	1.690	108.79	101.9000	101.04	100.3300
3	0.525	0.1200	1.7400	1.0800	7.308	22295.7	0.120	1.064	1.728	110.94	101.4200	101.97	100.1000
							0.170	1.062	1.727	110.90	101.4100	101.95	100.1000
							0.270	1.057	1.730	110.88	101.4100	101.94	100.1200
4	0.500	0.0900	1.9200	1.1500	5.879	16698.6	0.095	1.013	1.793	113.43	101.0300	104.25	99.8900
							0.145	1.009	1.791	113.40	101.0400	104.21	99.8800
							0.245	1.002	1.801	113.38	101.0300	104.19	99.9100
5	0.500	0.1250	1.5100	1.0000	7.271	25955.3	0.095	1.061	1.710	111.13	101.2700	102.09	100.1000
							0.145	1.060	1.720	111.10	101.3000	102.06	100.1000
							0.245	1.054	1.730	111.11	101.3000	102.05	100.1000
6	0.500	0.1050	1.2800	1.0800	6.520	19893.4	0.095	1.050	1.736	112.58	101.2000	103.15	99.9500
							0.145	1.045	1.740	112.55	101.2300	103.11	99.9700
							0.245	1.044	1.750	112.53	101.1800	103.13	99.9900

Table (5.8) - Wave Characteristics and Wave Force Coefficients for Structure (E) with 2020 Grs at the Top.

STR. DIAMETER=0.0574 M STR. LENGTH=0.8350 M
NATURAL FREQ. IN AIR=17.650 DAMPING COEFF. IN AIR=0.063

TEST NO.	WATER DEPTH METER	WAVE HEIGHT METER	WAVE LENGTH METER	WAVE PERIOD SEC.	KEULEGAN-CARPENTER NUMBER	REYNOLDS NUMBER	LEVEL FROM SWL	CD FIXED	CM FIXED	CD FR % CD FX	CD FR(R) % CD FR	CM FR % CM FX	CM FR(R) % CM FR
1	0.525	0.1080	1.9600	1.1500	6.742	19148.2	0.120	1.073	1.718	114.76	101.3500	104.35	100.0000
							0.170	1.068	1.720	114.74	101.4000	104.32	100.2000
							0.270	1.056	1.735	114.72	101.3000	104.30	100.2000
2	0.525	0.1360	1.5200	1.0000	8.130	26787.7	0.120	1.133	1.650	112.30	102.2000	102.38	100.4000
							0.170	1.128	1.650	112.25	102.4000	102.35	100.5000
							0.270	1.104	1.690	112.24	102.4000	102.33	100.2000
3	0.525	0.1200	1.7400	1.0800	7.308	22295.7	0.120	1.075	1.720	113.90	101.5000	103.35	100.1500
							0.170	1.074	1.723	113.88	101.1000	103.32	100.3000
							0.270	1.058	1.730	113.86	101.0000	103.31	100.4000
4	0.500	0.0900	1.9200	1.1500	5.879	16698.6	0.095	1.007	1.770	115.90	101.0700	106.11	99.9100
							0.145	1.008	1.780	115.87	101.1100	106.08	99.9500
							0.245	1.002	1.801	115.86	101.2000	106.07	99.9400
5	0.500	0.1250	1.5100	1.0000	7.271	25955.3	0.095	1.060	1.710	114.00	101.3200	103.53	100.1000
							0.145	1.056	1.714	113.98	101.2000	103.50	100.2000
							0.245	1.054	1.730	113.96	101.5000	103.50	100.3000
6	0.500	0.1050	1.2800	1.0300	6.520	19893.4	0.095	1.048	1.737	115.11	101.2400	104.76	99.9700
							0.145	1.050	1.740	115.08	101.6000	104.73	99.7000
							0.245	1.044	1.750	115.09	101.6000	104.72	99.8000

Table (5.9) - Wave Characteristics and Wave Force Coefficients for Structure (F) with no Mass at the Top.

STR. DIAMETER=0.0574 M STR. LENGTH=0.8350 M
NATURAL FREQ. IN AIR=13.460 DAMPING COEFF. IN AIR=0.065

TEST NO.	WATER DEPTH METER	WAVE HEIGHT METER	WAVE LENGTH METER	WAVE PERIOD SEC.	KEULEGAN-CARPENTER NUMBER	REYNOLDS NUMBER	LEVEL FROM SWL	CD FIXED	CM FIXED	CD FR % CD FX	CD FR(R) % CD FR	CM FR % CM FX	CM FR(R) % CM FR
1	0.525	0.1080	1.9600	1.1500	6.742	19148.2	0.120	1.096	1.710	115.37	101.3100	104.98	99.9500
							0.170	1.092	1.710	115.34	101.2000	104.95	99.9900
							0.270	1.068	1.725	115.32	101.2400	104.93	99.9900
2	0.525	0.1360	1.5200	1.0000	8.130	26787.7	0.120	1.118	1.619	112.85	102.0700	103.09	100.3700
							0.170	1.116	1.620	112.82	101.9000	103.07	100.2000
							0.270	1.092	1.670	112.81	101.8000	103.06	100.2000
3	0.525	0.1200	1.7400	1.0800	7.308	22295.7	0.120	1.103	1.705	114.41	101.4800	104.05	100.0100
							0.170	1.100	1.700	114.38	101.2000	104.02	100.2000
							0.270	1.074	1.720	114.36	101.2500	104.00	100.1700
4	0.500	0.0900	1.9200	1.1500	5.879	16698.6	0.095	1.037	1.797	116.45	101.0500	106.62	99.9000
							0.145	1.032	1.800	116.42	101.1300	106.59	99.9900
							0.245	1.008	1.810	116.40	101.1600	106.57	100.1000
5	0.500	0.1250	1.5100	1.0000	7.271	25955.3	0.095	1.086	1.630	114.63	101.3000	104.05	99.9300
							0.145	1.080	1.680	114.60	101.7000	104.01	99.9000
							0.245	1.056	1.730	114.58	101.5000	104.00	99.9700
6	0.500	0.1050	1.2800	1.0800	6.520	19893.4	0.095	1.068	1.742	115.66	101.1800	105.37	99.9500
							0.145	1.056	1.740	115.62	101.2000	105.34	100.0000
							0.245	1.044	1.750	115.63	101.1000	105.32	100.1000

Table (5.10) - Wave Characteristics and Wave Force Coefficients for Structure (F) with 1010 Grs at the Top.

STR. DIAMETER=0.0574 M STR. LENGTH=0.8350 M
NATURAL FREQ. IN AIR=11.520 DAMPING COEFF. IN AIR=0.070

TEST NO.	WATER DEPTH METER	WAVE HEIGHT METER	WAVE LENGTH METER	WAVE PERIOD SEC.	KEULEGAN-CARPENTER NUMBER	REYNOLDS NUMBER	LEVEL FROM SWL	CD FIXED	CM FIXED	CD FR %----- CD FX	CD FR(R) %-------- CD FR	CM FR %----- CM FX	CM FR(R) %-------- CM FR
1	0.525	0.1080	1.9600	1.1500	6.742	19148.2	0.120	1.080	1.719	115.98	101.2800	105.53	99.8000
							0.170	1.074	1.721	115.95	101.3000	105.49	99.9500
							0.270	1.062	1.737	115.96	101.2000	105.47	100.0100
2	0.525	0.1360	1.5200	1.0000	8.130	26787.7	0.120	1.128	1.680	113.46	101.9800	103.62	100.3200
							0.170	1.126	1.660	113.41	101.8000	103.59	100.2000
							0.270	1.120	1.670	113.40	101.5000	103.57	100.1700
3	0.525	0.1200	1.7400	1.0800	7.308	22295.7	0.120	1.092	1.716	115.07	101.4500	104.61	100.1100
							0.170	1.080	1.720	115.05	101.3700	104.58	100.2000
							0.270	1.074	1.737	115.03	101.3000	104.56	100.2200
4	0.500	0.0900	1.9200	1.1500	5.879	16698.6	0.095	1.038	1.790	117.08	101.0300	107.24	99.8700
							0.145	1.032	1.801	117.05	101.1000	107.21	99.9000
							0.245	1.008	1.812	117.04	101.0200	107.19	100.0200
5	0.500	0.1250	1.5100	1.0000	7.271	25955.3	0.095	1.074	1.726	115.22	101.2000	104.67	100.1500
							0.145	1.068	1.730	115.19	101.2000	104.63	100.2000
							0.245	1.056	1.739	115.17	101.1000	104.64	100.1500
6	0.500	0.1050	1.2800	1.0300	6.520	19893.4	0.095	1.055	1.753	116.28	101.2100	105.91	99.9600
							0.145	1.056	1.755	116.27	101.1100	105.88	99.9900
							0.245	1.050	1.760	116.25	101.0900	105.86	100.0500

Table (5.11) - Wave Characteristics and Wave Force Coefficients for Structure (F) with 2020 Grs at the Top.

depth while the structure velocity decreases sinusoidally with depth, for the case of the drag coefficient and similarly for the acceleration in the case of the inertia coefficient.

The Reynolds numbers of the tests varied from (1.09×10^4 to 4.7×10^4). Figures (5.3.1a) and (5.3.1b) show the variation of C_D and C_M with Reynolds numbers, 1/D.

The relationship between the waves and the corresponding forces can be shown by plotting the results of the dimensionless parameter (Keulegan-Carpenter number) and the coefficients representing the forces.

Figure (5.3.2a) and Figure (5.3.2b) show the variation of the drag coefficients and the variation of the inertia coefficients for the fixed structure versus the Keulegan-Carpenter numbers respectively.

Figures (5.3.3a through 5.3.12a) show the variation of the percentage ratio of C_D for free structures to that for fixed structures for each structure tested and Figure (5.3.13) shows the same variation of the ratio for all the tested structures with Keulegan-Carpenter numbers. Figures (5.3.3b through 5.3.12b) show the variation of the percentage ratio of the

C_M for free structures to that for fixed structures for each structure tested and Figure (5.3.14) shows the same variation of the ratio for all tested structures with Keulegan-Carpenter number.

Figure (5.3.15a) and Figure (5.3.15b) show the variation of the percentage ratio of the C_D calculated by the modified Morison's equation to that calculated by Morison's equation and similarly for C_M with Keulegan-Carpenter numbers respectively.

The data presented in the figures cited above exhibit certain characteristics which will be discussed below.

(a) THE VARIATION OF THE FORCE COEFFICIENTS (C_D AND C_M) FOR FIXED STRUCTURE WITH KEULEGAN-CARPENTER NUMBER

The drag coefficient C_D increases with increasing K-C numbers used in the tests (the range of K-C number used was (1.874 - 10.604)), while the C_M decreases with increasing K-C number for the same range. These results agree with the previous results of Keulegan and Carpenter 1968 [27]. The results show less scattering and that is because of the method used in measuring the local force for a segment of the structure which minimizes the error; arising due to the

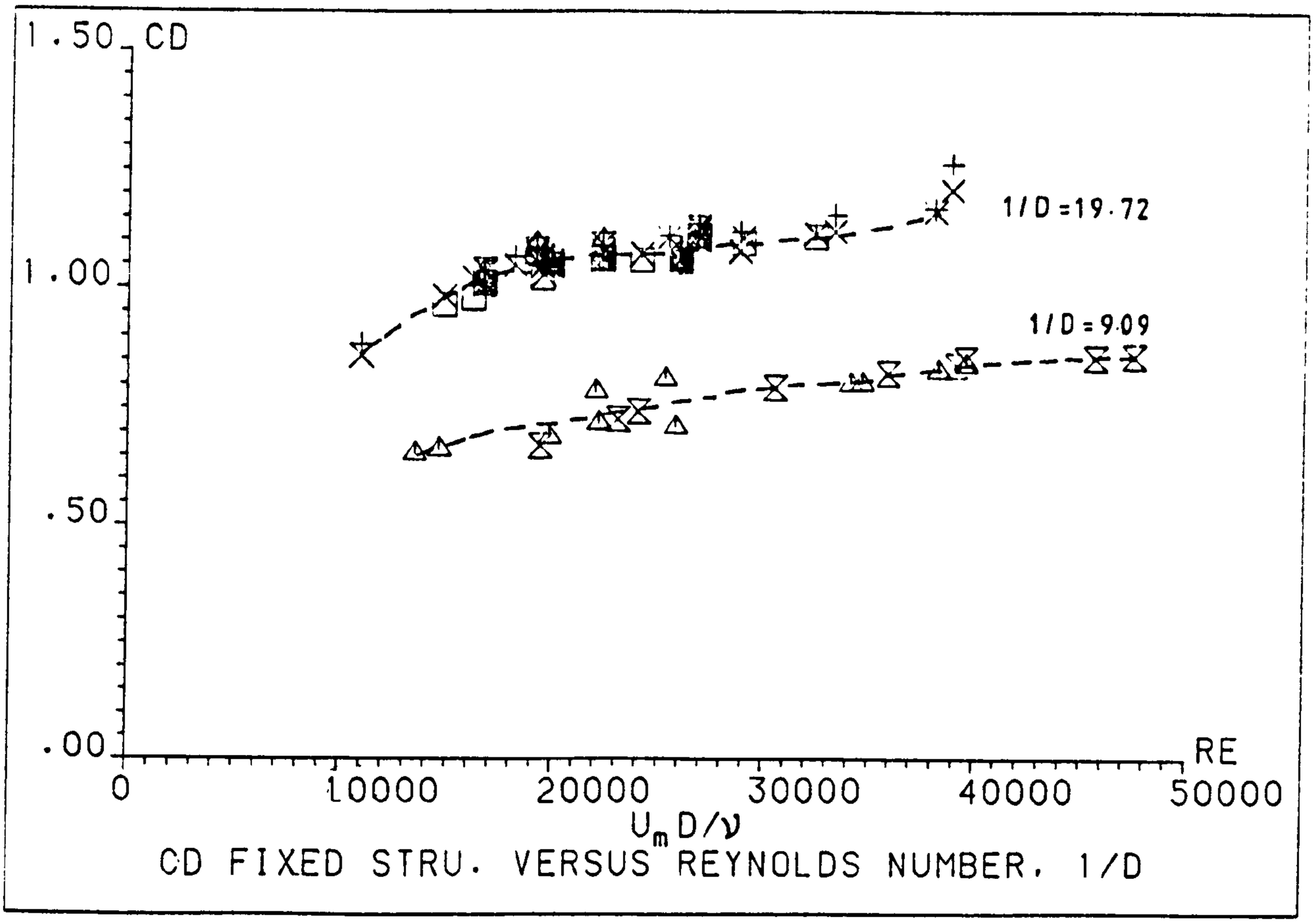

Figure (5.3.1a) - Variation of C_D with Reynolds Number, 1/D.

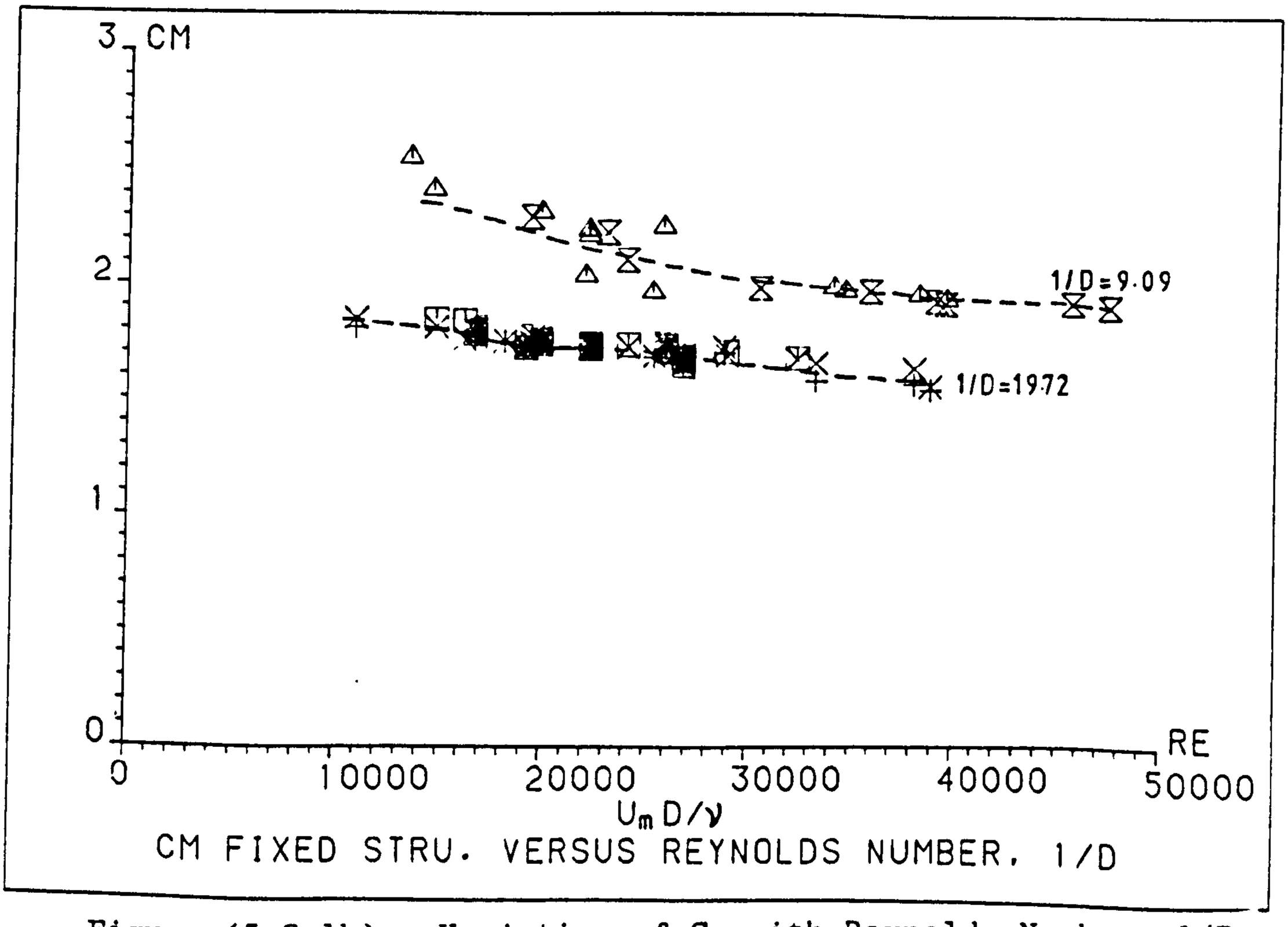

Figure (5.3.1b) - Variation of C_M with Reynolds Number, 1/D.

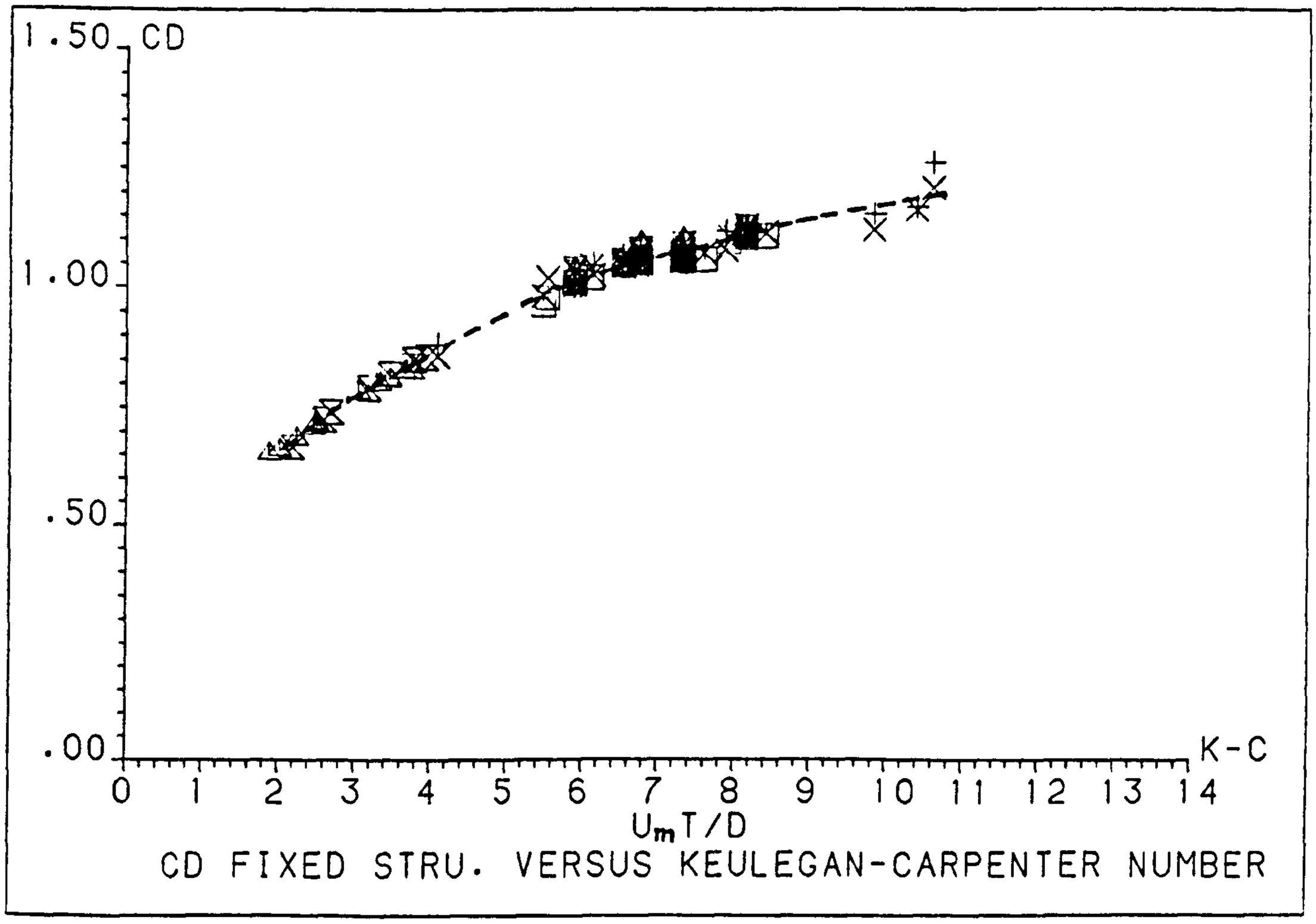

Figure (5.3.2a) - Variation of C_D with Keulegan-Carpenter Number.

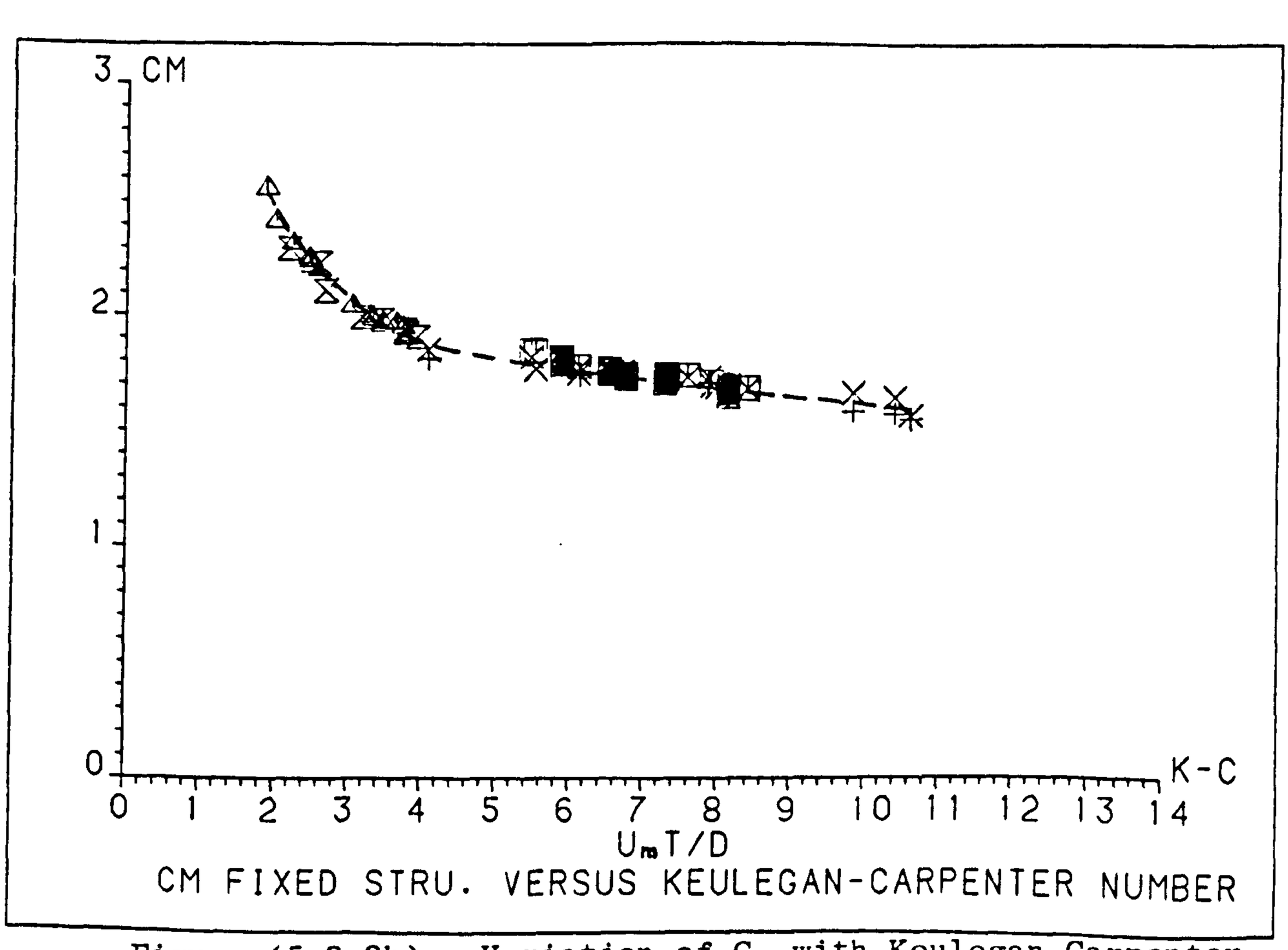

Figure (5.3.2b) - Variation of C_M with Keulegan-Carpenter Number.

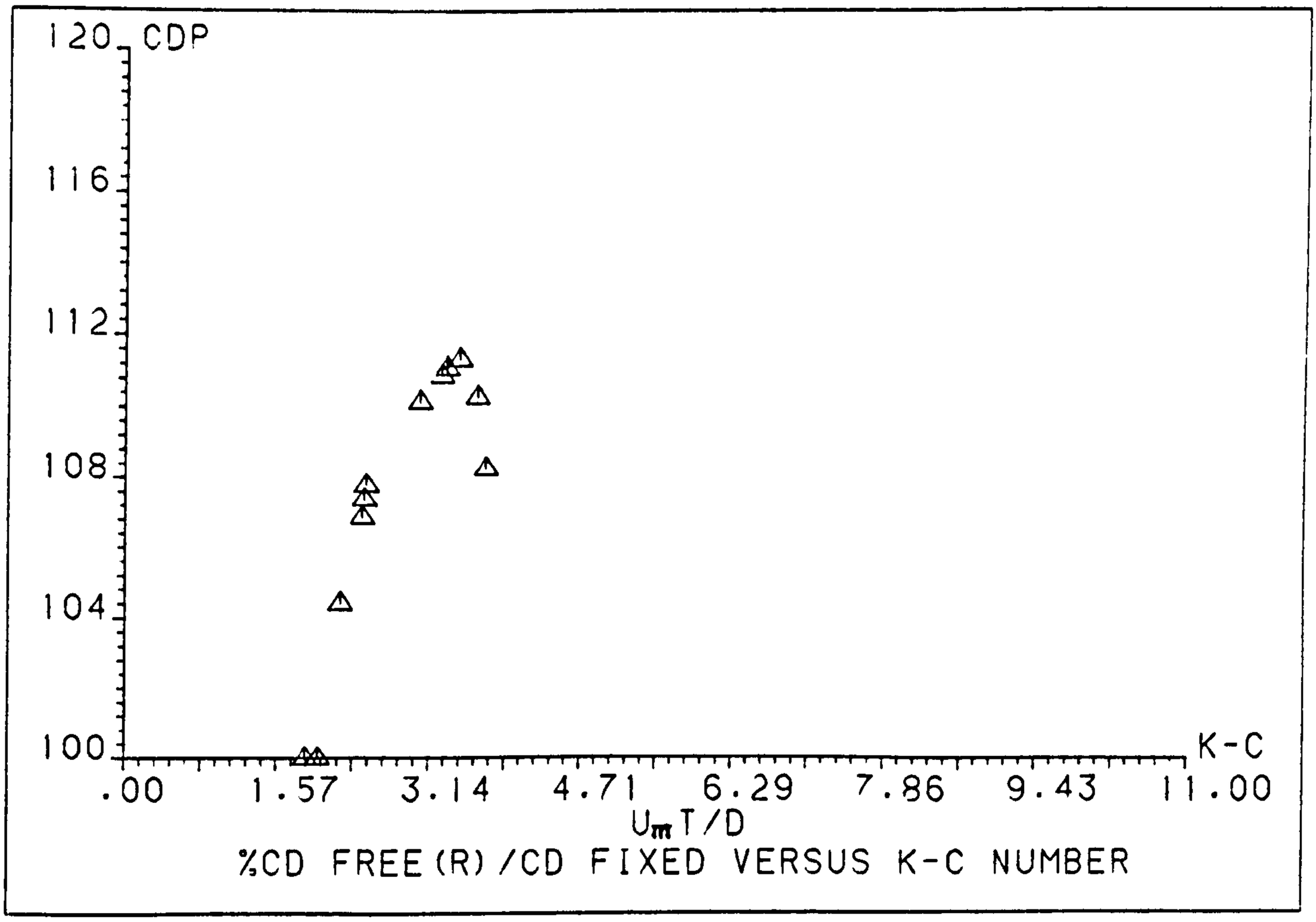

Figure (5.3.3a) - Variation of % (C_D Free (R)/C_D Fixed) with K-C Number for Structure A.

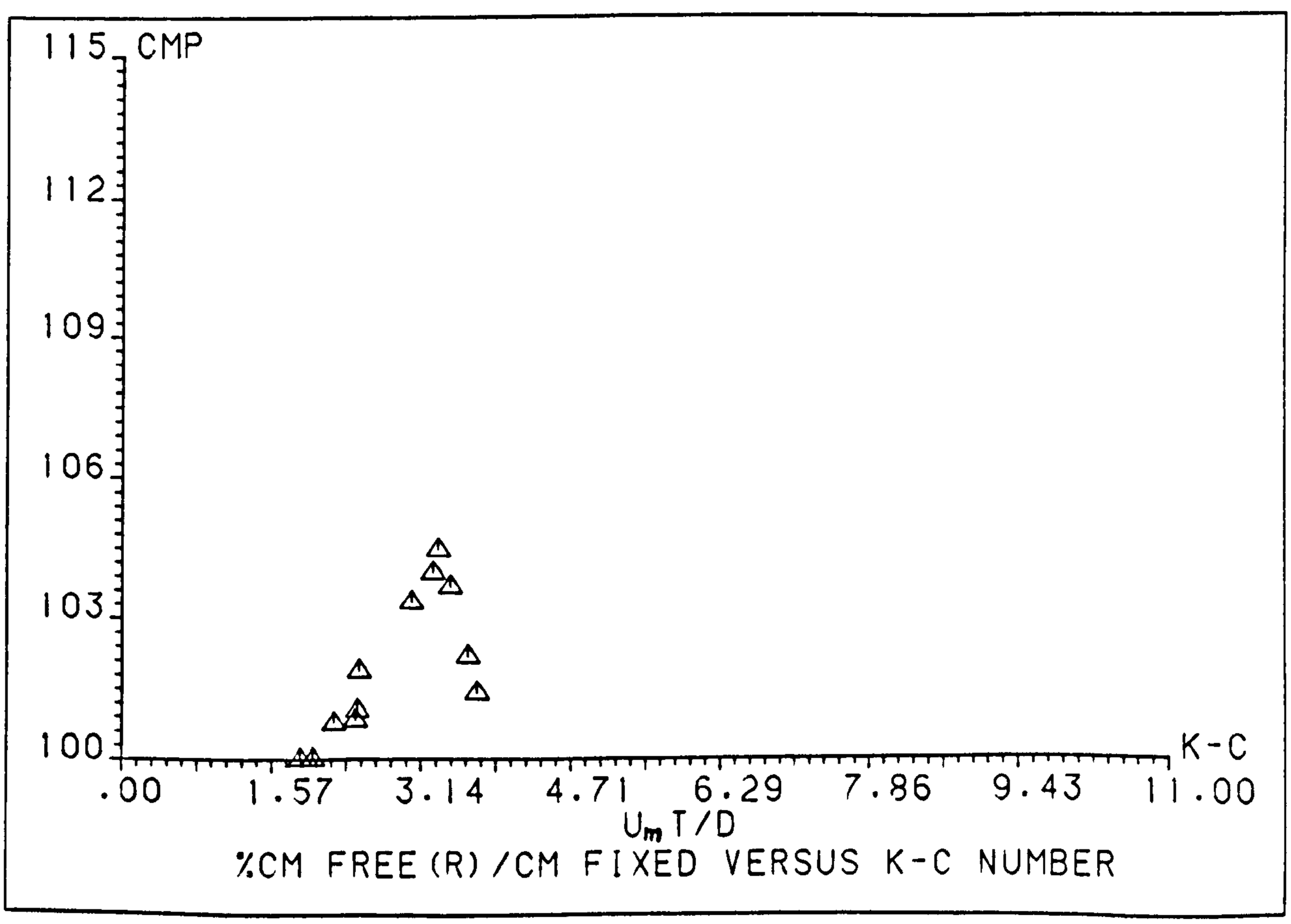

Figure (5.3.3b) - Variation of % (C_M Free (R)/C_M Fixed) with K-C Number for Structure A.

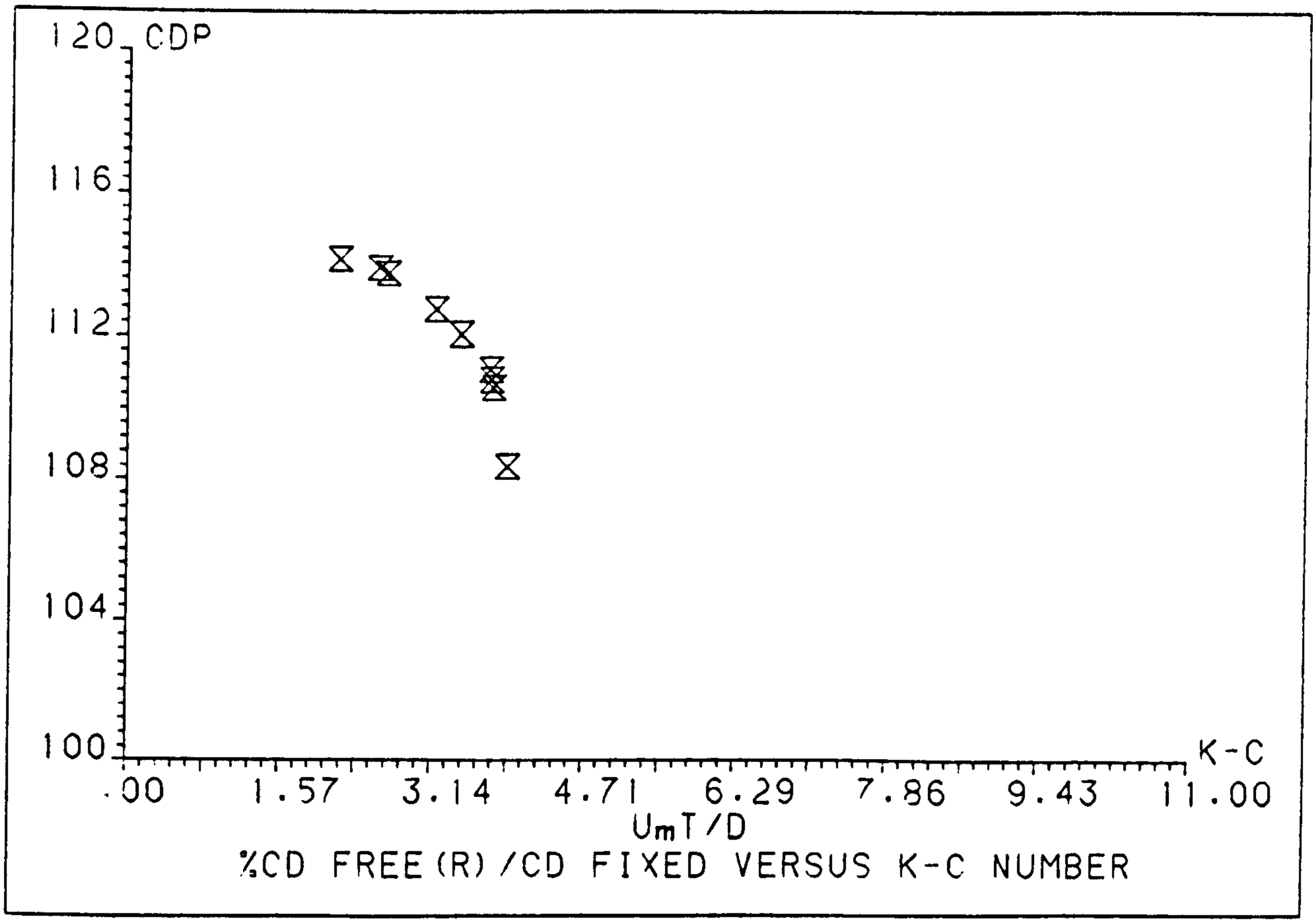

Figure (5.3.4a) - Variation of % (C_D Free (R)/C_D Fixed) with K-C Number for Structure B.

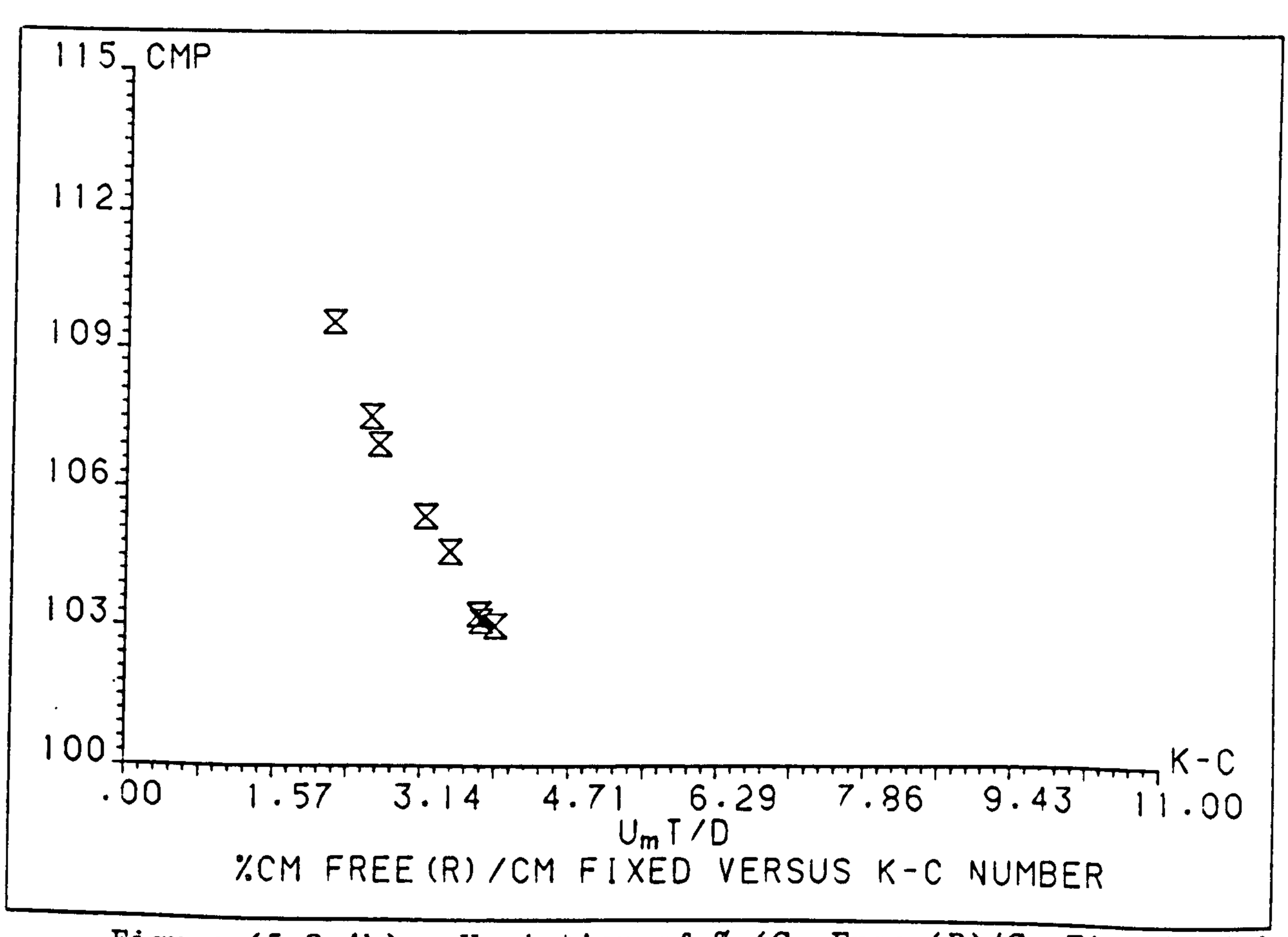

Figure (5.3.4b) - Variation of % (C_M Free (R)/C_M Fixed) with K-C Number for Structure B.

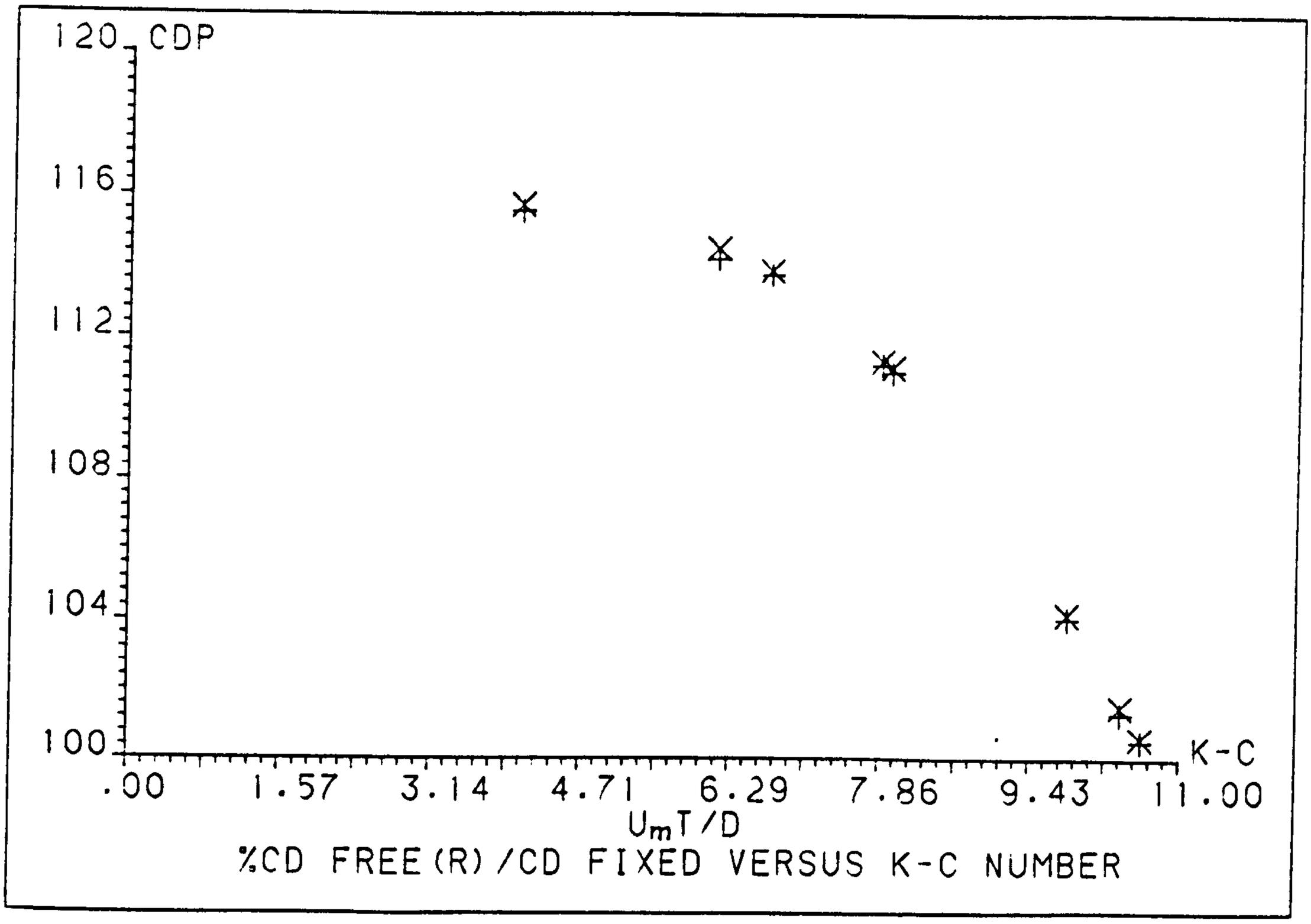

Figure (5.3.5a) - Variation of % (C_D Free (R)/C_D Fixed) with K-C Number for Structure C.

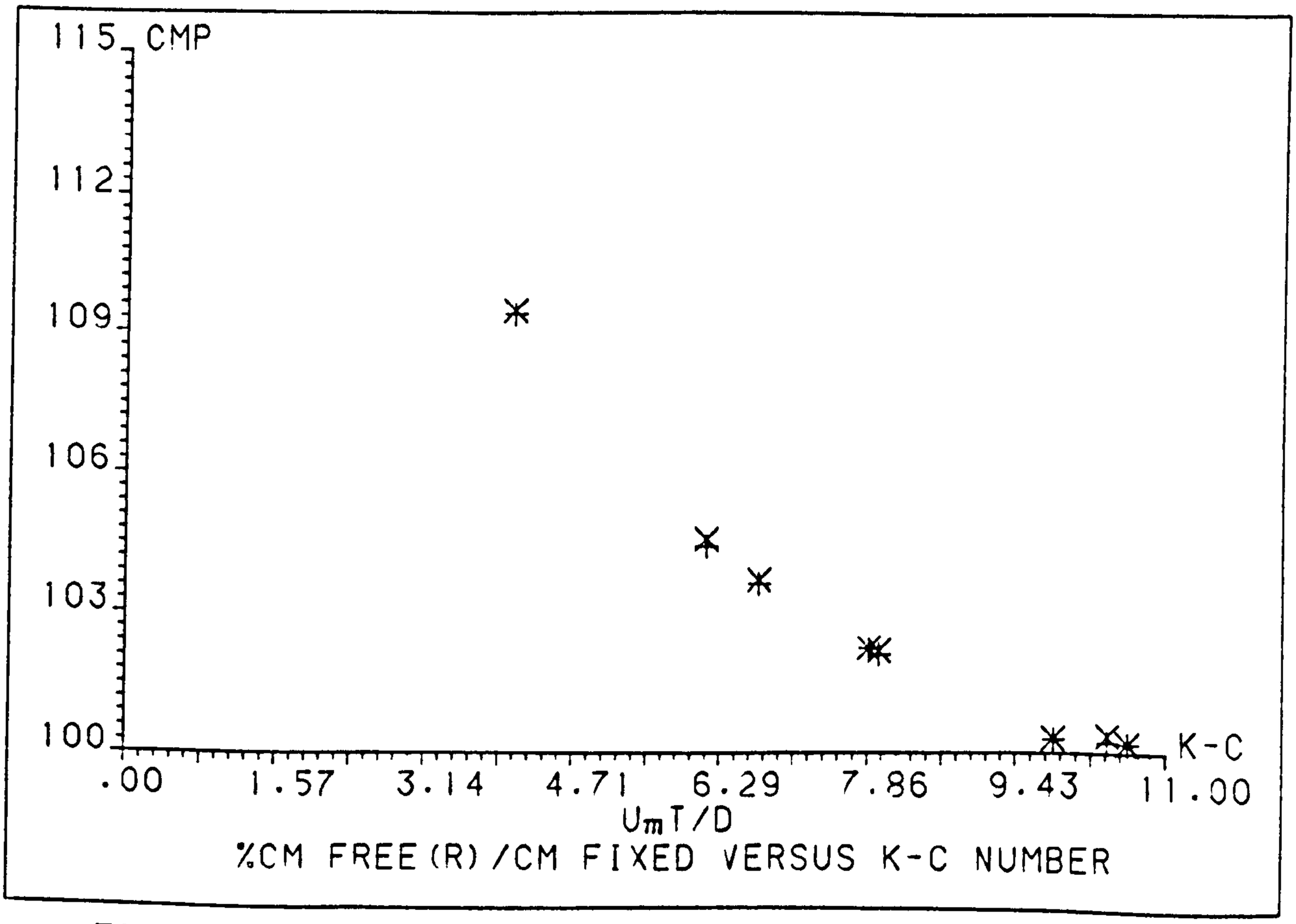

Figure (5.3.5b) - Variation of % (C_M Free (R)/C_M Fixed) with K-C Number for Structure C.

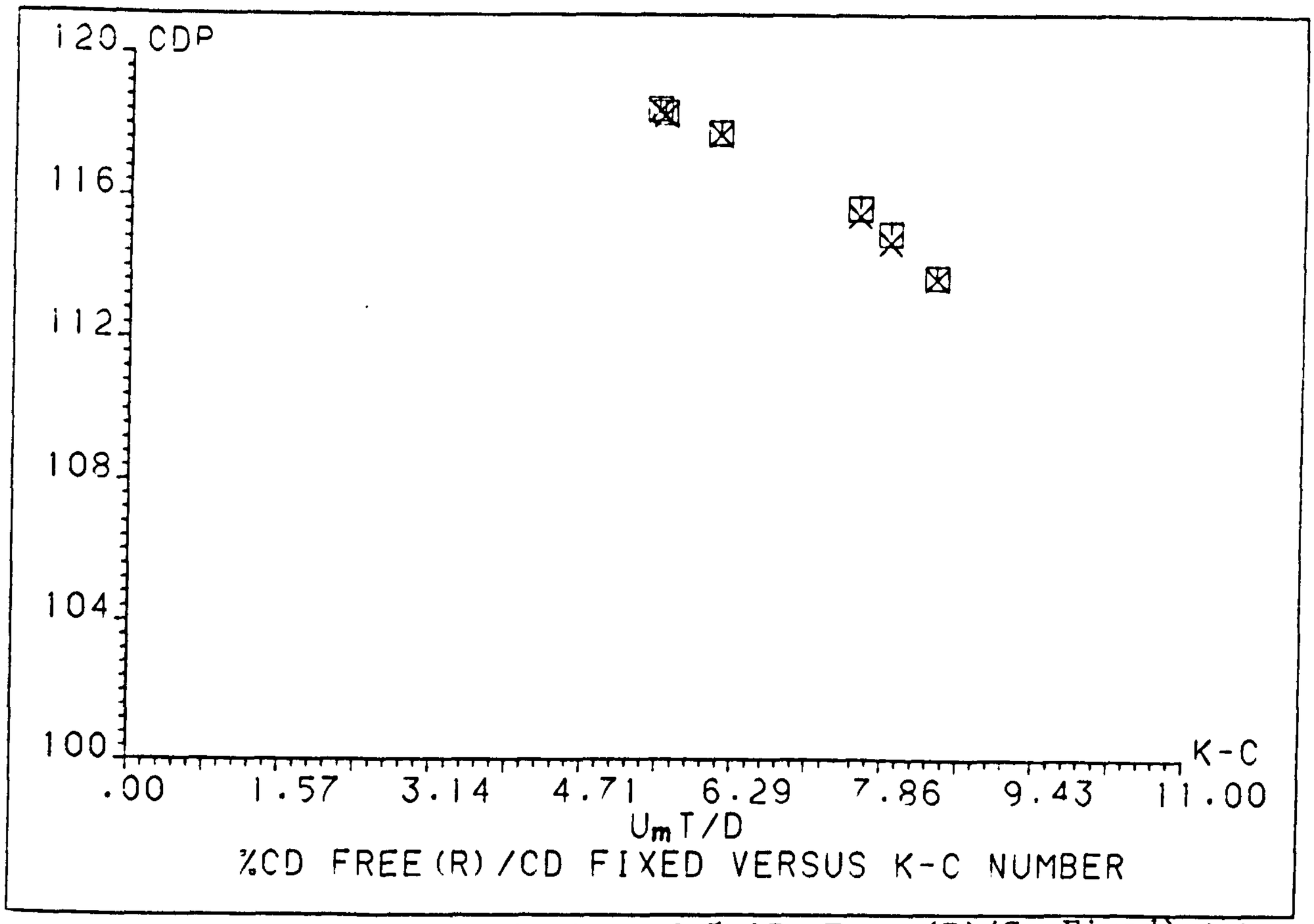

Figure (5.3.6a) - Variation of % (C_D Free (R)/C_D Fixed)with K-C Number for Structure D.

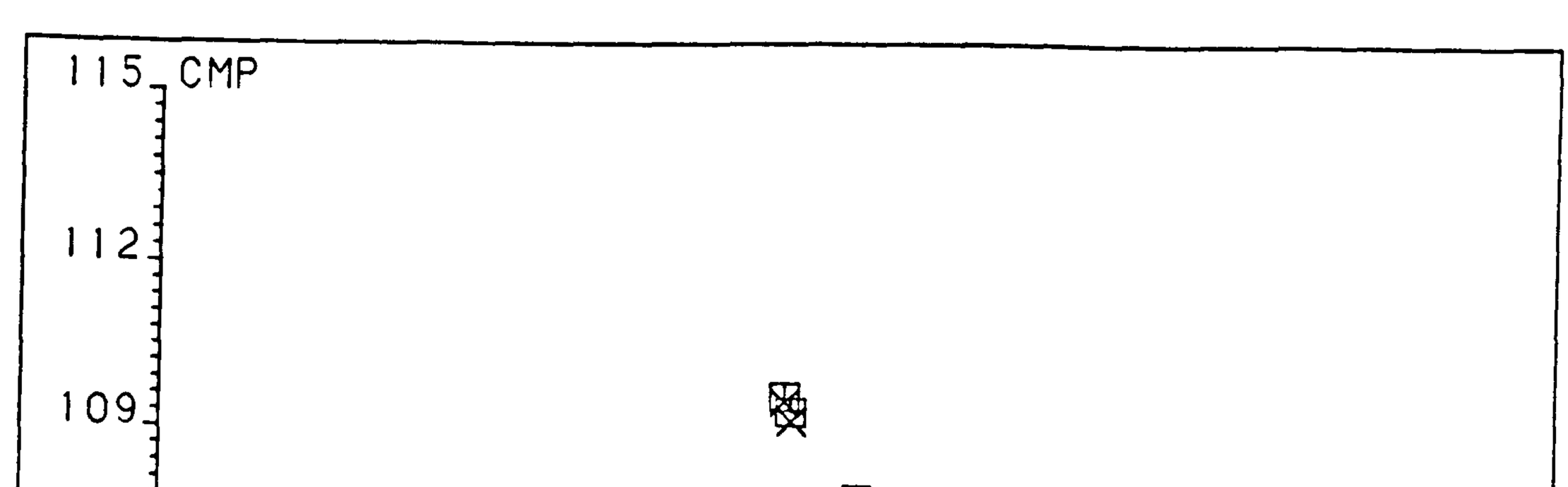

Figure (5.3.6b) - Variation of % (C_M Free (R)/C_M Fixed) with K-C Number for Structure D.

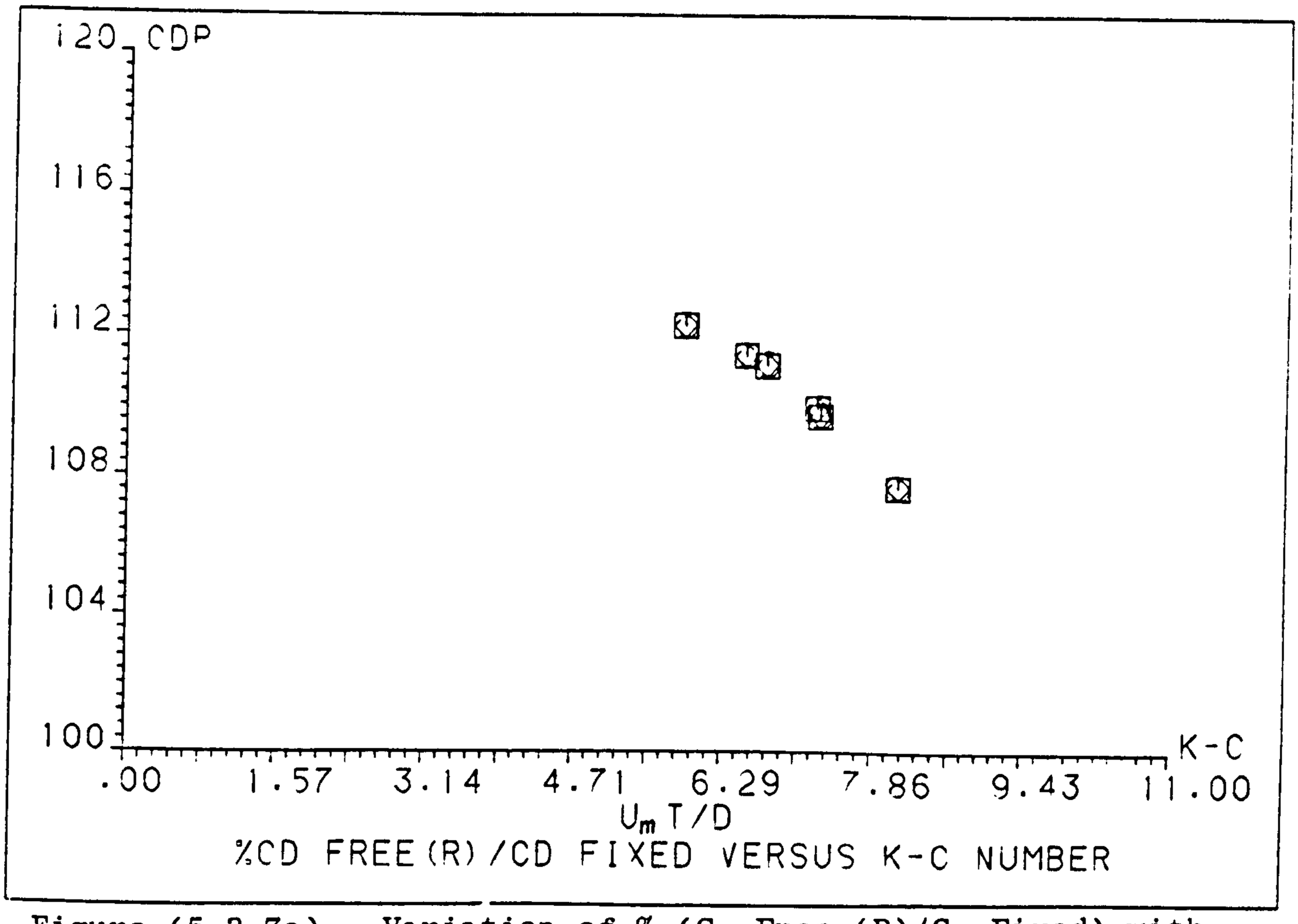

Figure (5.3.7a) - Variation of % (C_D Free (R)/C_D Fixed) with K-C Number for Structure E with no Mass at Top.

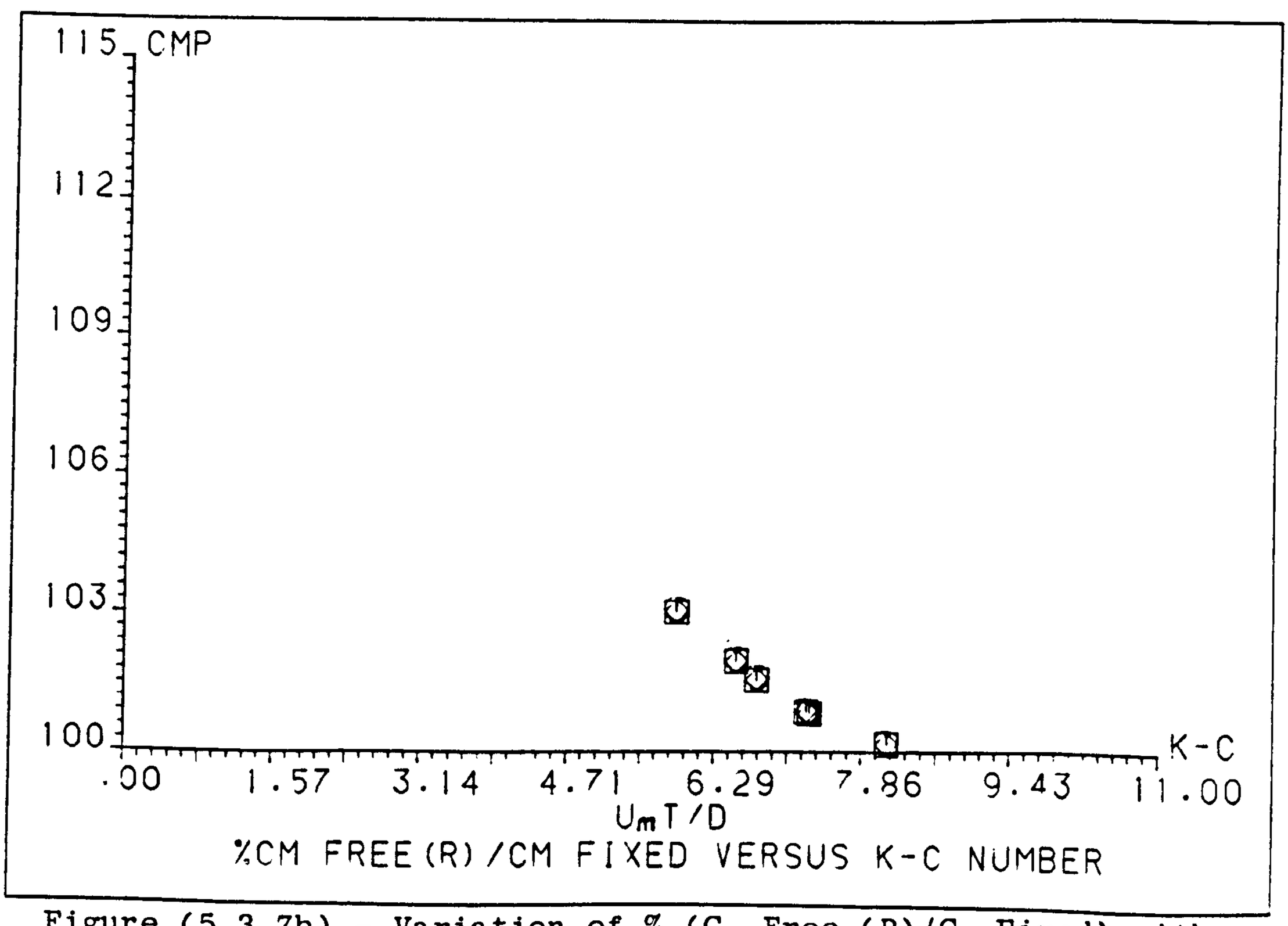

Figure (5.3.7b) - Variation of % (C_M Free (R)/C_M Fixed) with K-C Number for Structure E with no Mass at Top.

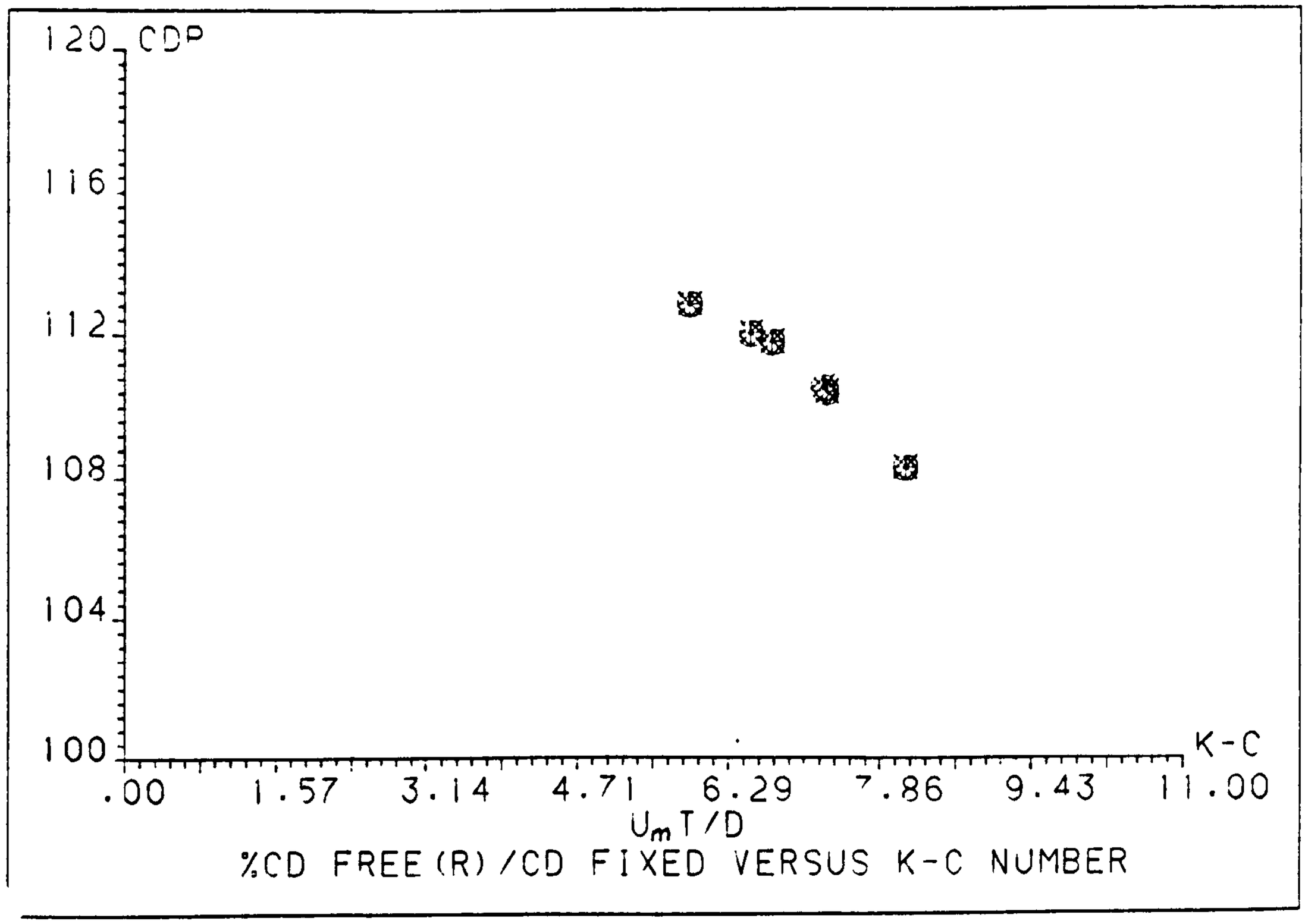

Figure (5.3.8a) - Variation of % (C_D Free (R)/C_D Fixed) with K-C Number for Structure E with 1010 Grs at Top.

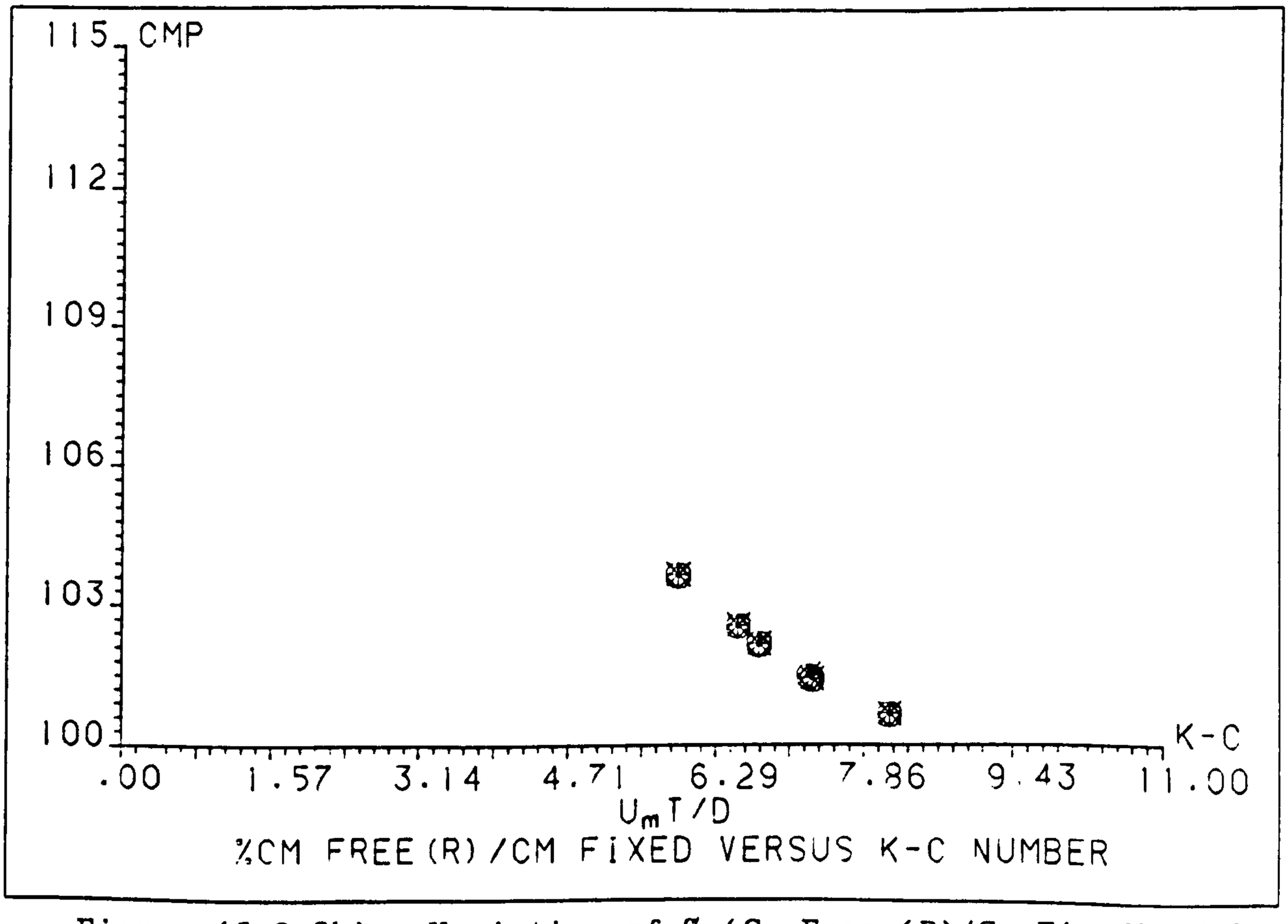

Figure (5.3.8b) - Variation of % (C_M Free (R)/C_M Fixed) with K-C Number for Structure E with 1010 Grs at Top.

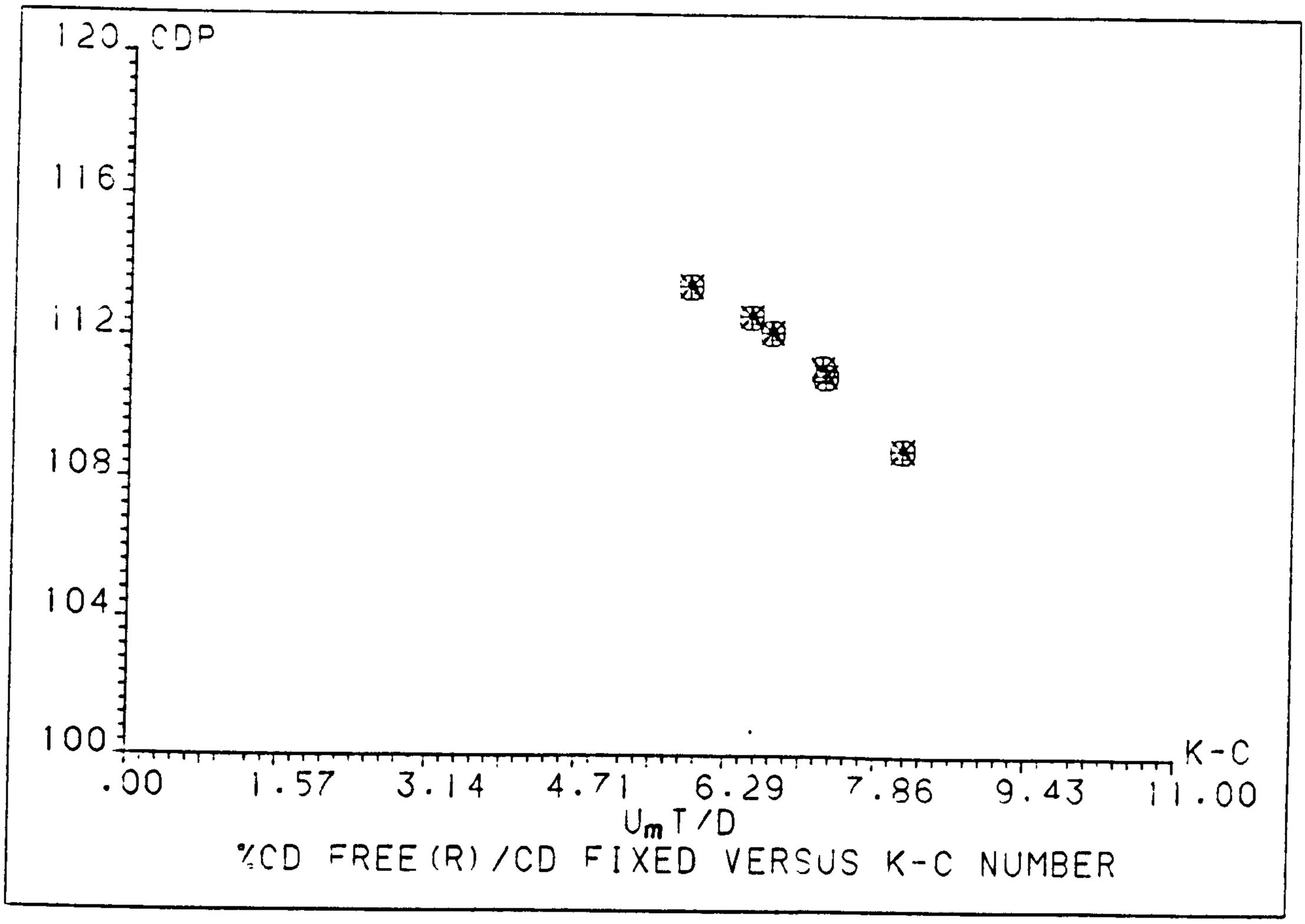

Figure (5.3.9a) - Variation of % (C_D Free (R)/C_D Fixed) with K-C Number for Structure E with 2020 Grs at Top.

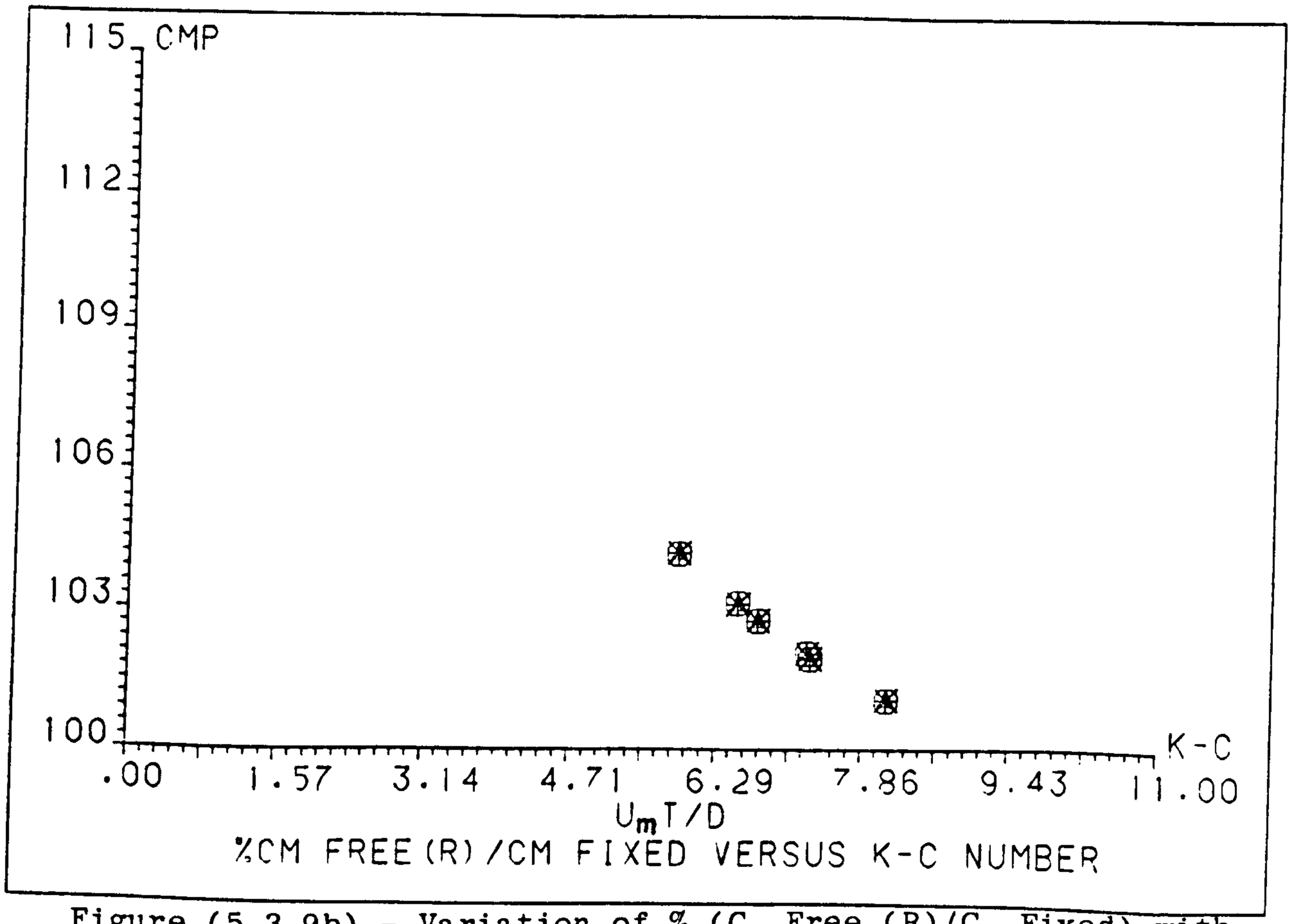

Figure (5.3.9b) - Variation of % (C_M Free (R)/C_M Fixed) with K-C Number for Structure E with 2020 Grs at Top.

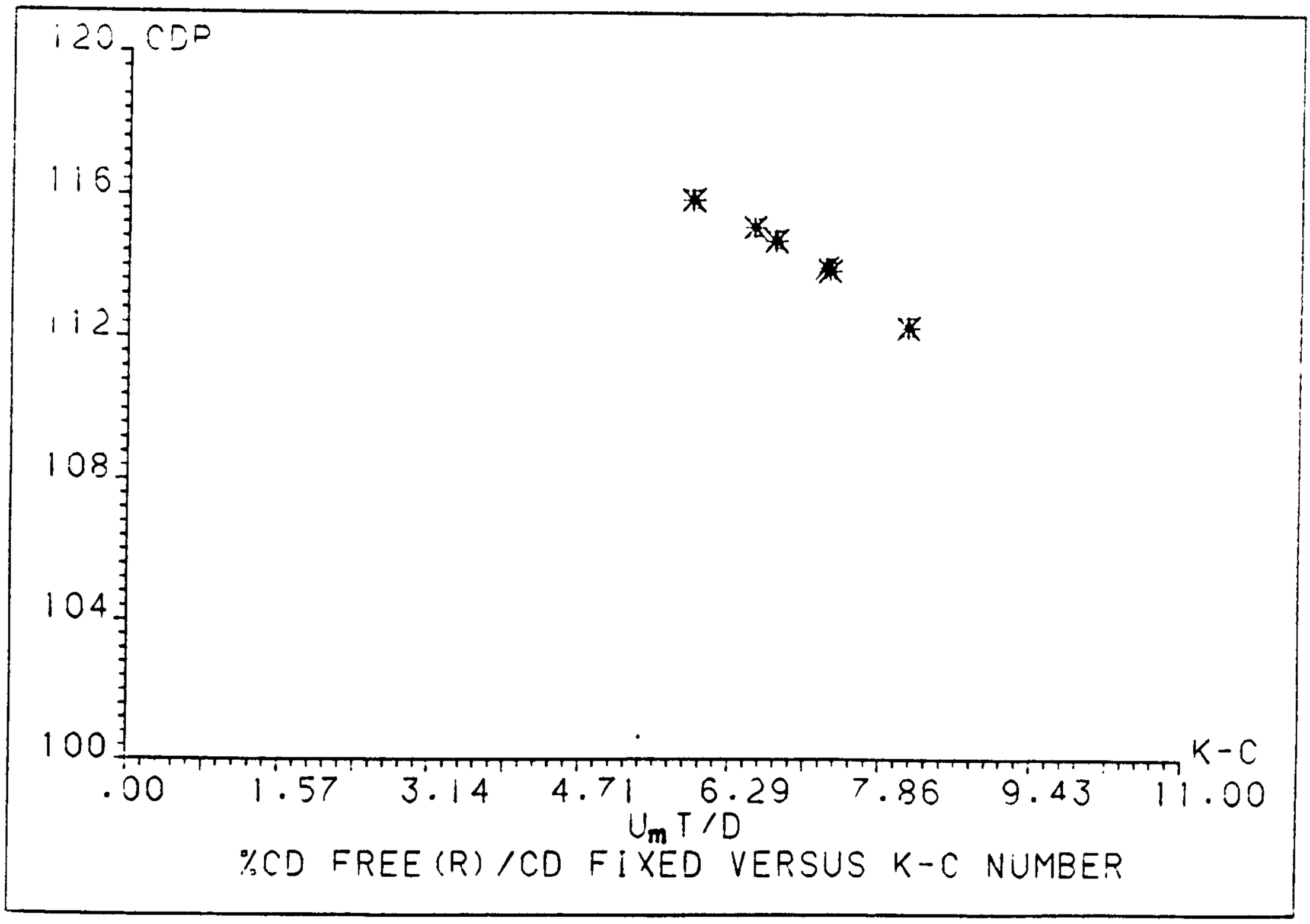

Figure (5.3.10a) - Variation of % (C_D Free (R)/C_D Fixed) with K-C Number for Structure F with no Mass at Top.

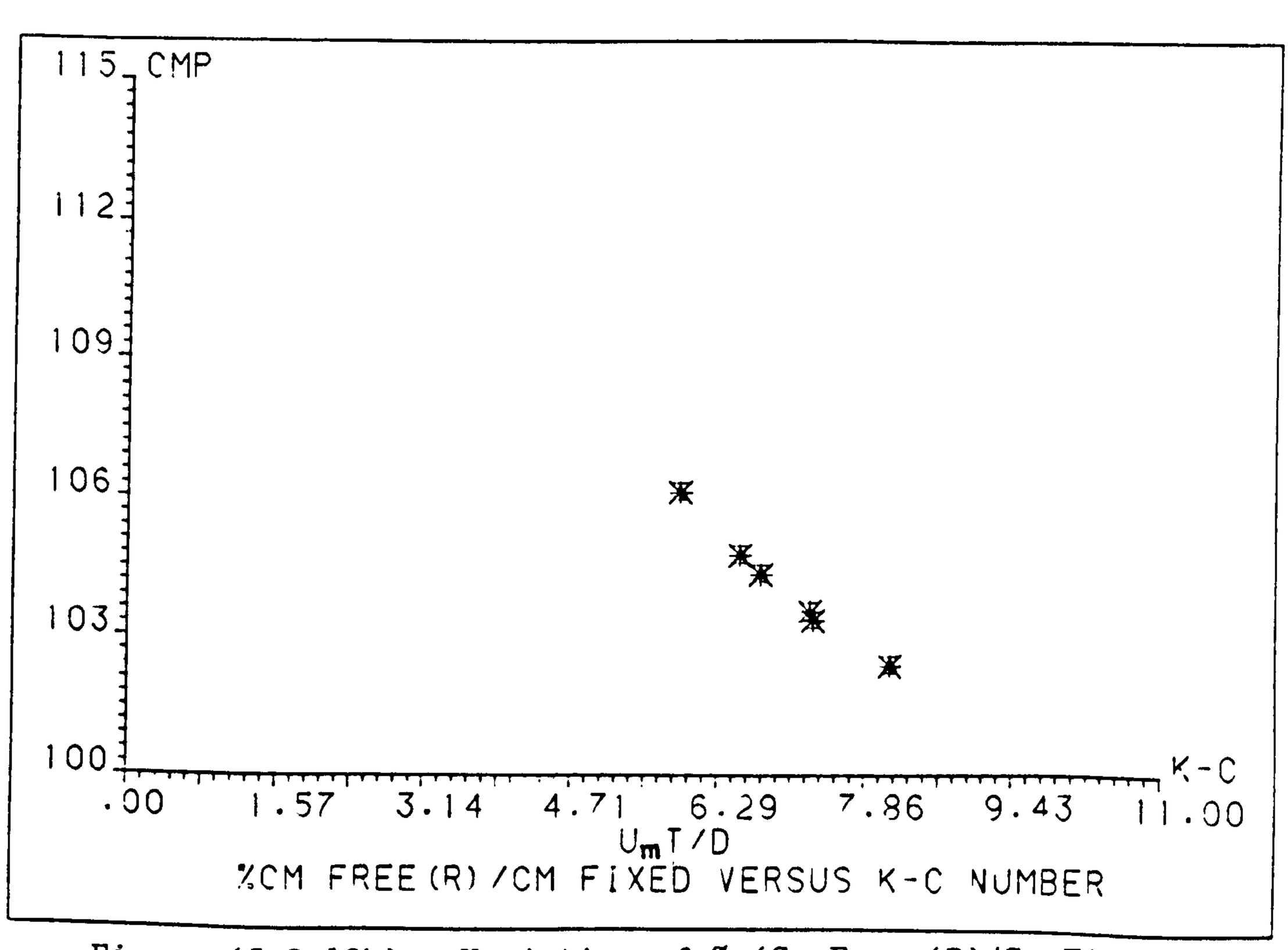

Figure (5.3.10b) - Variation of % (C_D Free (R)/C_D Fixed) with K-C Number for Structure F with no Mass at Top.

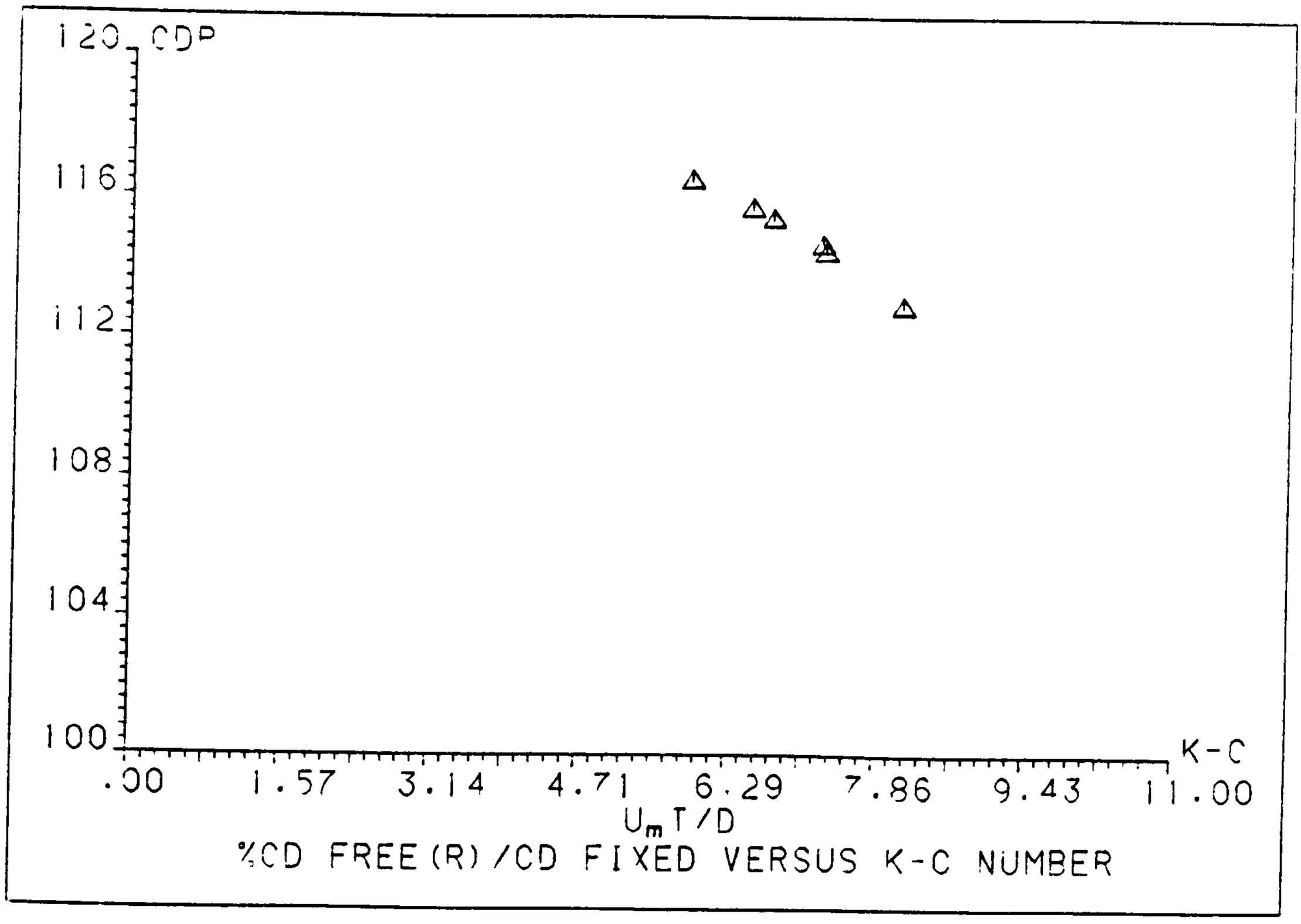

Figure (5.3.11a) - Variation of % (C_D Free (R)/C_D Fixed) with K-C Number for Structure F with 1010 Grs at Top.

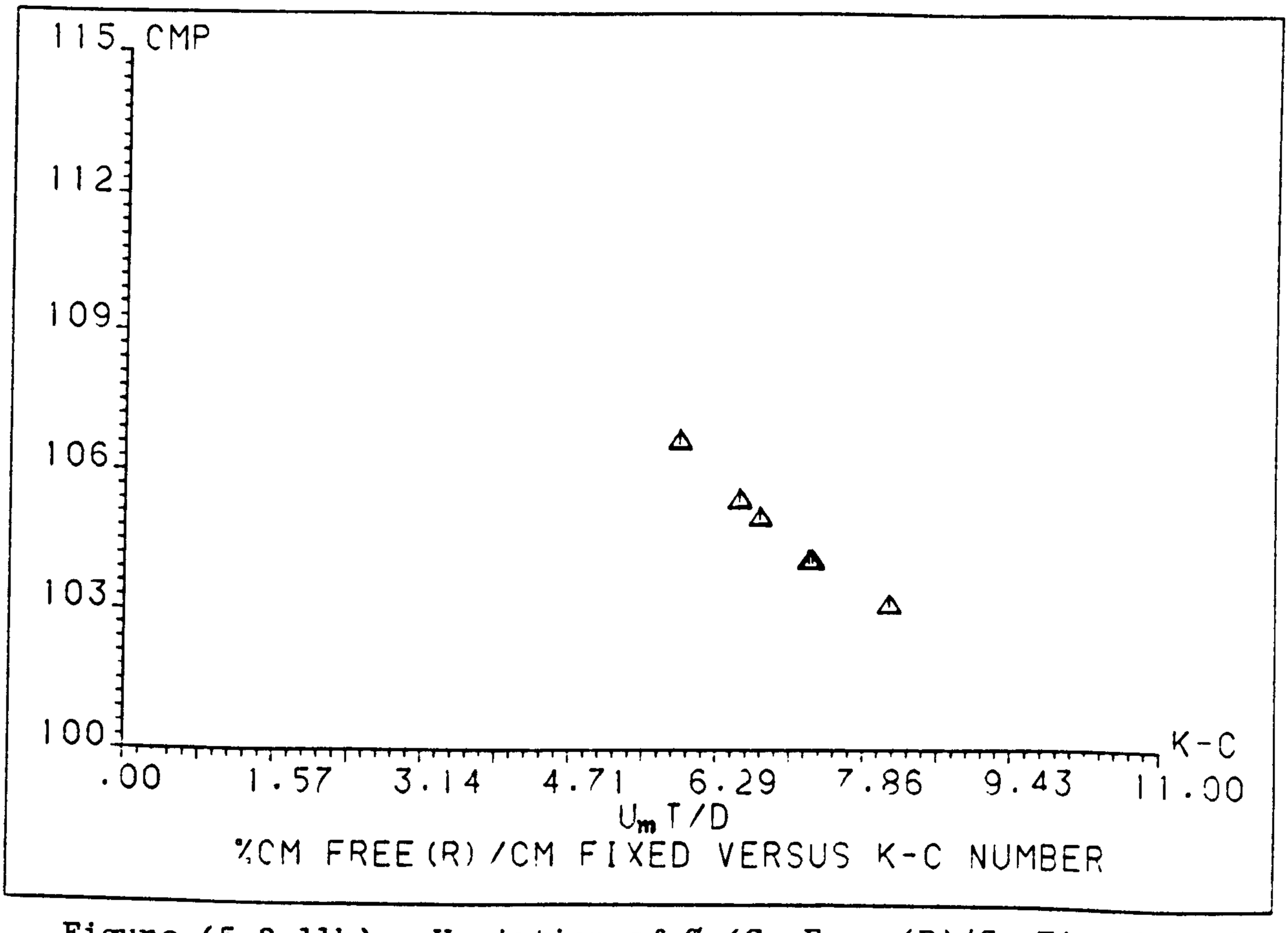

Figure (5.3.11b) - Variation of % (C_M Free (R)/C_M Fixed) with K-C Number for Structure F with 1010 Grs at Top.

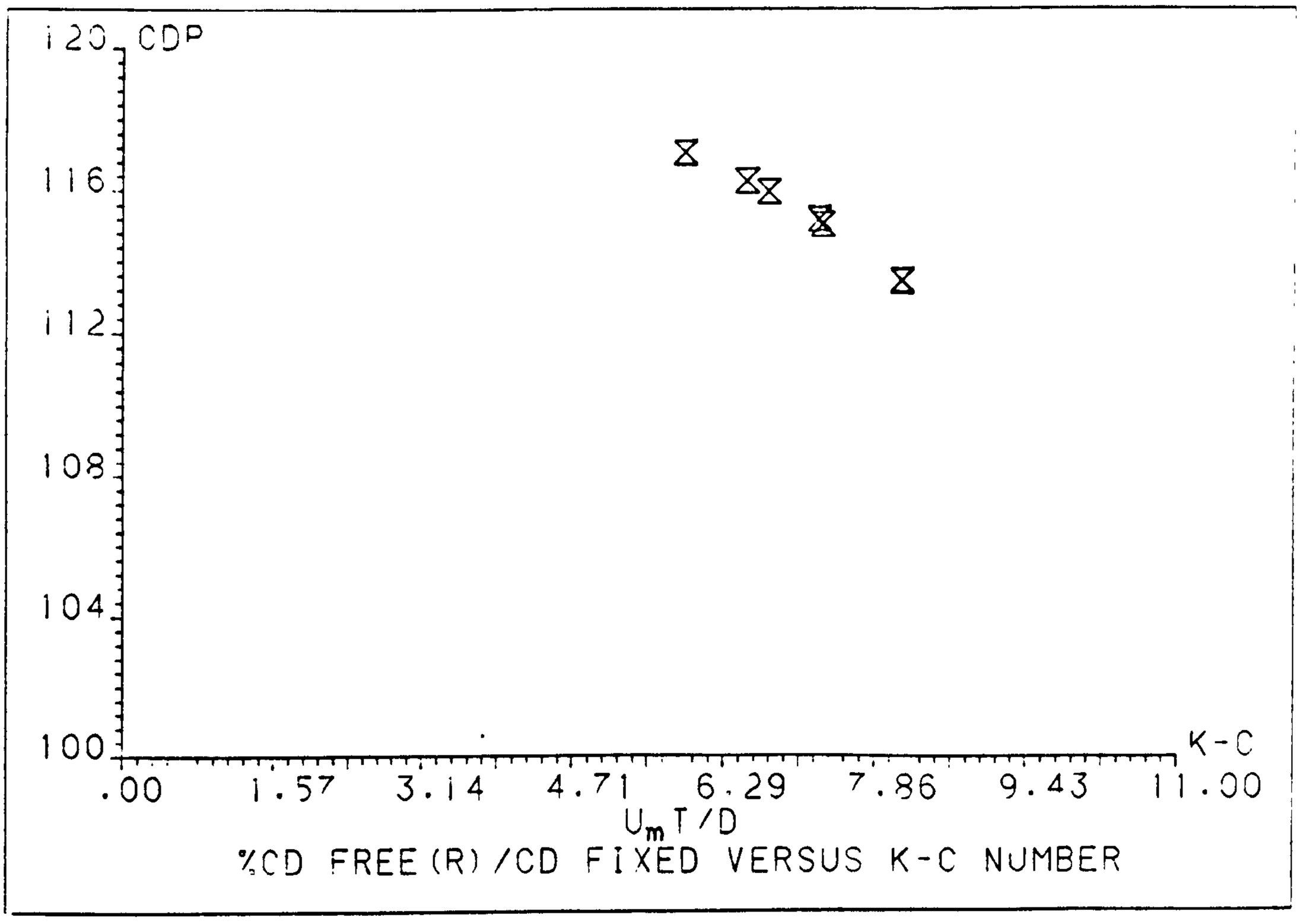

Figure (5.3.12a) - Variation of % (C_D Free (R)/C_D Fixed) with K-C Number for Structure F with 2020 Grs at Top.

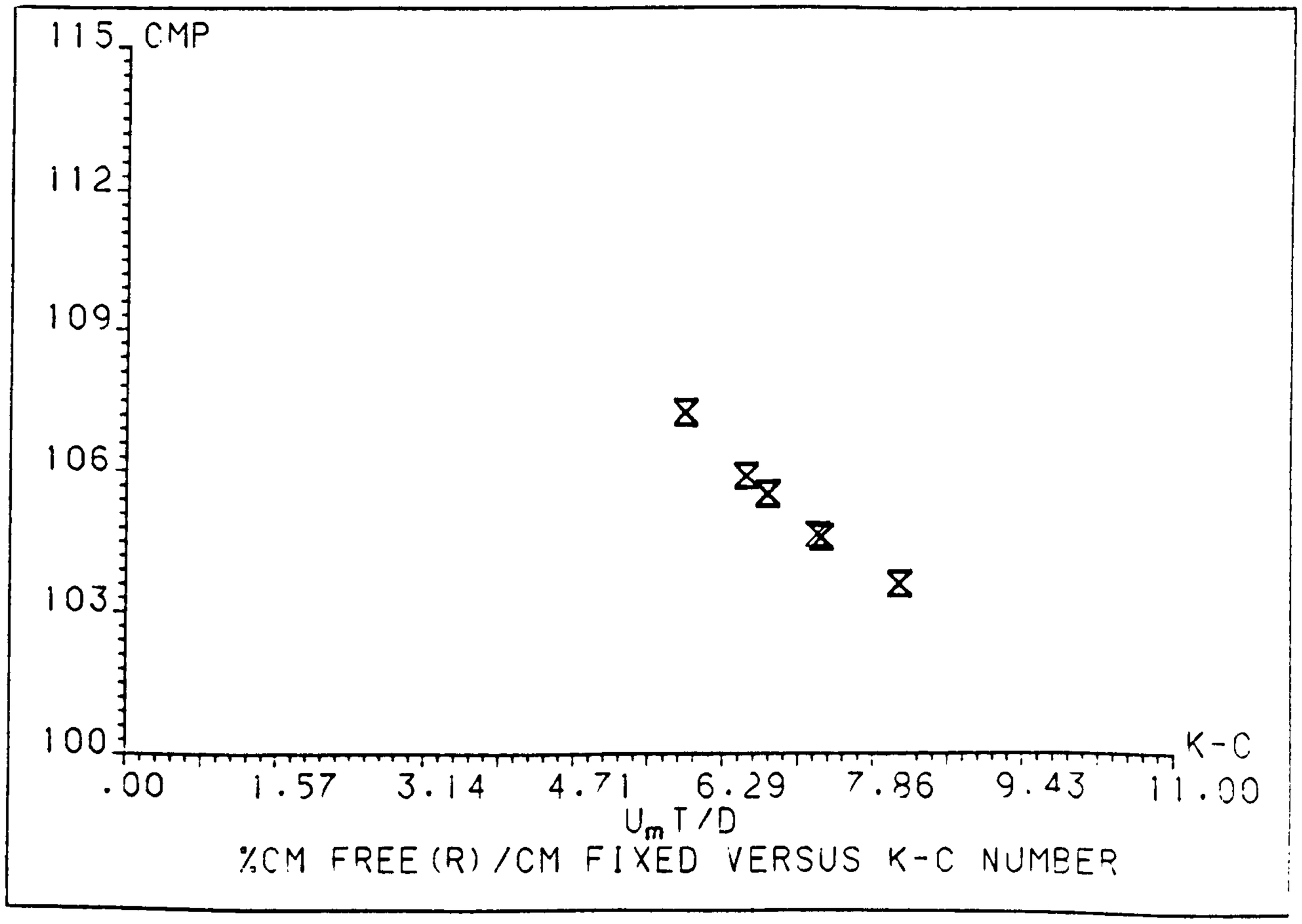

Figure (5.3.12b) - Variation of % (C_M Free (R)/C_M Fixed)with K-C Number for Structure F with 2020 Grs at Top.

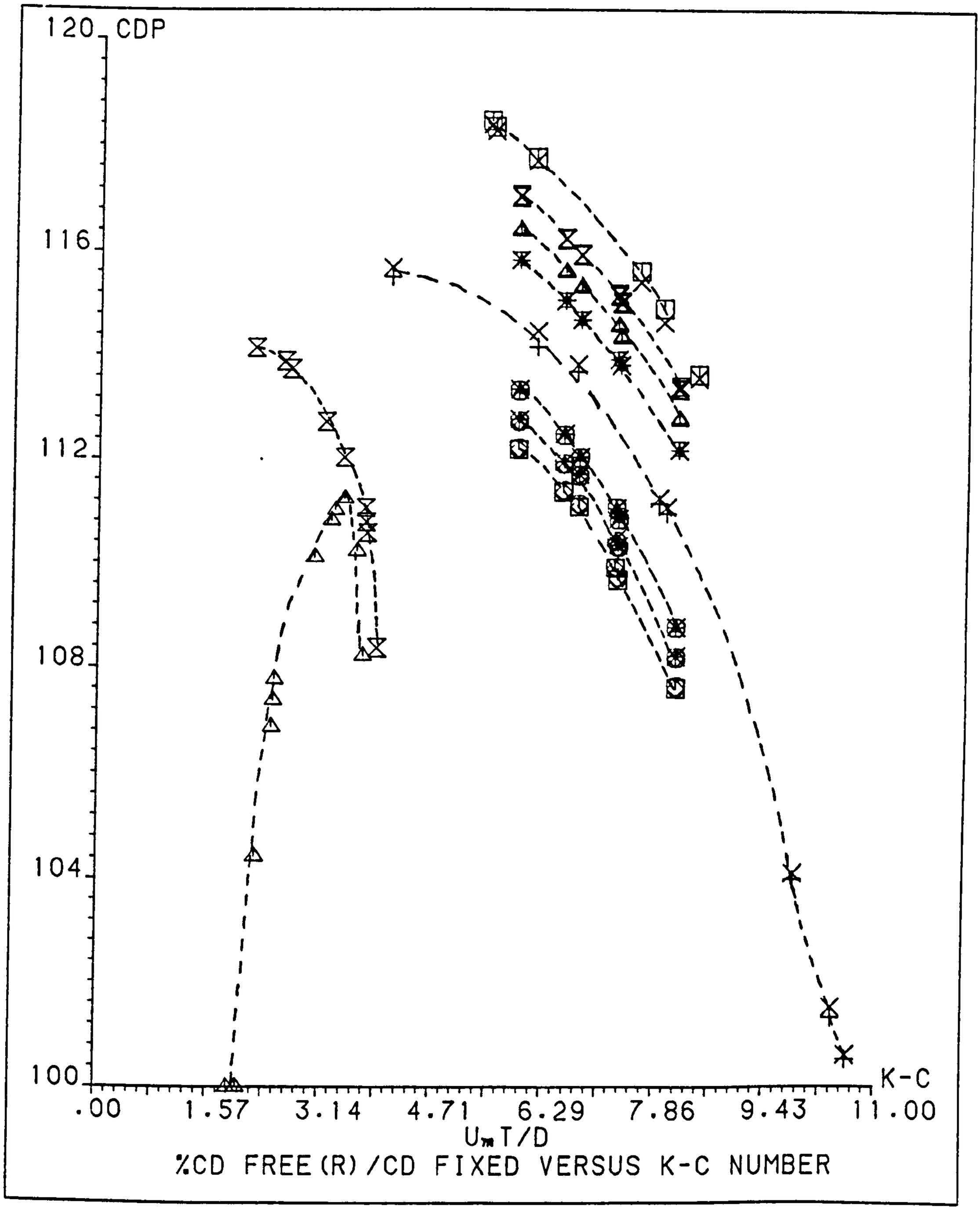

Figure (5.3.13) - Variation of % (C_D Free (R)/C_D Fixed) with K-C Number for all Structures.

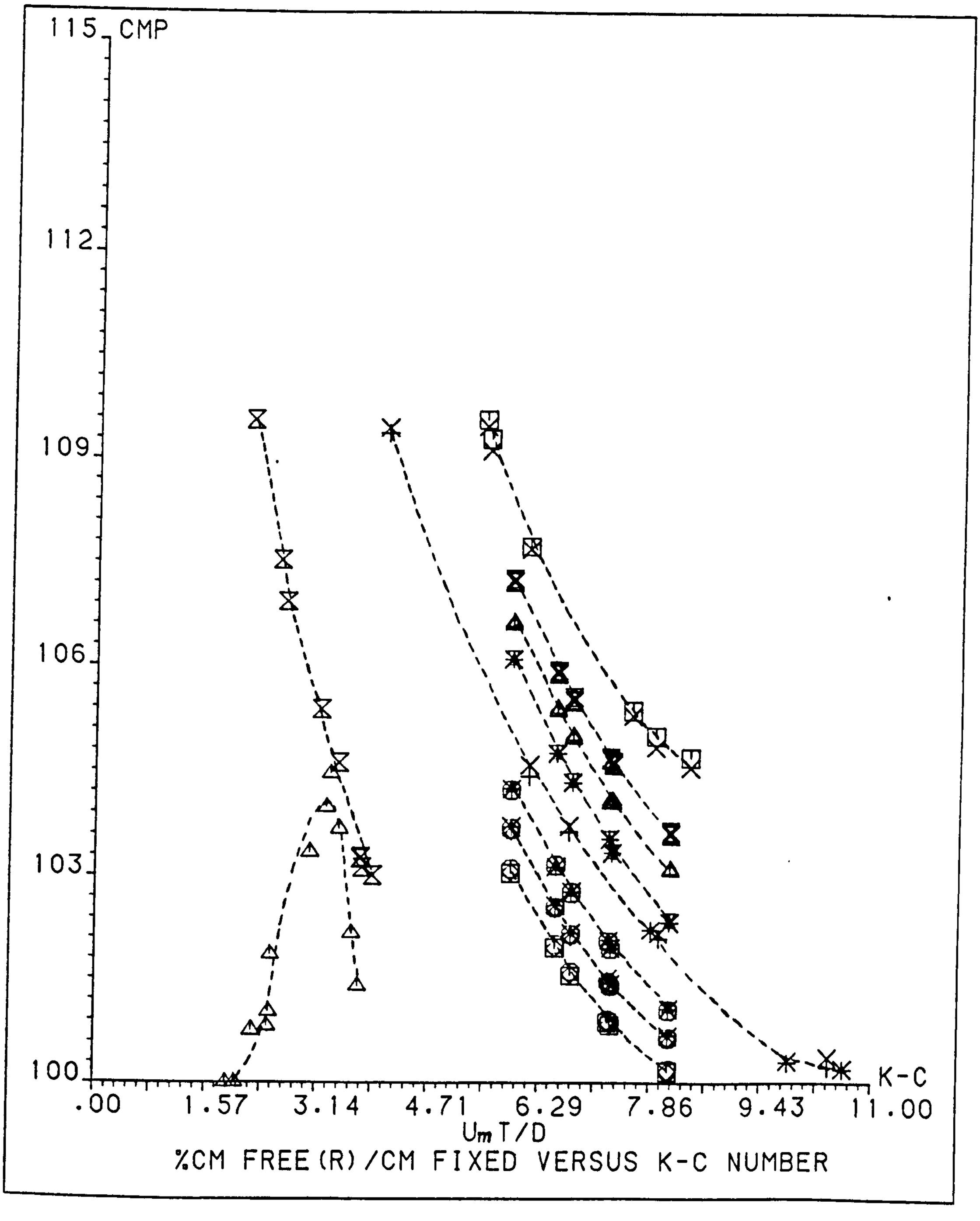

Figure (5.3.14) - Variation of % (C_M Free (R)/C_M Fixed) with K-C Number for all Structures.

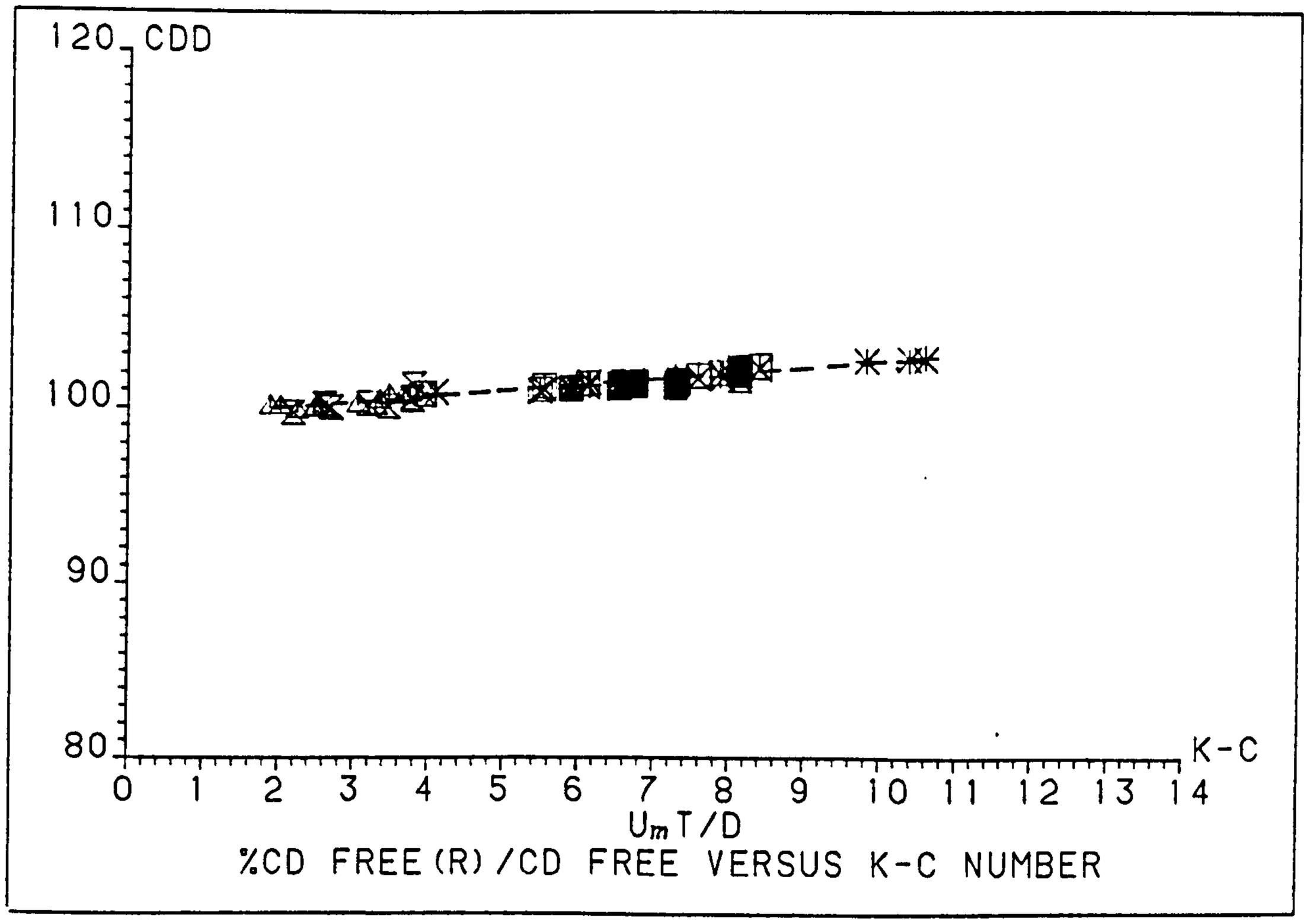

Figure (5.3.15a) - Variation of % (C_D Free (R)/C_D Free) with K-C Number for all Structures.

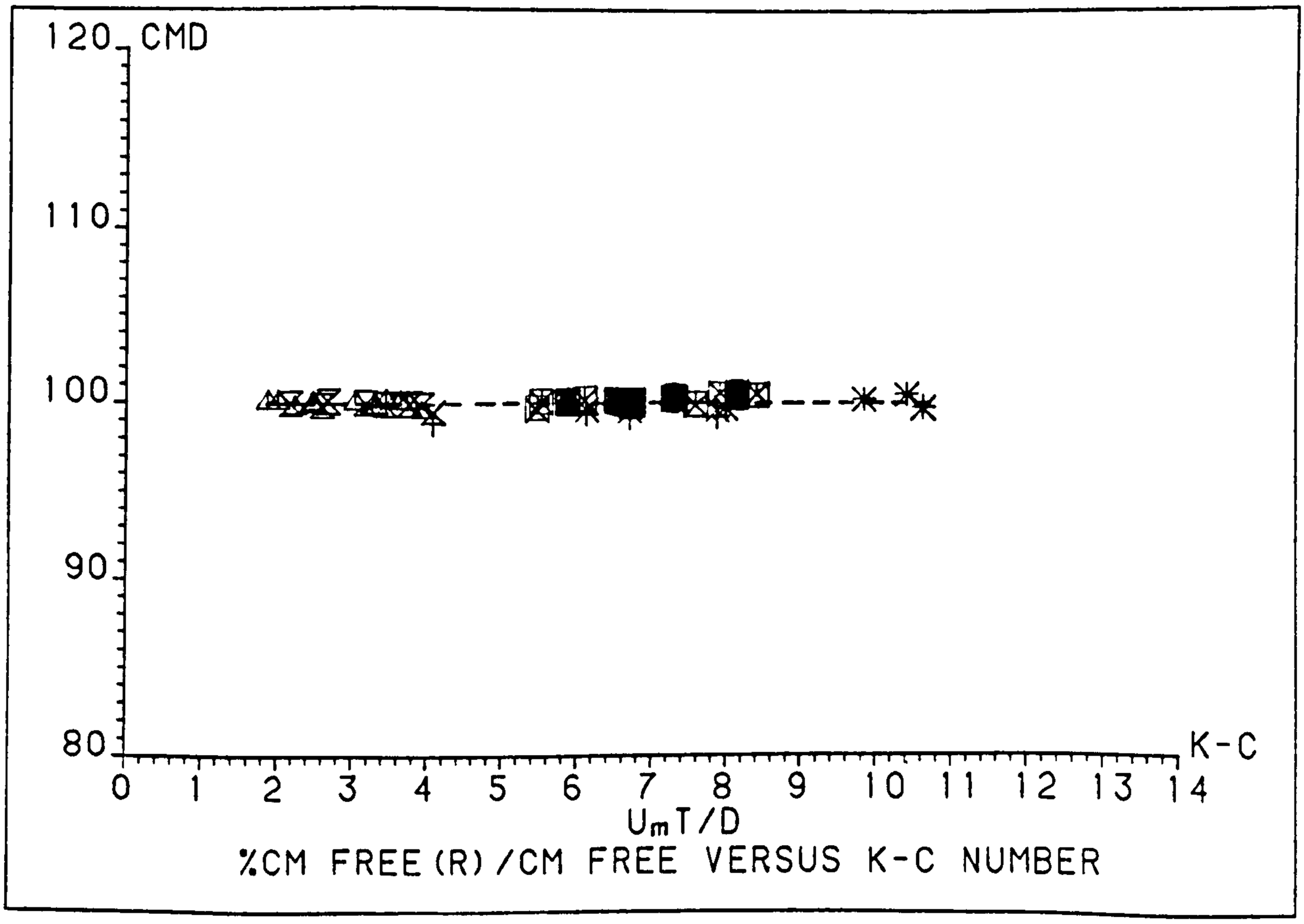

Figure (5.3.15b) - Variation of % (C_M Free (R)/C_M Free) with K-C Number for all Structures.

variation of the value of C_D and C_M along the submerged part of the cylinder.

(b) THE VARIATION OF THE PERCENTAGE RATIO OF THE FORCE COEFFICIENTS FOR FREE STRUCTURES TO THOSE FOR FIXED STRUCTURES WITH KEULEGAN-CARPENTER NUMBERS.

The variations of the percentage ratio of C_D for free structures to that for fixed structures with the K-C numbers have negative gradients with negative second derivatives for all structures with the exception of structure A (relatively stiff). In this case, for low K-C numbers there is no difference between C_D (free) and C_D (fixed) with the increasing value of K-C numbers. The percentage ratio C_D (free):C_D(fixed)then increases up to a definite maximum. Thereafter the behaviour of the C_D ratio is similar to that described above for other less stiff structures. The variations of the percentage ratio of the C_M have the same behaviour as that of the C_D for all structures with the exception that they have positive second derivatives.

The percentage ratios of the drag coefficient are greater than the corresponding percentage ratios of the inertia coefficient.

With decreasing structural stiffness, it is apparent that the rate of variation of the force coefficient ratios with the K-C numbers decreases.

For a particular value of K-C number it is apparent that, with decreasing structural stiffness, the ratios of the force coefficients increase.

This increase in the value of the force coefficients for free structures compared to those for fixed structures could be attributed to the movement of the structure and consequently the shift of the separation points in the upper stream direction towards the front of the cylinder and the increase of the wake size, thereby leading to an increase of the force. Similarly this increase occurs when the flow reverses its direction (see Figures (5.3.3) and (5.3.4)).

The negative gradient variation of the force coefficients with the K-C numbers could be due to the change of the level of vorticity in the fluid near the structure. The change of the level of vorticity is due to the increase of the K-C numbers combined with the increase of the amplitude of the structure's response

The reason for the lessening of the rate of change of the force coefficients ratio with the decreasing of the structural stiffness could be due to the increase of the phase angle shift of the structure's response to the wave force.

(c) THE VARIATION OF THE PERCENTAGE RATIO OF THE FORCE COEFFICIENTS CALCULATED BY MODIFIED MORISON'S EQUATION TO THOSE CALCULATED BY MORISON'S EQUATION WITH THE K-C NUMBERS

The variation of the percentage ratio of the C_D slightly increased with K-C numbers while no variation is noticed for the percentage ratio of C_M with the K-C numbers. This shows that for the tested waves and structures the ratio of the squared values of the relative velocity to the squared values of the water wave particle velocity is very small and decreases with the increase of the K-C numbers while the value of the ratio of the relative accelerations to the water wave particle accelerations is negligible.

5.3.2 STRUCTURE'S RESPONSE TO THE WAVE FORCES

The structures' response results obtained from the experiments conducted for the structures together with the computed values are presented below. The structures' responses were represented by the non-dimensional parameter X/D where X is the amplitude of vibration and D is structure diameter.

As the structures' responses should be related to non-dimensional parameters which contain the necessary variables to represent the structures' dynamic properties as well as the water waves. As the reduced velosity is normally used to relate the structure's response in the transverse direction due to the transverse force, because the drag component of the in-line force and the transverse force are due to the separation of flow around the structure and the formation of eddies, the structure's response in the direction of the wave can also be related to the reduced velocity. Therefore the reduced velocity was used as the non-dimensional parameter which satisfied the above condition. The reduced velocity was calculated for each test by:

$$VD = U_m/D \times FN \quad \ldots\ldots\ldots\ldots\ldots \quad (5.3.1)$$

where U_m = maximum water particle velocity

FN = structure natural frequency

Also the reciprocal of reduced velocity was calculated for each test, as it was used by T. Sarpkaye and R. L. Shoaff 1979 (113).

Because the reduced velocity does not include the effect of damping of the structure, the dimensionless damping was calculated for each test by:

$$2\bar{M}\xi/\rho\ D^2L_s \qquad \ldots\ \ldots\ \ldots\ \ldots\ \ldots\ \ldots \qquad (5.3.2)$$

where $\bar{M}$ = effective mass

ξ = damping coefficient

ρ = water density

L_s = structure length

The response parameter has been introduced by multiplying the dimensionless damping by the reciprocal of the reduced velocity

$$DM = (2\bar{M}\xi/\rho\ D^2L_s)\ (D \times FN/U_m) \qquad \ldots\ \ldots\ \ldots \qquad (5.3.3)$$

Each tested structure is presented in a table (Tables 5 - 12 through 5 - 21). Each table shows the following parameters:-

1. The above non-dimensional parameters

2. The ratio of the force frequency to the structure natural frequency

3. The percentage ratio of C_D for the free structure to that for fixed structure

4. The percentage ratio of C_M for the free structure to that for fixed structure

5. The calculated phase shift between the structure's response and the applied force

6. The structure's response calculated by using the wave forces acting on the fixed structure (Method B)

7. The structure's response calculated by using the wave forces acting on the free structure (Method A)

8. The structure's response measured from the experimental tests.

The structures' responses calculated by using the wave forces acting on the structures when they were free to vibrate, the ratio of the frequency of the force to the natural frequency of the structure, the calculated phase shift between the structure's response and the applied load and the percentage ratio of C_D

STRUCTURE (A)

TEST NO.	REDUCED VELOCITY	1/REDUCED VELOCITY	DIMENSIONLESS DAMPING	RESPONSE PARAMETER	FORCE FREQ. / NATURAL FREQ.	CD FR % CD FX	CM FR 4 CM FX	PHASE SHIFT	STRUCTURE RESPONSE METHOD(B)	STRUCTURE RESPONSE METHOD(A)	STRUCTURE RESPONSE MEASURED
1	0.03285	30.4396	0.0000017	2.0400307	0.0106067	110.10	103.52	0.0010	0.006979	0.007210	0.007040
2	0.03307	30.2432	0.0733963	2.3704560	0.0133708	107.37	101.02	0.0013	0.007735	0.007820	0.007640
3	0.05706	17.5266	0.0024011	1.0747317	0.0156961	110.22	102.16	0.0015	0.013044	0.013380	0.012440
4	0.03793	26.3636	0.0692044	1.8244800	0.0109993	111.25	103.05	0.0011	0.006955	0.007208	0.006920
5	0.03547	28.1907	0.0545571	2.5837230	0.0132473	106.56	100.61	0.0013	0.010309	0.010390	0.010304
6	0.05935	16.8502	0.0645521	1.5877126	0.0159200	108.22	101.59	0.0016	0.013970	0.014148	0.013964
7	0.02208	45.2893	0.0789562	3.5756679	0.0109929	100.00	100.00	0.0010	0.005336	0.005360	0.005323
8	0.03344	29.9008	0.0702426	2.1003103	0.0134176	107.77	101.85	0.0012	0.006949	0.007073	0.006924
9	0.05224	19.1426	0.0597975	1.1446788	0.0157312	111.32	104.44	0.0014	0.009809	0.010242	0.009665
10	0.02047	48.8590	0.0811279	3.9638270	0.0109218	100.00	100.00	0.0010	0.004947	0.004950	0.004876
11	0.03011	33.2133	0.0697751	2.3174618	0.0134916	104.42	100.75	0.0012	0.006463	0.006509	0.006440
12	0.05169	19.3450	0.0562251	1.0876711	0.0158178	110.62	103.95	0.0014	0.008507	0.008841	0.008539

Table (5.12) - Structure's Response to the Wave Force for Structure A.

STRUCTURE (C)

TEST NO.	REDUCED VELOCITY	1/REDUCED VELOCITY	DIMENSIONLESS DAMPING	RESPONSE PARAMETER	FORCE FREQ. NATURAL FREQ.	CD FR % CD FR	CM FR % CM FR	PHASE SHIFT	STRUCTURE RESPONSE METHOD(B)	STRUCTURE RESPONSE METHOD(A)	STRUCTURE RESPONSE MEASURED
1	0.12834	7.7919	0.0353260	0.2752620	0.0314901	115.53	109.37	0.0021	0.025731	0.029215	0.027365
2	0.22822	4.3817	0.0293724	0.1387025	0.0372405	114.20	104.59	0.0025	0.047718	0.049740	0.048124
3	0.33674	2.9697	0.0267382	0.0794037	0.0428260	111.20	102.22	0.0029	0.061550	0.062691	0.061545
4	0.21642	4.6207	0.0326133	0.1506949	0.0323165	113.74	103.60	0.0023	0.035137	0.039564	0.038017
5	0.30253	3.3055	0.0268020	0.0952051	0.0379391	111.00	102.09	0.0027	0.060371	0.061323	0.060339
6	0.46262	2.1616	0.0221021	0.0477757	0.0436300	108.52	100.19	0.0031	0.087303	0.089977	0.089387
7	0.40437	2.4730	0.0253946	0.0628009	0.0411736	104.30	100.29	0.0031	0.083264	0.083523	0.082950
8	0.46199	2.1646	0.0249626	0.0540332	0.0444642	101.31	100.22	0.0033	0.092158	0.092142	0.091844

Table (5.14) - Structure's Response to the Wave Force for Structure C.

STRUCTURE (D)

TEST NO.	REDUCED VELOCITY	1/REDUCED VELOCITY	DIMENSIONLESS DAMPING	RESPONSE PARAMETER	FORCE FREQ. NATURAL FREQ.	CD FR CD FX	CM FR CM FX	PHASE SHIFT	STRUCTURE RESPONSE METHOD(B)	STRUCTURE RESPONSE METHOD(A)	STRUCTURE RESPONSE MEASURED
1	0.23348	4.2830	0.0208698	0.1151580	0.0422610	118.27	109.11	0.0026	0.045601	0.050618	0.045301
2	0.28142	3.5534	0.0255770	0.0908879	0.0459358	117.71	107.67	0.0029	0.066577	0.065118	0.061838
3	0.46705	2.1411	0.0220203	0.0471473	0.0556065	113.62	104.54	0.0035	0.083482	0.087226	0.084580
4	0.21842	4.5782	0.0321081	0.1469969	0.0399537	118.58	109.46	0.0026	0.046339	0.050748	0.048195
5	0.35536	2.8140	0.0266653	0.0750369	0.0469021	115.44	105.27	0.0031	0.073816	0.073142	0.077194
6	0.42633	2.3456	0.0251987	0.0591383	0.0539374	114.67	104.82	0.0036	0.090153	0.094442	0.093321

Table (5.15) - Structure's Response to the Wave Force for Structure D.

STRUCTURE (E)

TEST NO.	REDUCED VELOCITY	1/REDUCED VELOCITY	DIMENSIONLESS DAMPING	RESPONSE PARAMETER	FORCE FREQ. / NATURAL FREQ.	CD FR / CD FX	CM FR / CM FX	PHASE SHIFT	STRUCTURE RESPONSE METHOD(B)	STRUCTURE RESPONSE METHOD(A)	STRUCTURE RESPONSE MEASURED
1	0.37301	2.6760	0.0585157	0.1566238	0.0555988	111.12	101.54	0.0077	0.039402	0.040007	0.039143
2	0.51980	1.9236	0.0530964	0.1021353	0.0639386	107.64	100.16	0.0069	0.053455	0.053551	0.052905
3	0.43264	2.3114	0.0580511	0.1341777	0.0592024	109.70	100.82	0.0082	0.045076	0.045444	0.045026
4	0.31754	3.1493	0.0604615	0.1904719	0.0541111	112.24	103.04	0.0073	0.032108	0.033085	0.032196
5	0.45251	2.2099	0.0519618	0.1148311	0.0622273	109.93	100.86	0.0084	0.043873	0.044259	0.043593
6	0.37575	2.6613	0.0568427	0.1491482	0.0576183	111.41	101.95	0.0077	0.036541	0.037256	0.036014

Table (5.16) - Structure's Response to the Wave Force for Structure E with No Mass at Top.

STRUCTURE (E)

TEST NO.	REDUCED VELOCITY	1/REDUCED VELOCITY	DIMENSIONLESS DAMPING	RESPONSE PARAMETER	FORCE FREQ. / NATURAL FREQ.	CD FR % CD FX	CM FR / CM FX	PHASE SHIFT	STRUCTURE RESPONSE METHOD(B)	STRUCTURE RESPONSE METHOD(A)	STRUCTURE RESPONSE MEASURED
1	0.47622	2.0999	0.0617065	0.1295760	0.0708692	111.72	102.12	0.0101	0.039545	0.040381	0.039159
2	0.66265	1.5091	0.0504488	0.0851671	0.0614996	108.21	100.63	0.0116	0.053877	0.054231	0.053370
3	0.55147	1.8133	0.0613062	0.1111721	0.0754626	110.32	101.38	0.0106	0.045277	0.045909	0.044960
4	0.41385	2.4163	0.0653232	0.1576422	0.0705243	112.77	103.62	0.0099	0.032372	0.033551	0.032317
5	0.58976	1.6956	0.0555992	0.0942738	0.0811030	110.44	101.46	0.0114	0.044015	0.044662	0.043428
6	0.48973	2.0420	0.0598628	0.1222372	0.0750954	111.95	102.53	0.0106	0.038645	0.037574	0.036570

Table (5.17) - Structure's Response to the Wave Force for Structure E with 1010 Grs at Top.

STRUCTURE (E)

TEST NO.	REDUCED VELOCITY	1/REDUCED VELOCITY	DIMENSIONLESS DAMPING	RESPONSE PARAMETER	FORCE FREQ. / NATURAL FREQ.	CD FR % CD FX	CM FR % CM FX	PHASE SHIFT	STRUCTURE RESPONSE METHOD(B)	STRUCTURE RESPONSE METHOD(A)	STRUCTURE RESPONSE MEASURED
1	0.61833	1.6173	0.0680163	0.1100035	0.0920175	112.09	102.31	0.0141	0.039703	0.040796	0.039511
2	0.86039	1.1623	0.0620493	0.0721179	0.1058201	108.79	101.04	0.0163	0.054185	0.054733	0.054025
3	0.71604	1.3966	0.0675067	0.0942782	0.0979816	110.88	101.94	0.0150	0.045478	0.046376	0.045233
4	0.52069	1.9205	0.0709266	0.1362158	0.0887311	113.38	104.19	0.0132	0.032537	0.033902	0.032734
5	0.74202	1.3477	0.0612004	0.0824764	0.1020408	111.11	102.05	0.0153	0.044468	0.045373	0.044137
6	0.61616	1.6230	0.0647588	0.1051013	0.0944822	112.53	103.13	0.0141	0.036743	0.037687	0.036637

Table (5.18) - Structure's Response to the Wave Force for Structure E with 2020 Grs at Top.

STRUCTURE (F)

TEST NO.	REDUCED VELOCITY	1/REDUCED VELOCITY	DIMENSIONLESS DAMPING PARAMETER	RESPONSE NATURAL FREQ.	FORCE FREQ.	CD FR % CD FX	CM FR % CM FX	PHASE SHIFT	STRUCTURE RESPONSE METHOD(B)	STRUCTURE RESPONSE METHOD(A)	STRUCTURE RESPONSE MEASURED
1	0.42342	2.3617	0.0530358	0.1252556	0.0630120	114.72	104.30	0.0085	0.050194	0.052369	0.050770
2	0.58913	1.6973	0.0401540	0.0783563	0.0724638	112.24	102.33	0.0096	0.067256	0.068805	0.066399
3	0.49033	2.0394	0.0520433	0.1061593	0.0670961	113.80	103.31	0.0090	0.057294	0.059196	0.056984
4	0.34831	2.8710	0.0541597	0.1554913	0.0593560	115.86	106.17	0.0077	0.040683	0.043147	0.041223
5	0.49037	2.0146	0.0476178	0.0959327	0.0582594	113.96	103.50	0.0089	0.056244	0.058215	0.055750
6	0.41217	2.4262	0.0506136	0.1227972	0.0632031	115.09	104.72	0.0082	0.048538	0.050736	0.047014

Table (5.19) - Structure's Response to the Wave Force for Structure F with no Mass at Top.

STRUCTURE (F)

TEST NO.	REDUCED VELOCITY	1/REDUCED VELOCITY	DIMENSIONLESS DAMPING	RESPONSE PARAMETER	FORCE FREQ. / NATURAL FREQ.	CD FR % CD FX	CM FR % CM FX	PHASE SHIFT	STRUCTURE RESPONSE METHOD(B)	STRUCTURE RESPONSE METHOD(A)	STRUCTURE RESPONSE MEASURED
1	0.54866	1.8226	0.0551943	0.1005987	0.0816493	115.32	104.93	0.0113	0.050118	0.052573	0.050453
2	0.76344	1.3099	0.0466760	0.0611388	0.0933967	112.81	103.00	0.0131	0.060254	0.068300	0.066210
3	0.63536	1.5739	0.0537652	0.0846189	0.0669414	114.56	104.00	0.0121	0.056992	0.059269	0.055773
4	0.47424	2.1087	0.0395726	0.1256178	0.0808145	116.40	106.57	0.0111	0.041445	0.044171	0.040824
5	0.67582	1.4797	0.0488272	0.0722493	0.0929368	114.58	104.00	0.0127	0.055498	0.057711	0.054666
6	0.56118	1.7820	0.0544557	0.0970374	0.0860526	113.63	105.32	0.0118	0.048850	0.049352	0.047214

Table (5.20) - Structure's Response to the Wave Force for Structure F with 1010 Grs at Top.

STRUCTURE (F)

TEST NO.	REDUCED VELOCITY	1/REDUCED VELOCITY	DIMENSIONLESS DAMPING	RESPONSE PARAMETER	FORCE FREQ. NATURAL FREQ.	CD FR CD FX	CM FR CM FX	PHASE SHIFT	STRUCTURE RESPONSE METHOD(B)	STRUCTURE RESPONSE METHOD(A)	STRUCTURE RESPONSE MEASURED
1	0.70400	1.4205	0.0521627	0.0882993	0.1047009	115.96	105.47	0.0159	0.050637	0.053419	0.050692
2	0.97960	1.0208	0.0367555	0.0579375	0.1204819	113.40	103.57	0.0183	0.069199	0.071671	0.069468
3	0.81525	1.2266	0.0387554	0.0745239	0.1115573	115.33	104.56	0.0169	0.057685	0.060312	0.057597
4	0.59335	1.6854	0.0648626	0.1093165	0.1011122	117.04	107.19	0.0149	0.041470	0.044463	0.040180
5	0.84556	1.1827	0.0571037	0.0675339	0.1162791	115.17	104.64	0.0172	0.057344	0.060000	0.056327
6	0.70213	1.4242	0.0665368	0.0862167	0.1076658	116.25	105.66	0.0159	0.047361	0.050167	0.047727

Table (5.21) - Structure's Response to the Wave Force for Structure F with 2020 Grs at Top.

for the free structures to that of the fixed structures as well as the percentage ratio of C_M were plotted against the two non-dimensional parameters (the reciprocal of the reduced velocity and the response parameter).

The variation of the structures' responses with the reciprocal of the reduced velocity is shown in Figure (5.3.16a). The variation of the structures' responses the response parameter is shown in Figure (5.3.16b). The figures show that the structures' responses can be represented by these non-dimensional parameters. The structural responses increase with the decrease of these non-dimensional parameters. The response parameter represents more accurately the structural response as the graph shows less scattered results.

Figures (5.3.17a) and (5.3.17b) show the variation of the phase shift with the reciprocal of the reduced velocity and the response parameter respectively.

The variation of the ratio of the force frequency to the structure frequency with the reciprocal of the reduced velocity and with the response parameter are shown in Figures(5.3.18a) and (5.3.18b), respectively.

Although the variation of the structure's response is better represented by the response parameter, the variation of the phase shift and the variation of the ratio of the force frequency to the structure frequency are better represented by the reciprocal of the reduced velocity. The variations of these three parameters with respect to the reciprocal of the reduced velocity and response parameter have the same pattern.

The variation of the percentage ratio of the drag coefficient calculated for the free structures to that calculated for fixed structures with the reciprocal of the reduced velocity is shown in Figure (5.3.19). The variation of this percentage ratio with the response parameter is shown in Figure (5.3.20). This variation has a positive gradient with a negative second derivative for all the structures.

Figures (5.3.21) and (5.3.22) show the variation of the percentage ratio of the inertia coefficient calculated for free structures to that calculated for fixed structures with reciprocal of the reduced velocity and with the response parameter respectively. The variations have positive gradients with positive second derivatives for all structures. The variation of the percentage ratio of the force coefficients with the reciprocal of the reduced velocity is clearer than that with the response parameter.

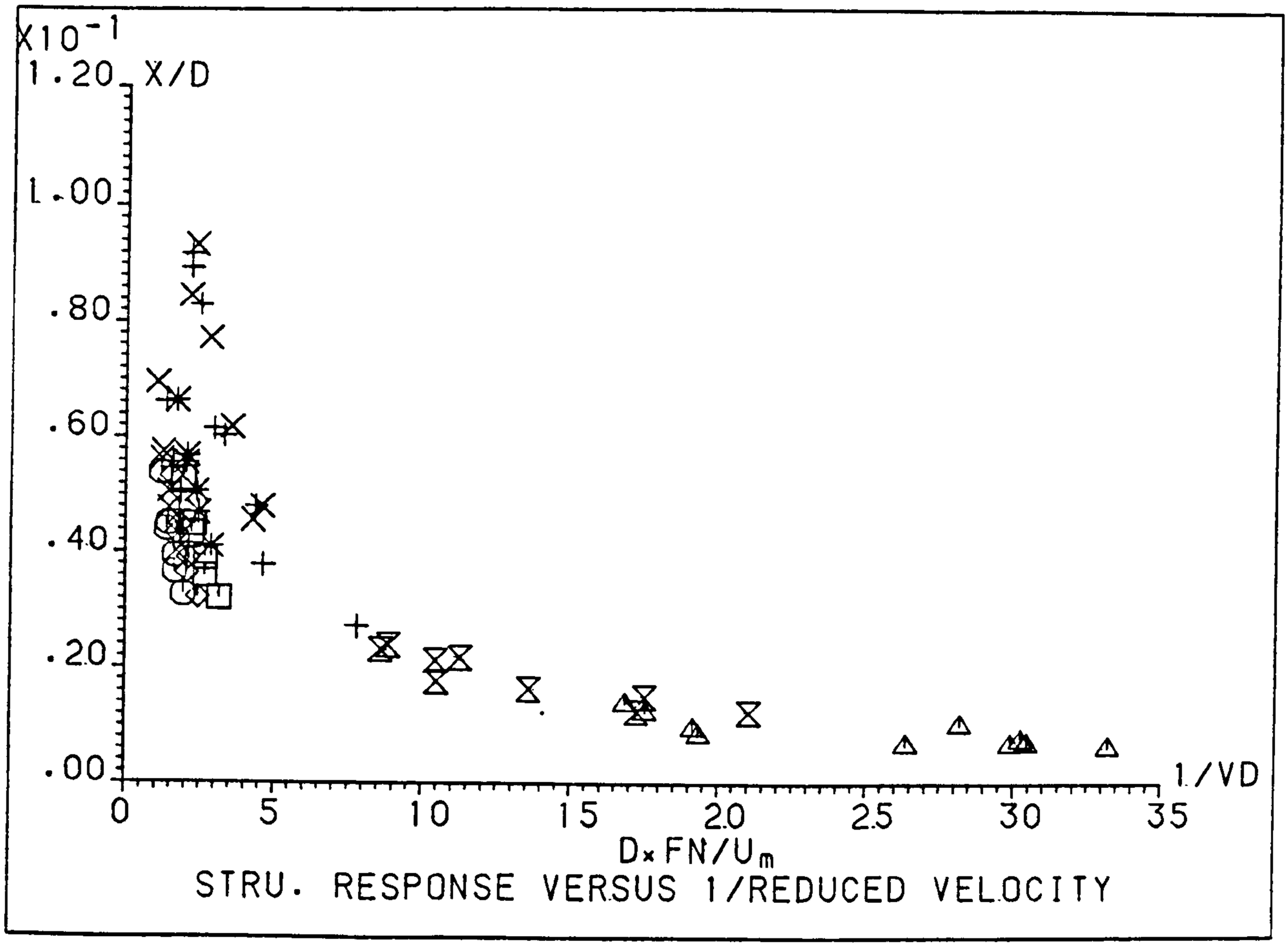

Figure (5.3.16a) - Variation of the Structures' Response (X/D) with the Reciprocal of Reduced Velocity.

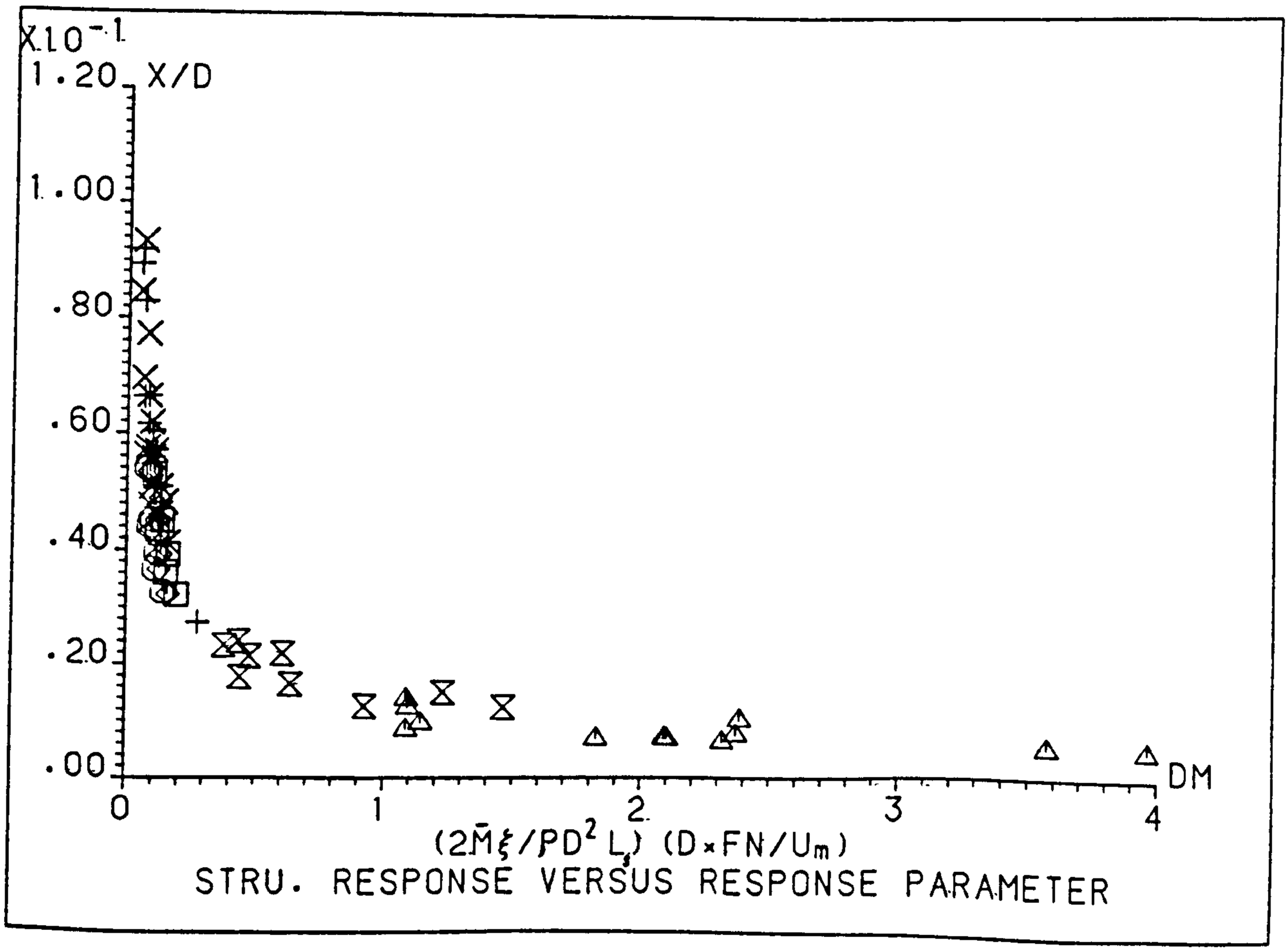

Figure (5.3.16b) - Variation of the Structures' Response (X/D) with the Response Parameters.

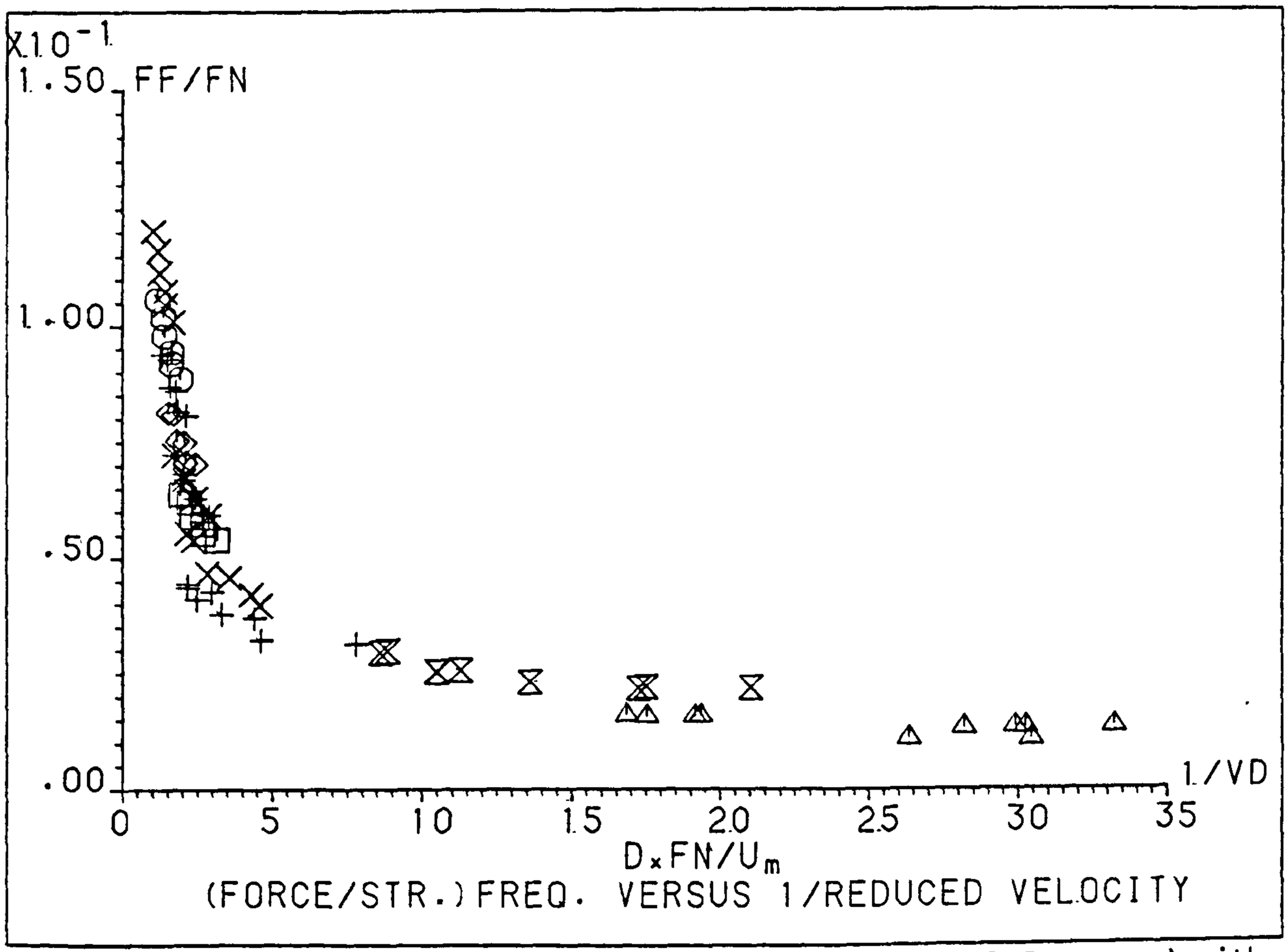

Figure (5.3.17a) - Variation of the (Force Frequency/Natural Frequency) with the Reciprocal of Reduced Velocity.

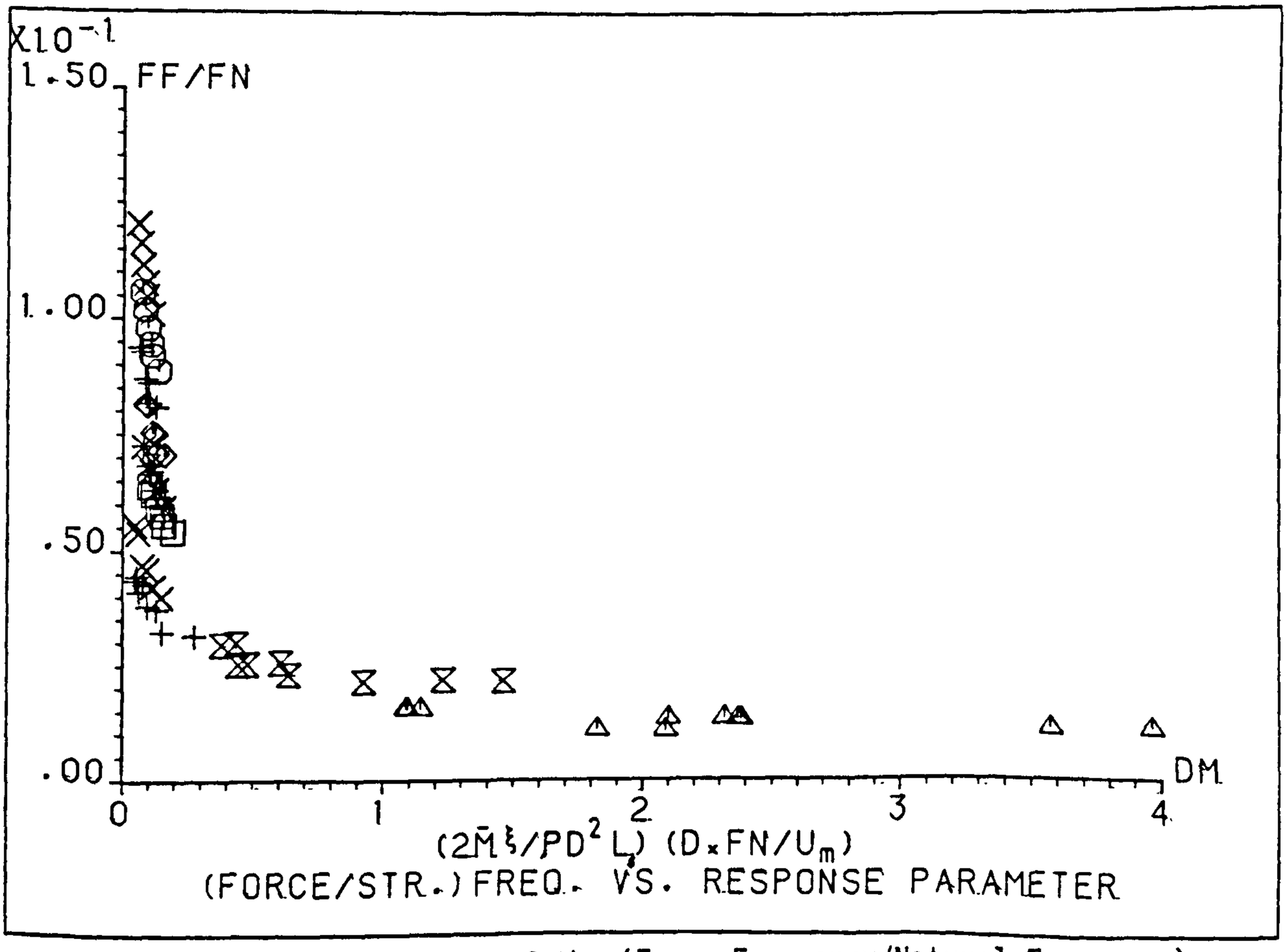

Figure (5.3.17b) - Variation of the (Force Frequency/Natural Frequency) with the Response Parameters.

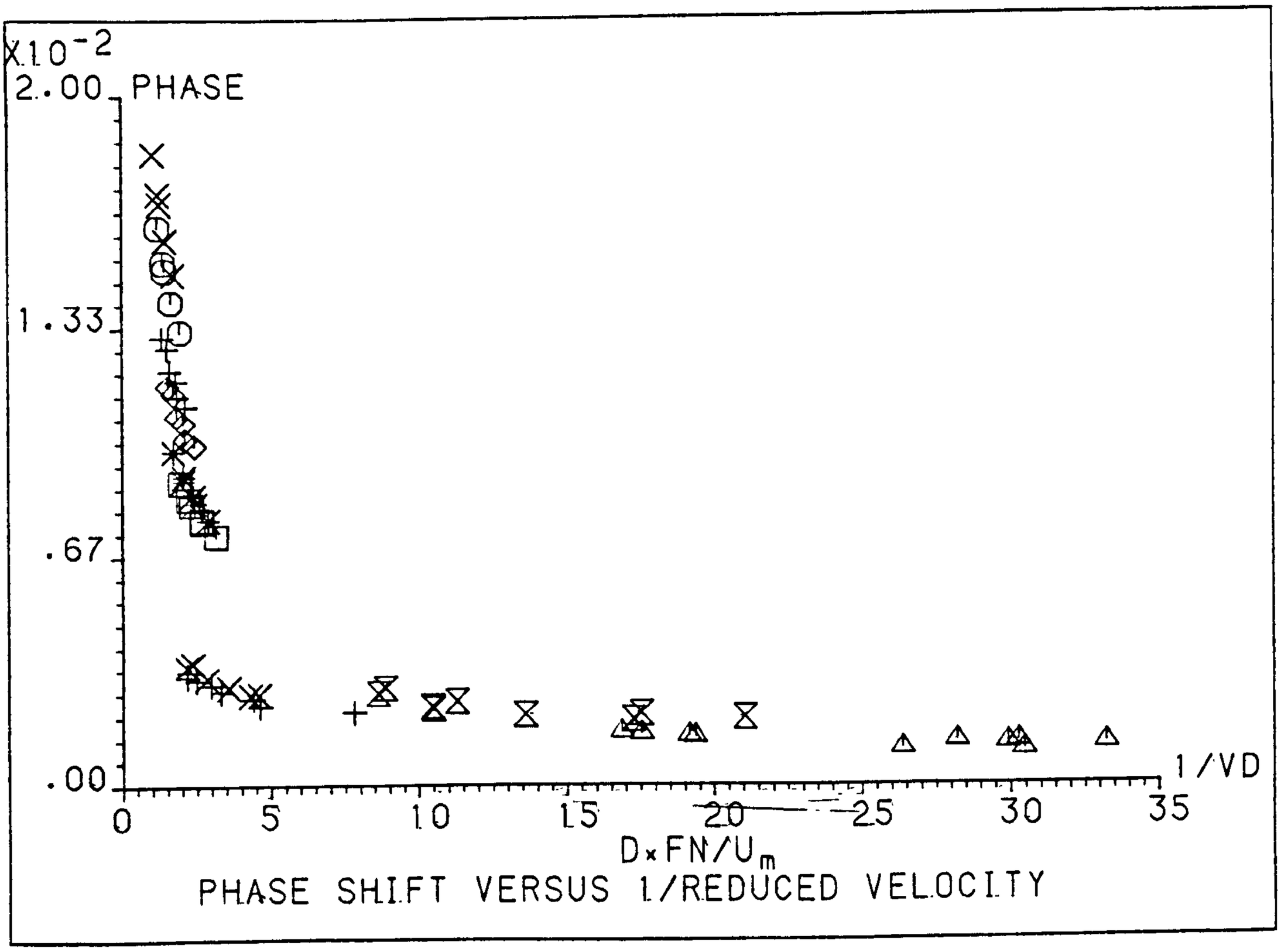

Figure (5.3.18a) - Variation of the Phase Shift between the Structure's Response and the Force with the Reciprocal of Reduced Velocity.

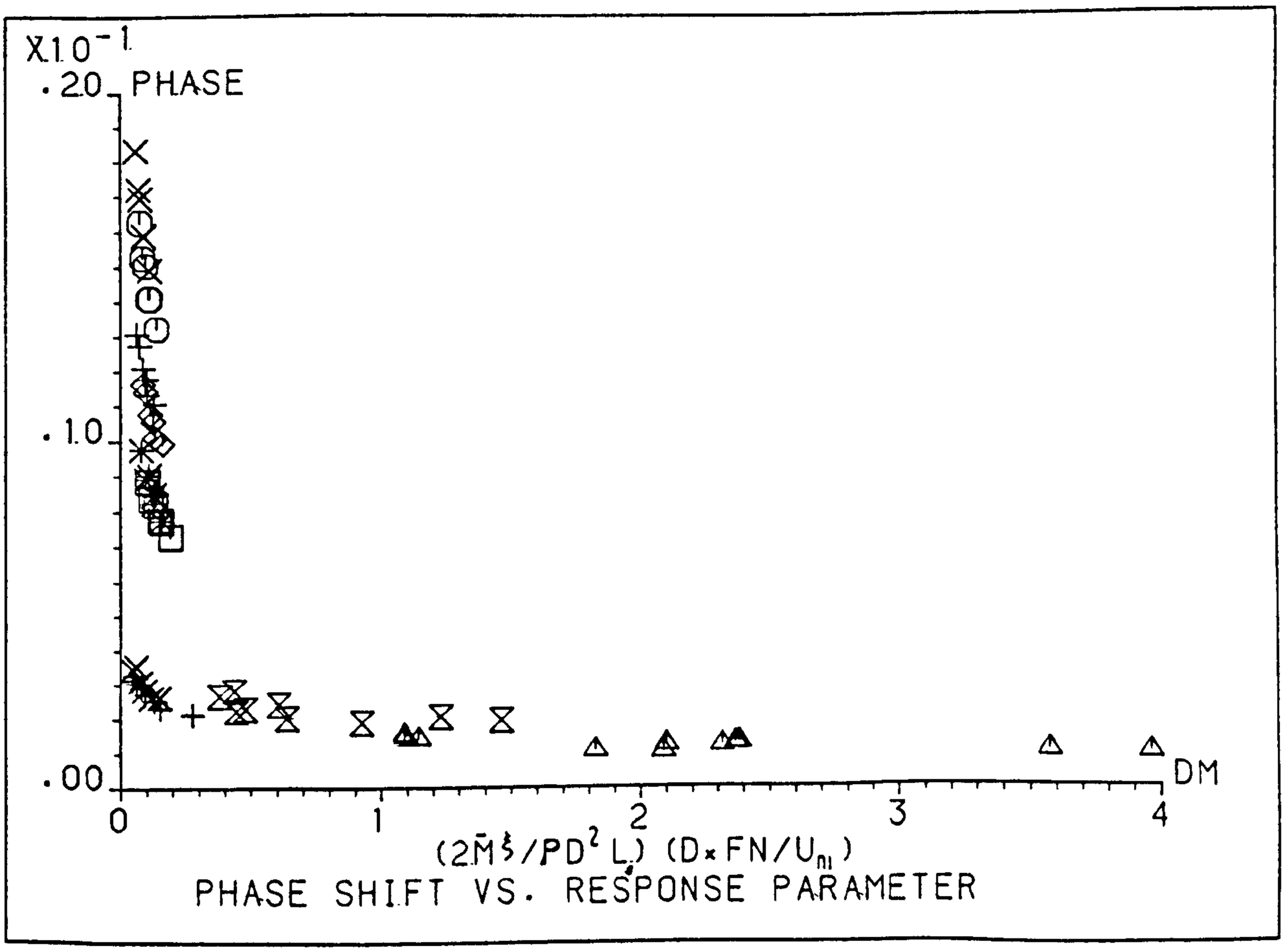

Figure (5.3.18b) - Variation of the Phase Shift between the Structure's Response and the Force with the Response Parameters.

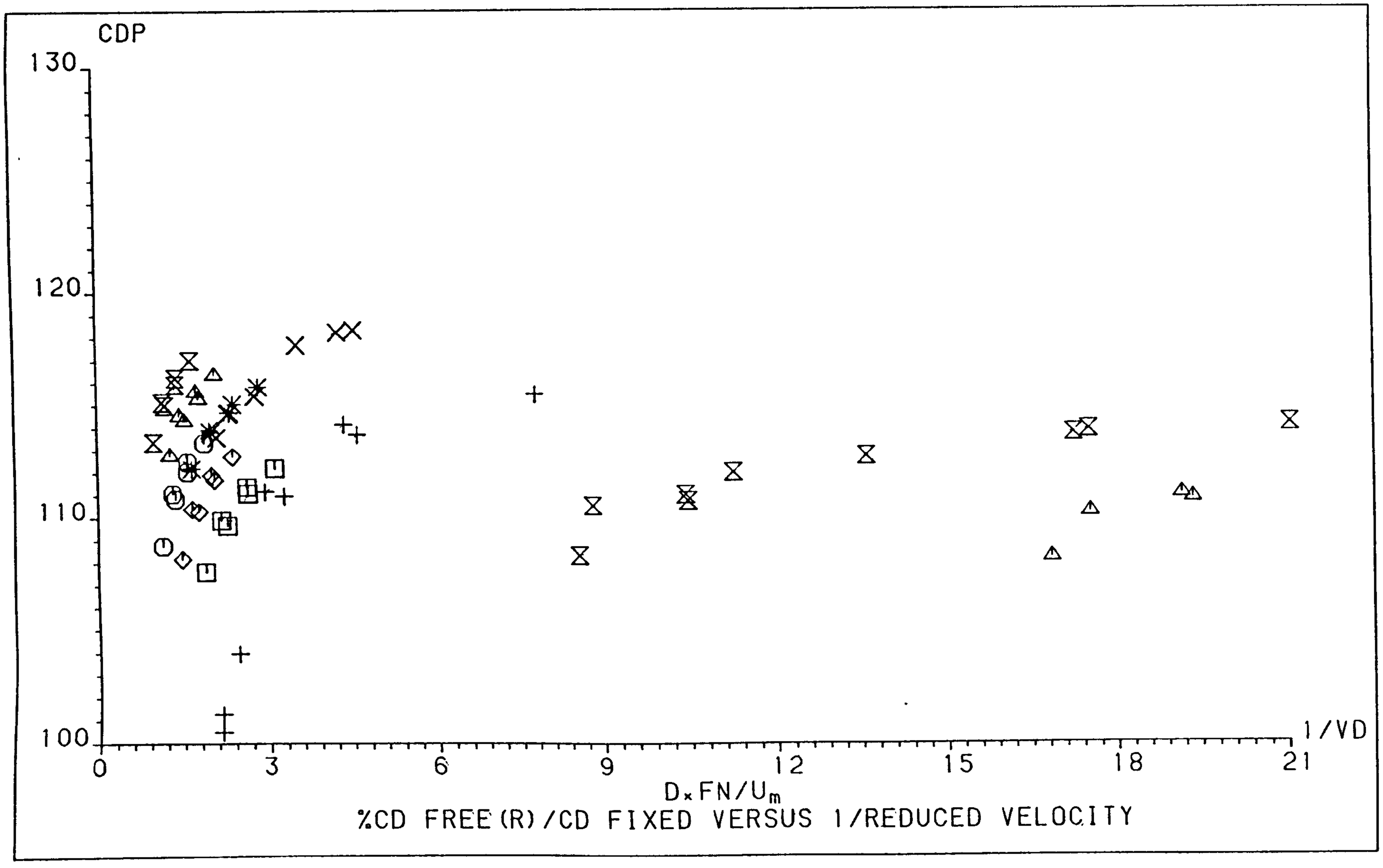

Figure (5.3.19) - Variation of (% C_D Free(R)/C_D Fixed) with the Reciprocal of Reduced Velocity.

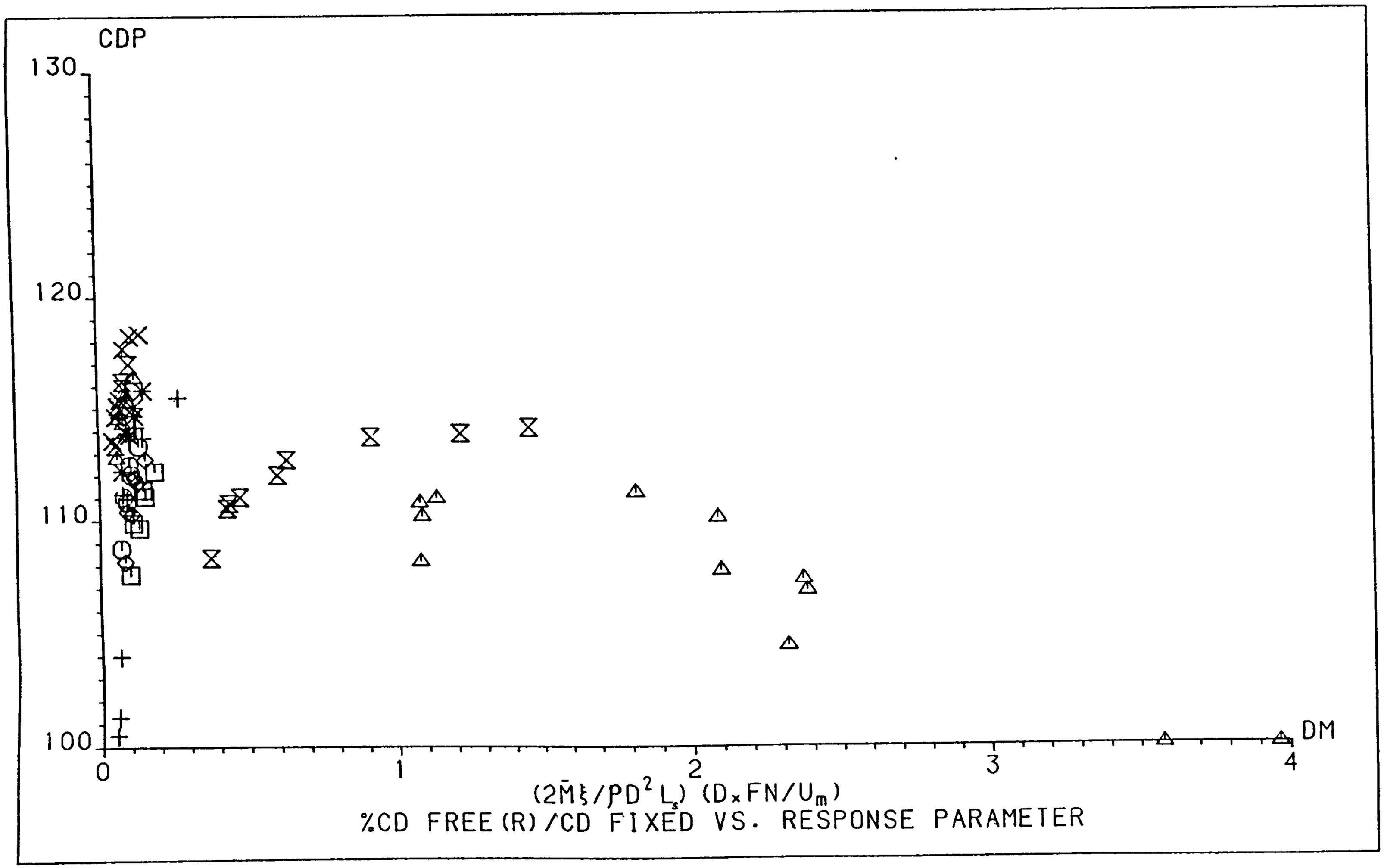

Figure (5.3.20) - Variation of (% C_D Free(R)/C_D Fixed) with the Response Parameter.

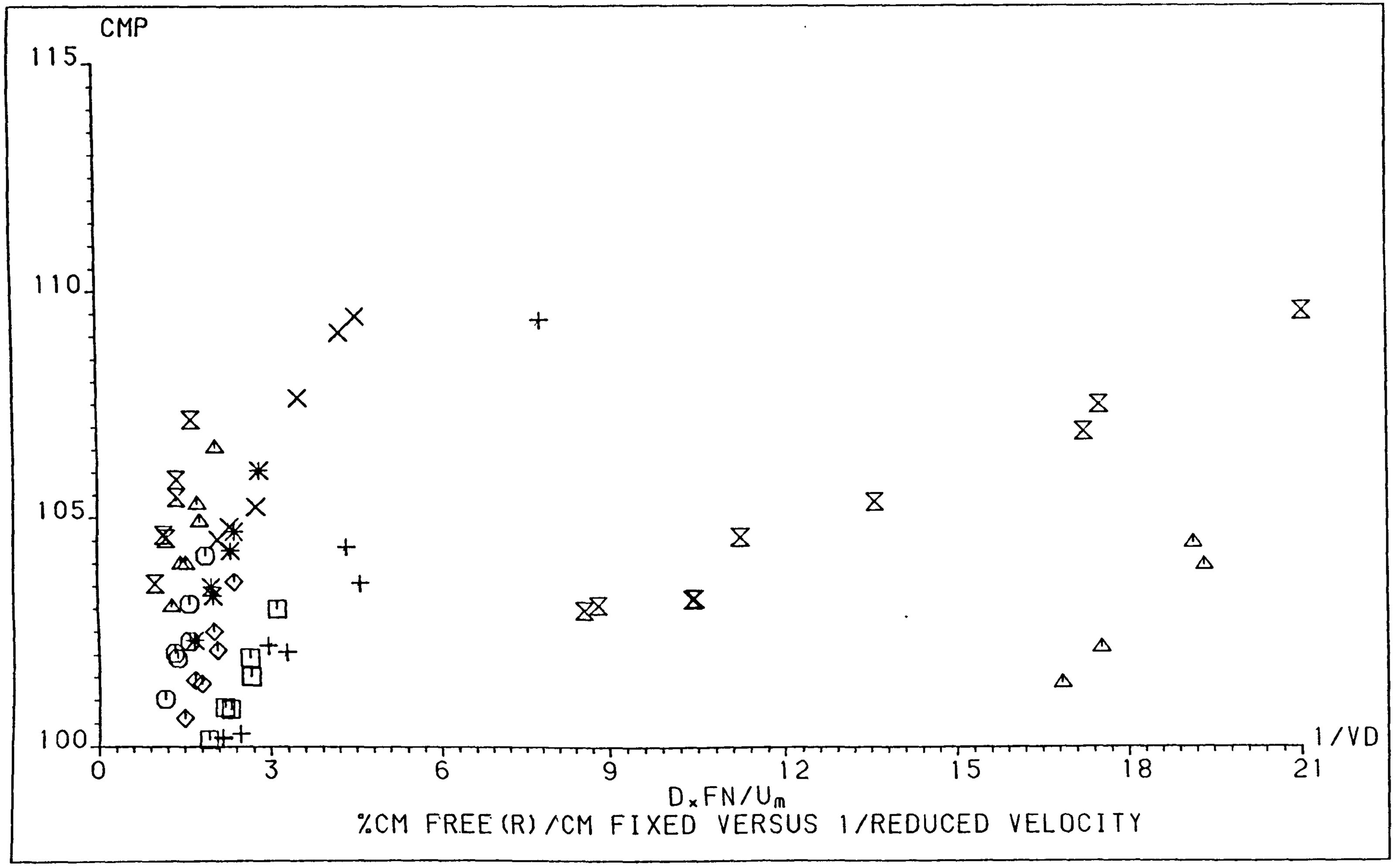

Figure (5.3.21) - Variation of (% C_M Free(R)/C_M Fixed) with the Reciprocal of Reduced Velocity.

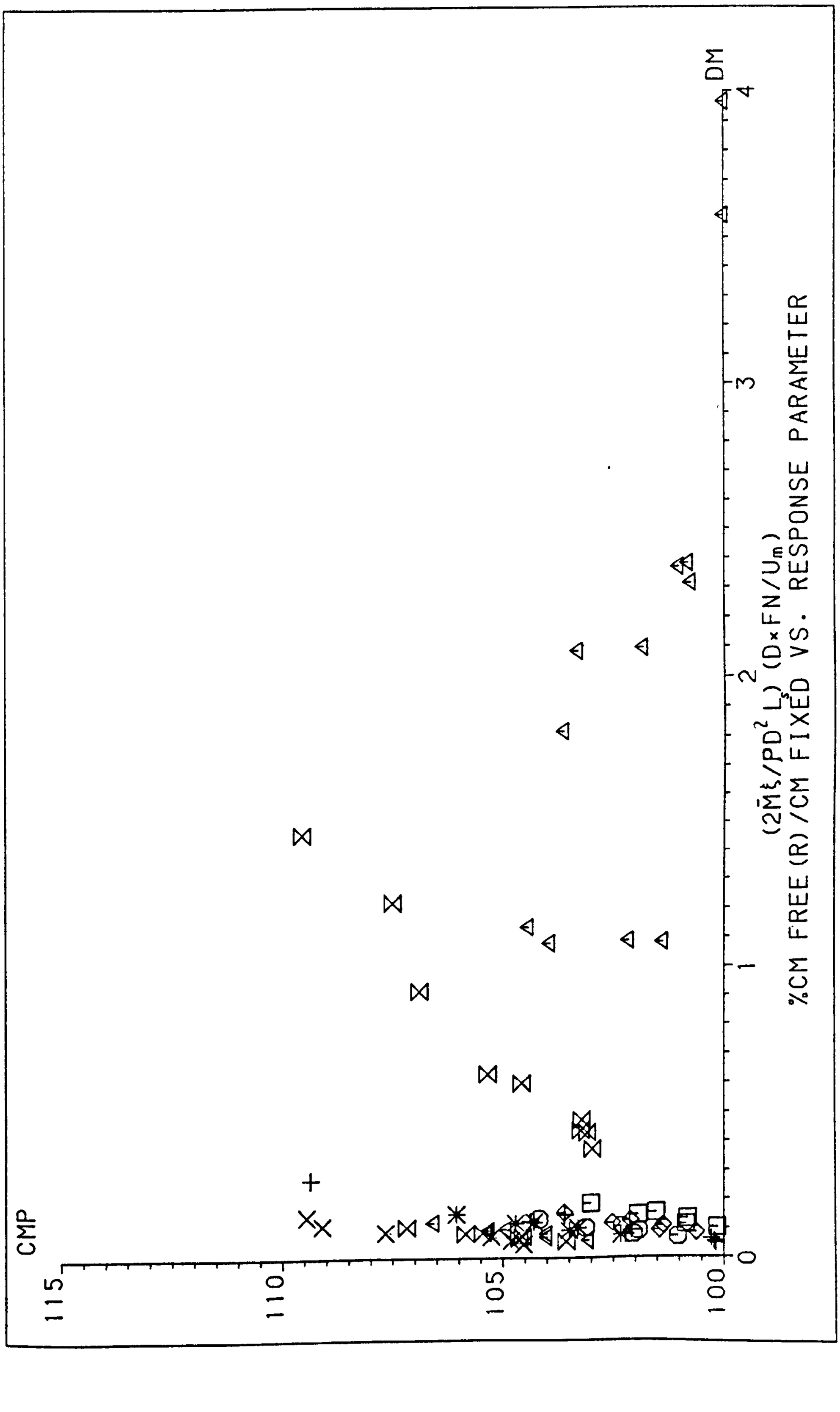

Figure (5.3.22) - Variation of (% C_M Free(R)/C_M Fixed) with the Response Parameter.

The variations of the percentage ratio of C_D and also of C_M with respect to the reduced velocity have a similar pattern to the variations with the K-C numbers. This similarity in the behaviour emphasises the previous explanation of the water wave structure interaction.

The variation of the relative displacement X/D with the time of one wave cycle is shown in Figures(5.3.23) through (5.3.32) for one test of each structure when the water depth was 0.5 meter and the wave period was 1.15 seconds. In these figures the relative displacements are represented by the following symbols:

- The measured displacement divided by the structure diameter.

o The calculated relative displacement using the values of C_D and C_M for free structure and taking into consideration the water wave/structure interaction in the dynamic equation.

\+ The calculated relative displacement using the values of C_D and C_M for fixed structure and using the damping coefficient in water without taking into consideration the water wave/ structure interaction in the dynamic equation.

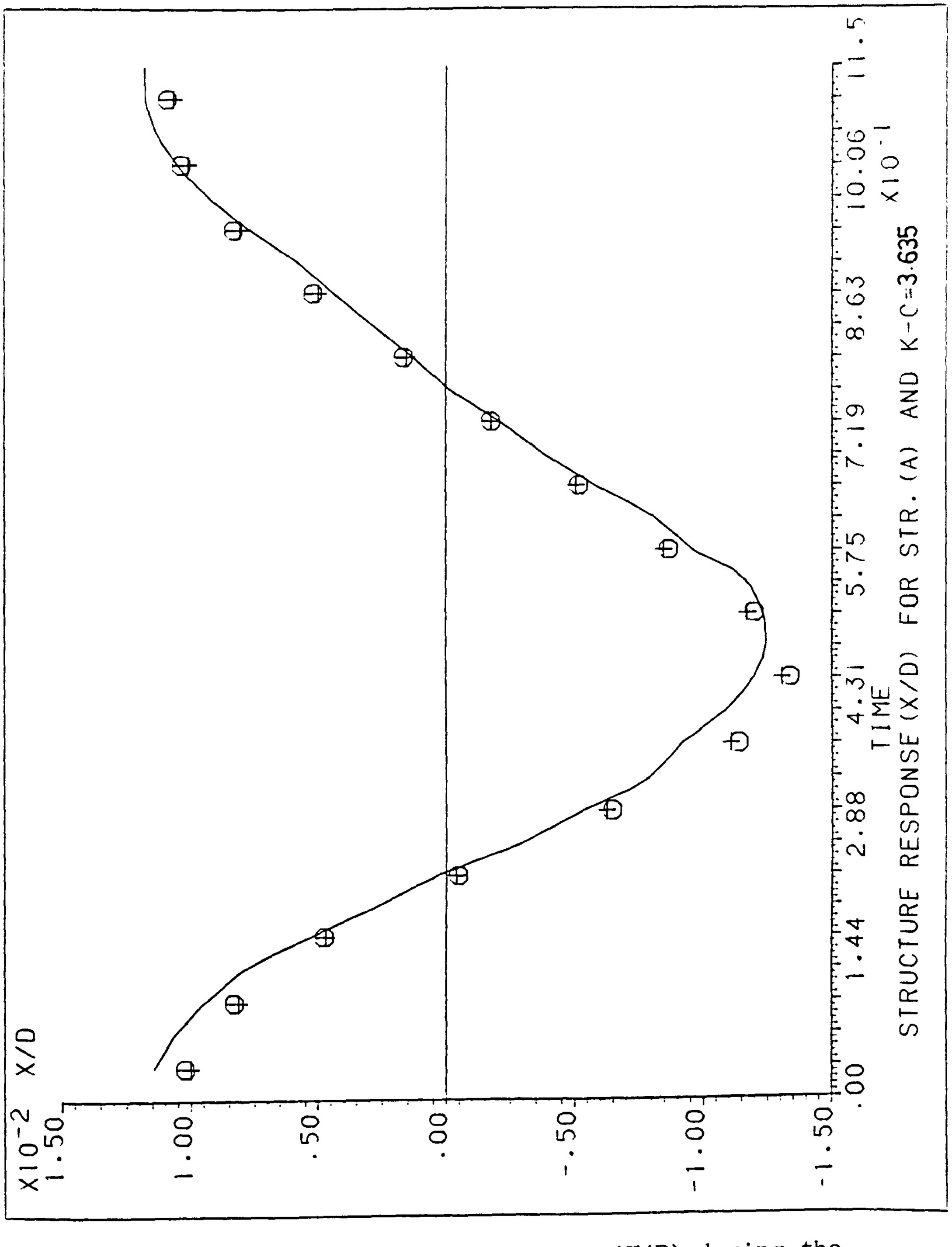

Figure (5.3.23) - Structure's Response (X/D) during the Wave Cycle for Structure A.

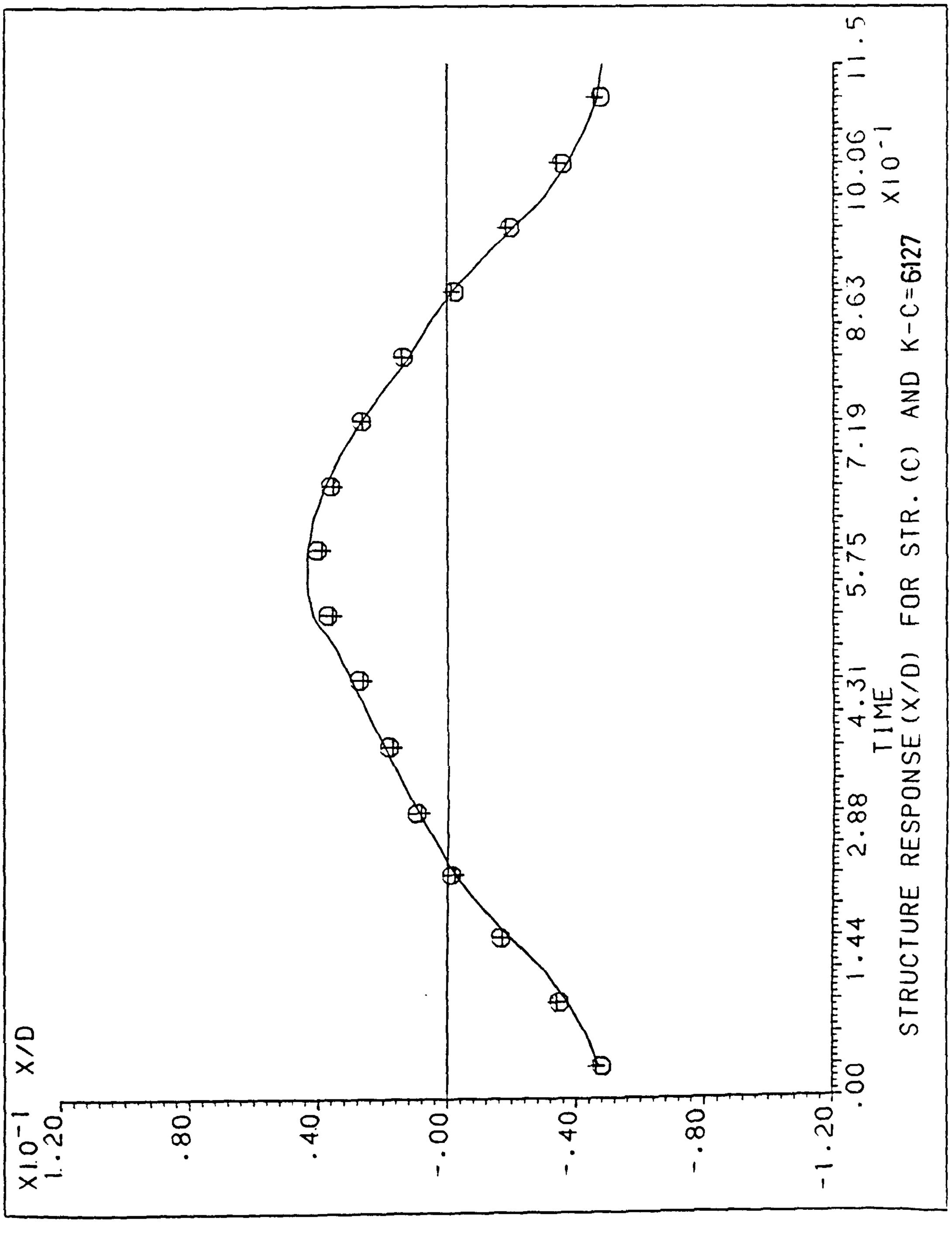

Figure (5.3.25) - Structure's Response (X/D) during the Wave Cycle for Structure C.

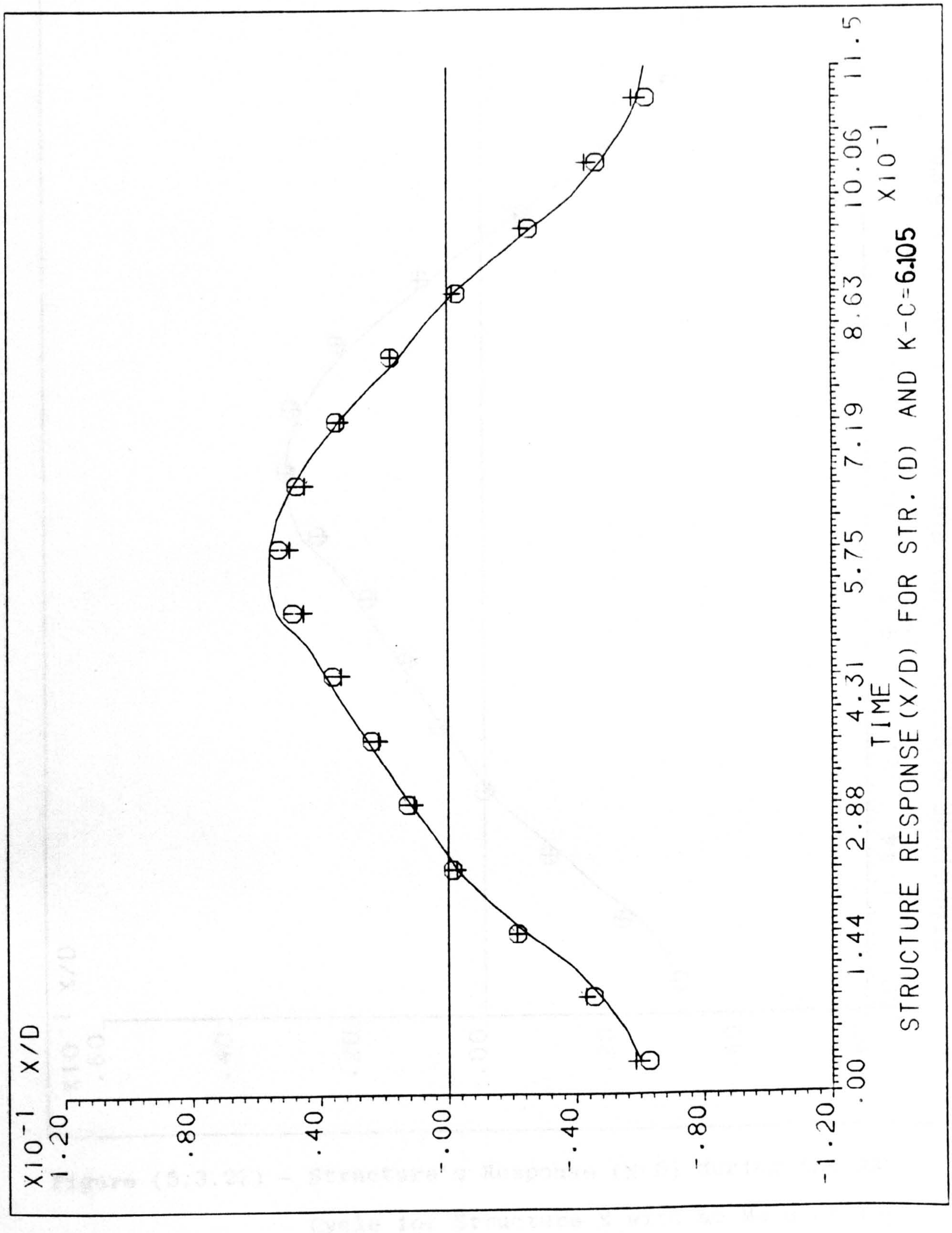

Figure (5.3.26) - Structure's Response (X/D) during the Wave Cycle for Structure D.

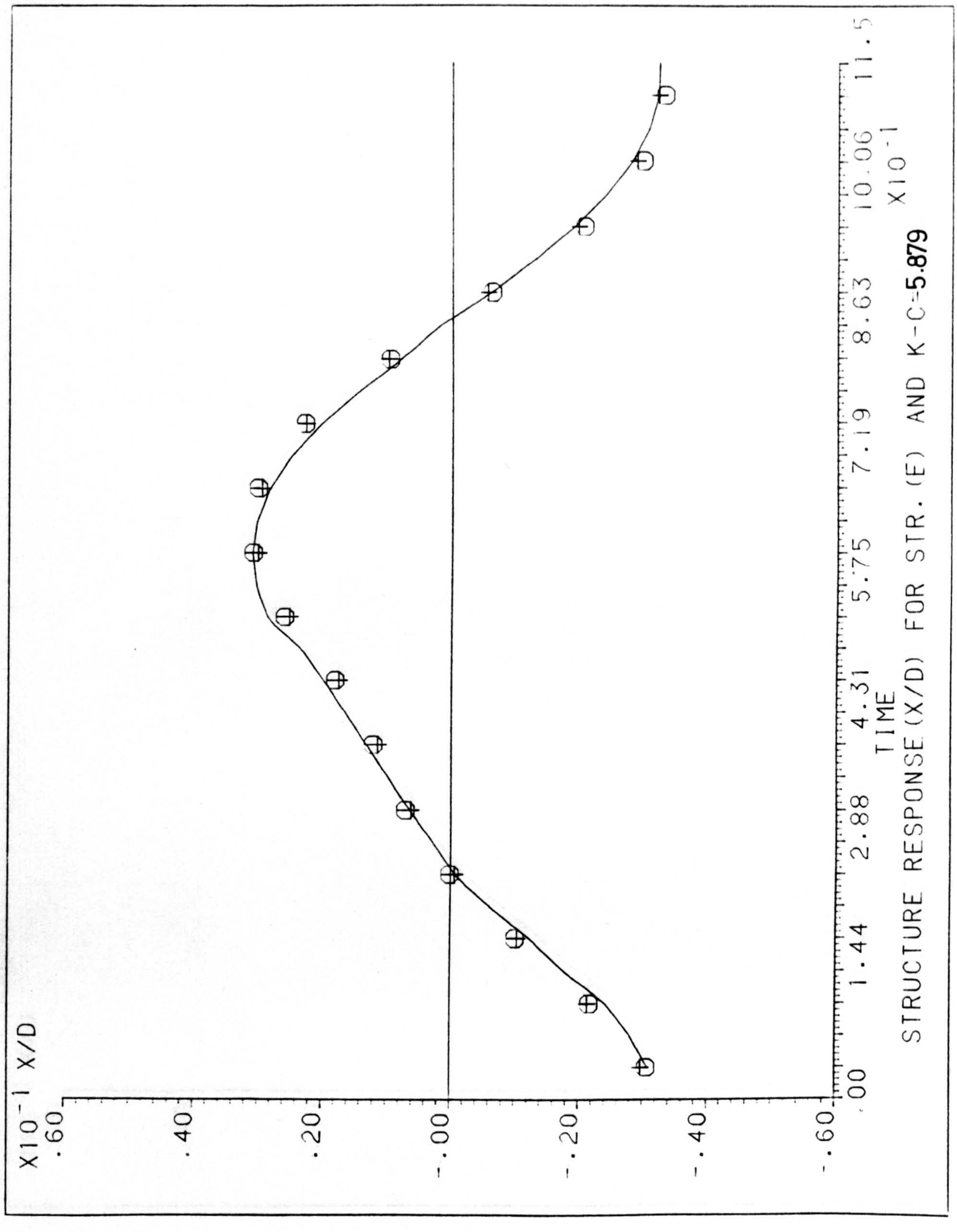

Figure (5.3.27) - Structure's Response (X/D) during the Wave Cycle for Structure E with no Mass at Top.

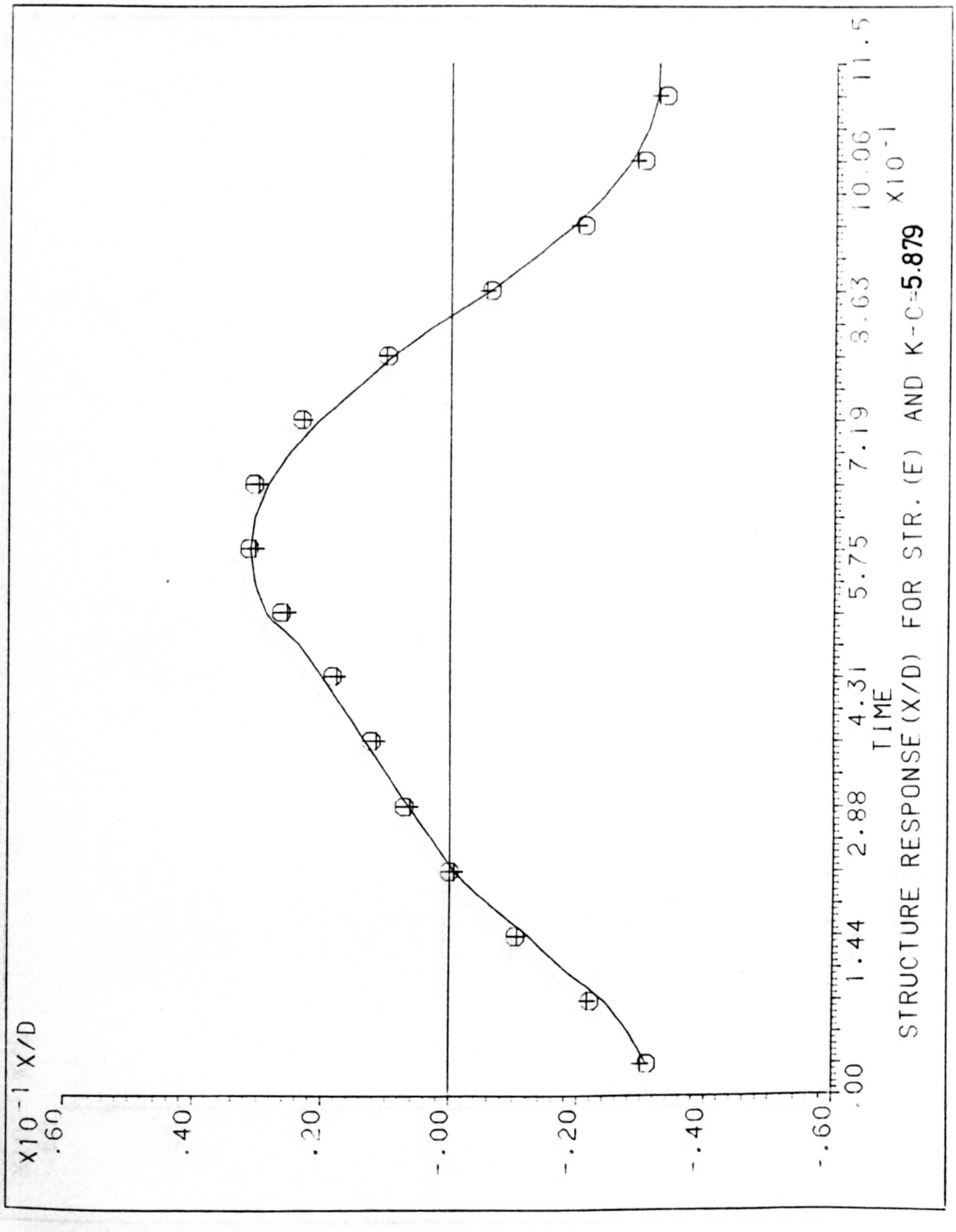

Figure (5.3.28) - Structure's Response (X/D) during the Wave Cycle for Structure E with 1010 Grs at Top.

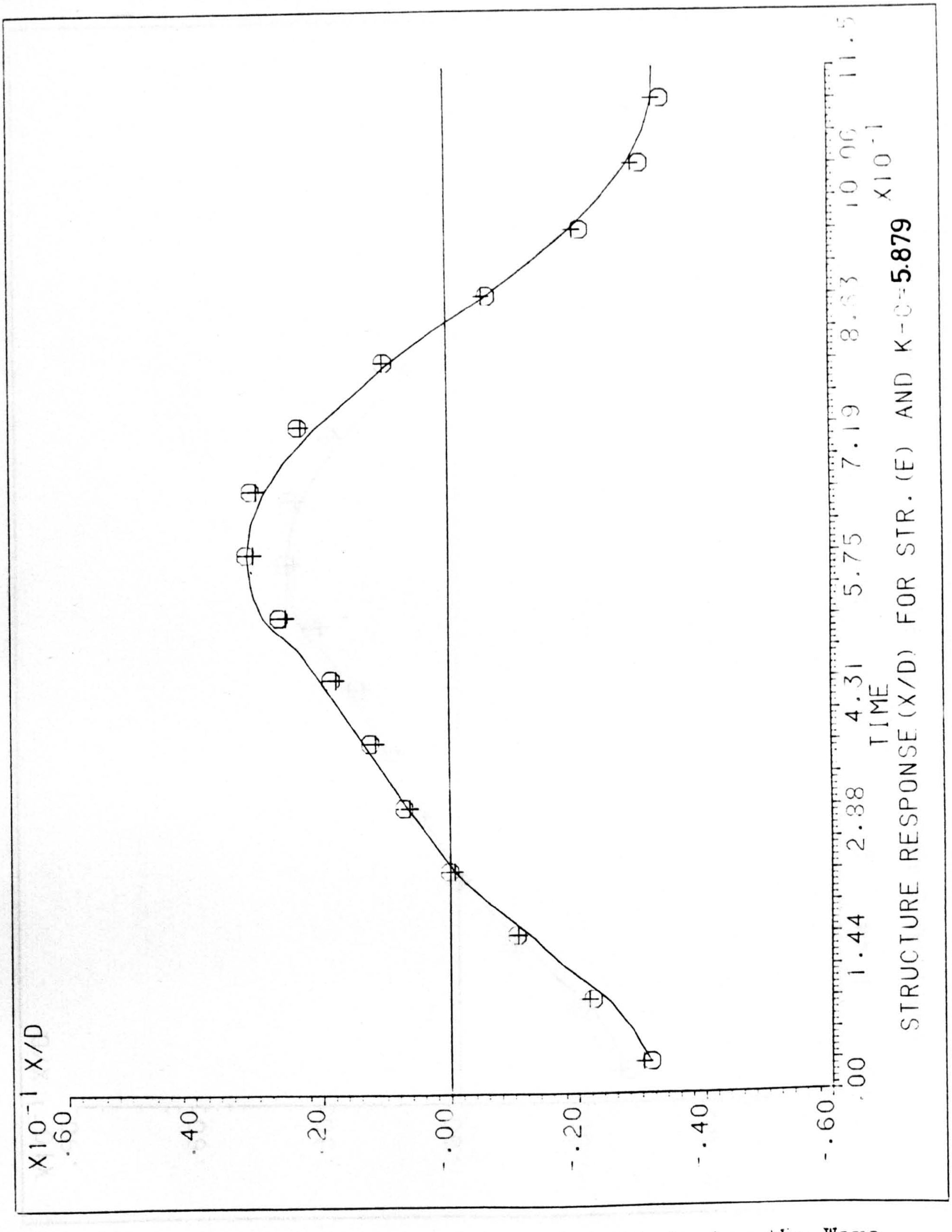

Figure (5.3.29) - Structure's Response (X/D) during the Wave Cycle for Structure E with 2020 Grs at Top.

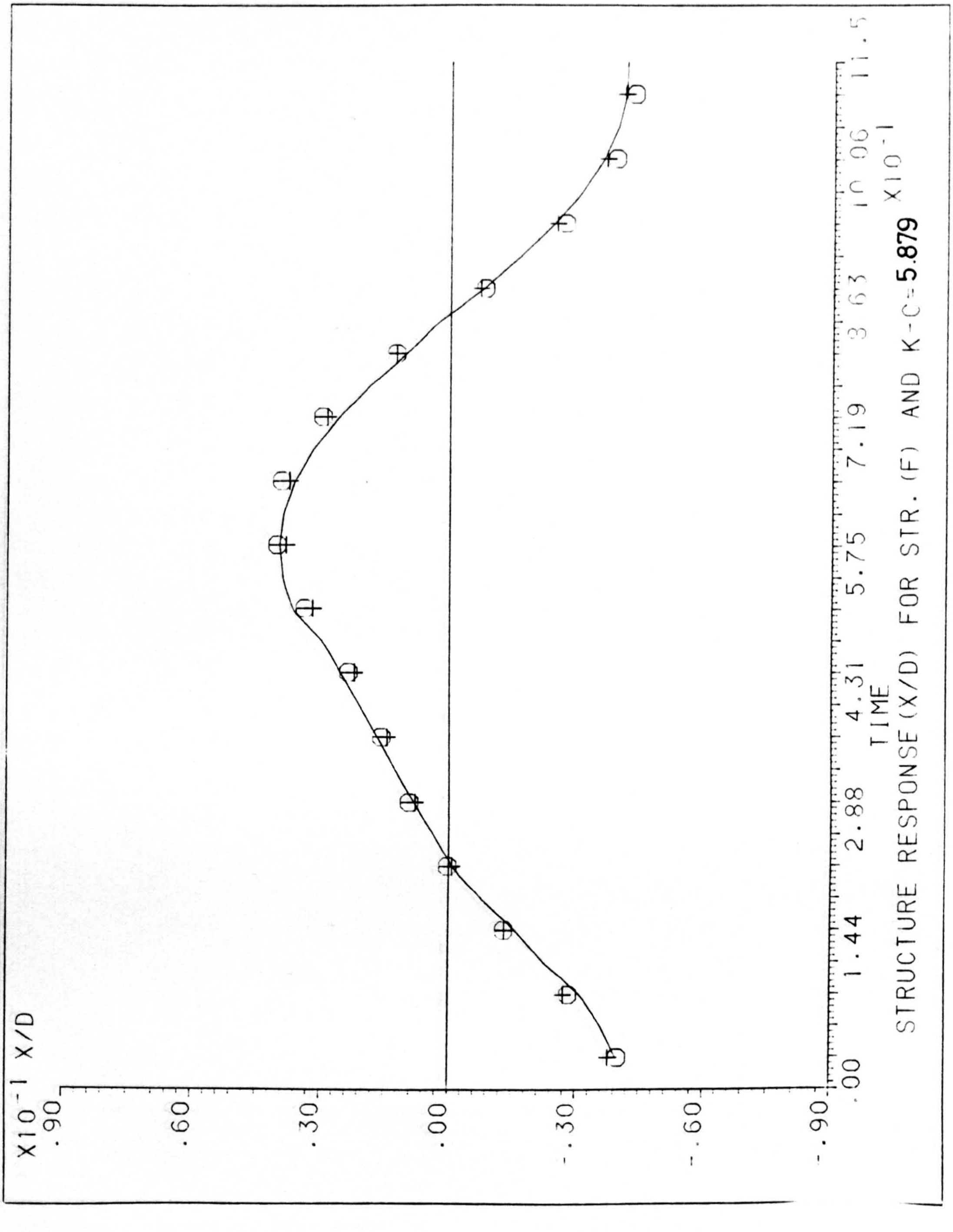

Figure (5.3.30) - Structure's Response (X/D) during the Wave Cycle for Structure F with no Mass at Top.

Results for the other tests are shown in Appendix (D)

The figures show that there are insignificant differences between the two calculated relative displacements, the greatest difference being 11.08%. Also the difference between the calculated relative displacement using finite elements and the measured relative displacement is small, with a maximum difference of 9.63%.

CHAPTER SIX

SUMMARY AND CONCLUSION

6.1.1 Theoretical and experimental investigations were carried out in a regular wave regime to study the wave force and the associated structural response on a circular cylinder pier. The cylinders were tested under three different conditions (a) fixed at the bottom and the top end free to vibrate as a cantilever (b) fixed at both ends top and bottom, (c) fixed at the bottom but free to vibrate, top end being free carrying an added mass. The experimental study was conducted in a wave regime defined by Keulegan-Carpenter number (Um T/D) i.e. from 1.87 to 10.69. The maximum amplitude of the structural response (X/D) at the top of the cylinder varied from 0.0049 to 0.098. The ratio of the force frequency to the structural natural frequency was low to avoid resonance.

6.1.2 The theoretical model used in this study were

(a) The structural model is an equivalent beam arrangement (pipe element) consisting of lumped masses, the force being applied to the mass node.

(b) The damping model for the structure is based on the mass and stiffness damping. In the case when the water wave/structure interaction is ignored and the value of C_D and C_M were taken for a fixed structure, the value of the relative damping coefficient taken from the water damping test was used.

(c) The direct approach method was used in the iterative solution of the nonlinear differential equation of motion in the case of water wave/structure interaction.

(d) The mode superposition technique with few modes was used in the two-dimensional analysis of the structural response in the time domain.

6.2.1 The local force coefficitnts C_D and C_M were calculated at certain depth of the structure (where the pressure could be measured). This depth varied from 0.06 to 0.12 metres from the still water level (see Tables 5.2 to 5.11). For some structures, there were more than one level where the local force coefficients were calculated. It was found that the magnitude of the drag coefficient decreases with depth and that of the

inertia coefficient increases with depth (e.g. in the Table 5.5 Test 1, the value of C_D at 0.05 metres from the still water level was 1.019 and at level 0.15 was 0.978 and the corresponding C_M value was 1.758 and 1.837).

6.2.2 When the structures were fixed at the top i.e. not allowed to vibrate, the value of the calculated local drag coefficient C_D varied from 0.654 to 1.261 and that of the inertia coefficient C_M varied from 2.548 to 1.544 as the Keulegan and Carpenter number varied from 1.87 to 10.69 (see Figures 5.3.2a and 5.3.2b). The order of the magnitude of the value of C_D and C_M obtained agree with the previous published results (Hogben et al 1977[53]) (Aagaard and Dean 1969[114] C_D varied from 0.5 - 1.2 and C_M was constant and equal to 1.33, Keulegan and Carpenter 1958[27] C_D varied from 0.7 to 1.3 and C_M varied from 2.6 to 1.6, Morison et al 1950[25] C_D varied from 1.626 ± 0.414 and C_M varied from 1.508 ± 0.197 and Paape and Breusers 1967[115] C_D varied from 0.5 to 2.0 and C_M varied from 1.0 to 2.5).

6.2.3 When the structures were free at the top i.e. free to vibrate, the values of C_D and C_M were calculated by two methods, first by Morison's equation and second by modified Morison's equation. The second

method producing the larger value of the drag coefficient with a maximum ratio (C_D modified Morison/C_D Morison) of 102.7% while the inertia coefficient C_M showed no differences between the above two methods (see Figures 5.3.15a and 5.3.15b).

6.2.4 There is a difference between the magnitude of the drag coefficient of the "free" structure and that of the "fixed" structure under a particular set of wave conditions, the former being greater in magnitude than the latter. The maximum ratio of the drag coefficient for the free structure to that of fixed structure (C_D Free/C_D Fixed) was 118.45% (see Figure 5.3.13). It was found that the inertia coefficient exhibited a similar pattern of behaviour. The maximum ratio of the inertia coefficient for free structure to that of fixed structure (C_M Free/C_M Fixed) was 109.56% (see Figure 5.3.14). These differences in the magnitude of the force coefficients decrease with increasing depth from the still water level (see Tables 5.4 to 5.11).

The difference in magnitude of the drag coefficients is greater than the corresponding difference in magnitude of the inertia coefficients. Both differences in value are attributed to the vibration of the structure and consequently the

change in the flow field round the structure (see Figures 5.2.3 and 5.2.4). Therefore these differences are dependent upon the structural response (X/D) and Keulegan-Carpenter number (Um T/D). There is a decrease in magnitude with increasing Keulegan-Carpenter number for the same structure because the increase of the level of vorticity in the fluid field near the structure. For constant Keulegan-Carpenter number their magnitude increase with decreasing structural stiffness.

No differences in the ratio of C_D Free/C_D Fixed and the ratio of C_M Free/C_M Fixed were noticed when the structural response (X/D) was less than 0.006. As the structural response increases to a value in excess of 0.006 the differences in coefficients start to increase sharply until a peak is reached at X/D of 0.01, thereafter the differences decrease with increasing structural response as shown in Table (5.12).

For the tested structures and waves, the wave force acting on the free structure was greater than that of the fixed structure. The maximum difference was 9.5%.

6.2.5 The magnitude of the relative damping coefficient is greater in water than in air and increases with increasing water depth. For the same structure, the magnitude of the relative damping coefficient increases with increasing added mass at the top of the structure i.e. decreasing structure stiffness (see Table (5.1)).

6.2.6 Also from Table (5.1) it can be seen that the natural frequency of the structure in air had a higher value than that in water. The value of the structural natural frequency in water decreases with increasing water depth.

6.2.7 The dynamic analysis using the first three modes in the mode superposition technique (as suggested by Molhatra and Penzien 1970[85]) and using the damping matrix as proportional to the mass and stiffness matrix, gives a more accurate representation of the problem of the structural response (deflection) for a cantilever structure when the ratio of the forcing frequency to the structure's frequency is very small (<< 1) (i.e. 0.13). The maximum difference between the measured and the calculated structural response using the modified Morison's equation (Method A) is 9.87%. The maximum difference between the measured and calculated structural response using Morison's equation (Method B) is 7.45% (see Tables 5.12 to 5.21).

6.2.81 The maximum difference between the calculated structural response using Mechod A and Method B is 9.56%.

6.2.9 The structural response (X/D) is better represented by the non-dimensional response parameter $(2\bar{M}\xi/\rho D^2 L_s)$ $(D\ FN/U_m)$ than by the reciprocal of the reduced velocity $(D\ FN/U_m)$ as shown in Figures(5.3.16a - 5.3.16b).

6.3 RECOMMENDATION

Additional research is needed, under the same condition (FF/FN << 1) for structures with dynamic properties differing from those investigated herein. In addition a continuation of research is needed under different condition (FF/FN > 1) and less than the second resonance for structure of similar dynamic properties to those investigated.

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A P P E N D I X A

PRESSURE TRANSDUCER

SENSING ELEMENT

The sensing element was a circular disc of 48 mm diameter which was a part of a brass plate. The brass plate was 100 mm square and 0.02 mm thick. It had to be absolutely flat and without initial strain. The plate was dipped in acid etch until its thickness was reduced to between 0.015 to 0.0125 mm. The plate was pulled from each corner to give initial tension and remained in this condition until the pressure transducer was completed. Four strain gauges were fixed at each side of the sensing element. They were FL-6-11 Tokyo Kenkyoio wireless type. Connecting wires had been soldered of 36 standard wire gauge enamel self-fluxing. The strain gauges were connected together through four pine holes which can be punched from the top before the strain gauges were mounted. After the connection of the strain gauge the four pine holes were sealed and the sensing element were sprayed with flexible varnish making the thickness of the sensing element increased to 0.025 mm.

Figure (1A) shows the strain gauges fixed on the sensing element and the connection of the strain gauges.

After the pressure transducer was completely built the outer part of the plate was discarded.

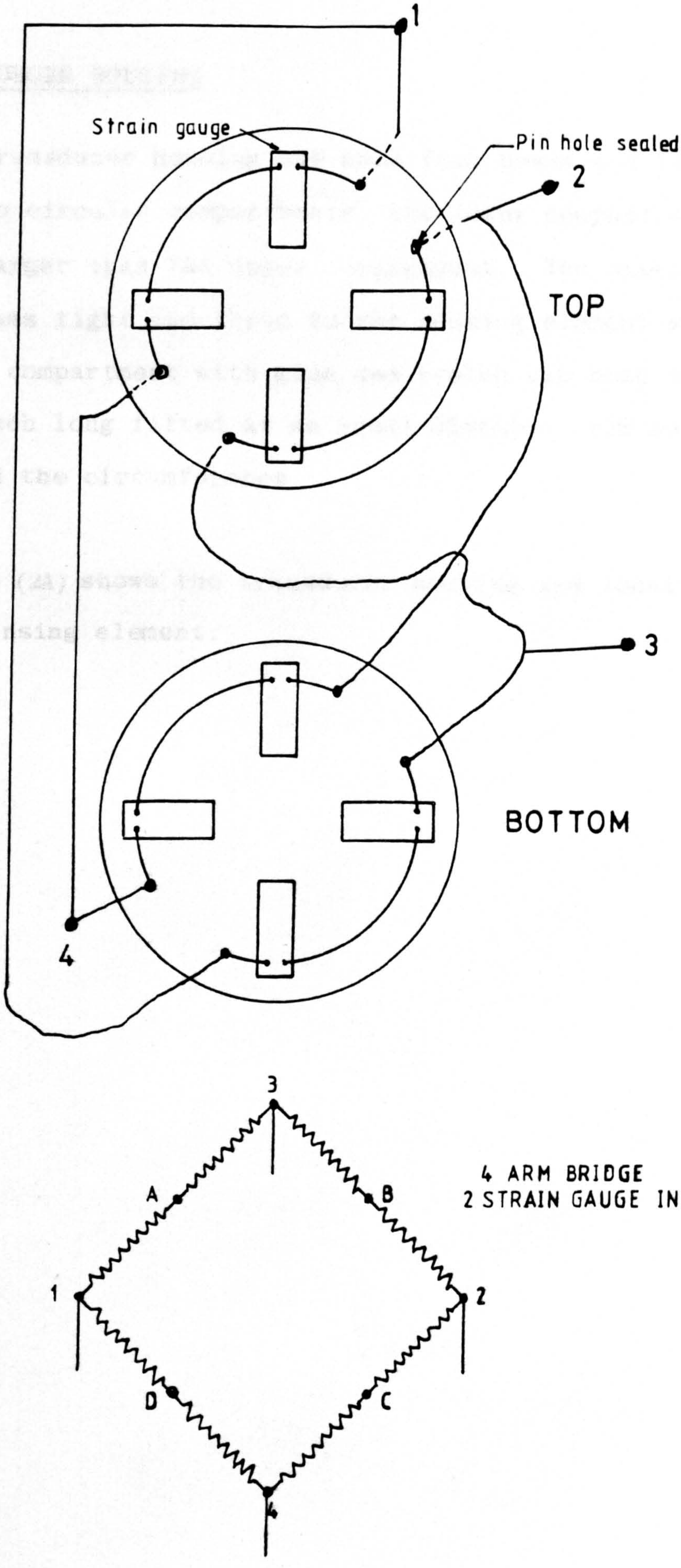


Figure (1A) - The Sensing Element.

TRANSDUCER HOUSING

The transducer housing was made from brass and it consisted of two circular compartments, the lower compartment which was larger than the upper compartment. The upper compartment was tight and fixed to the sensing element and the lower compartment with glue and twelve cap head screws 0.5 inch long fitted at an equal distance from each other around the circumference.

Figure (2A) shows the transducer housing and location of the sensing element.

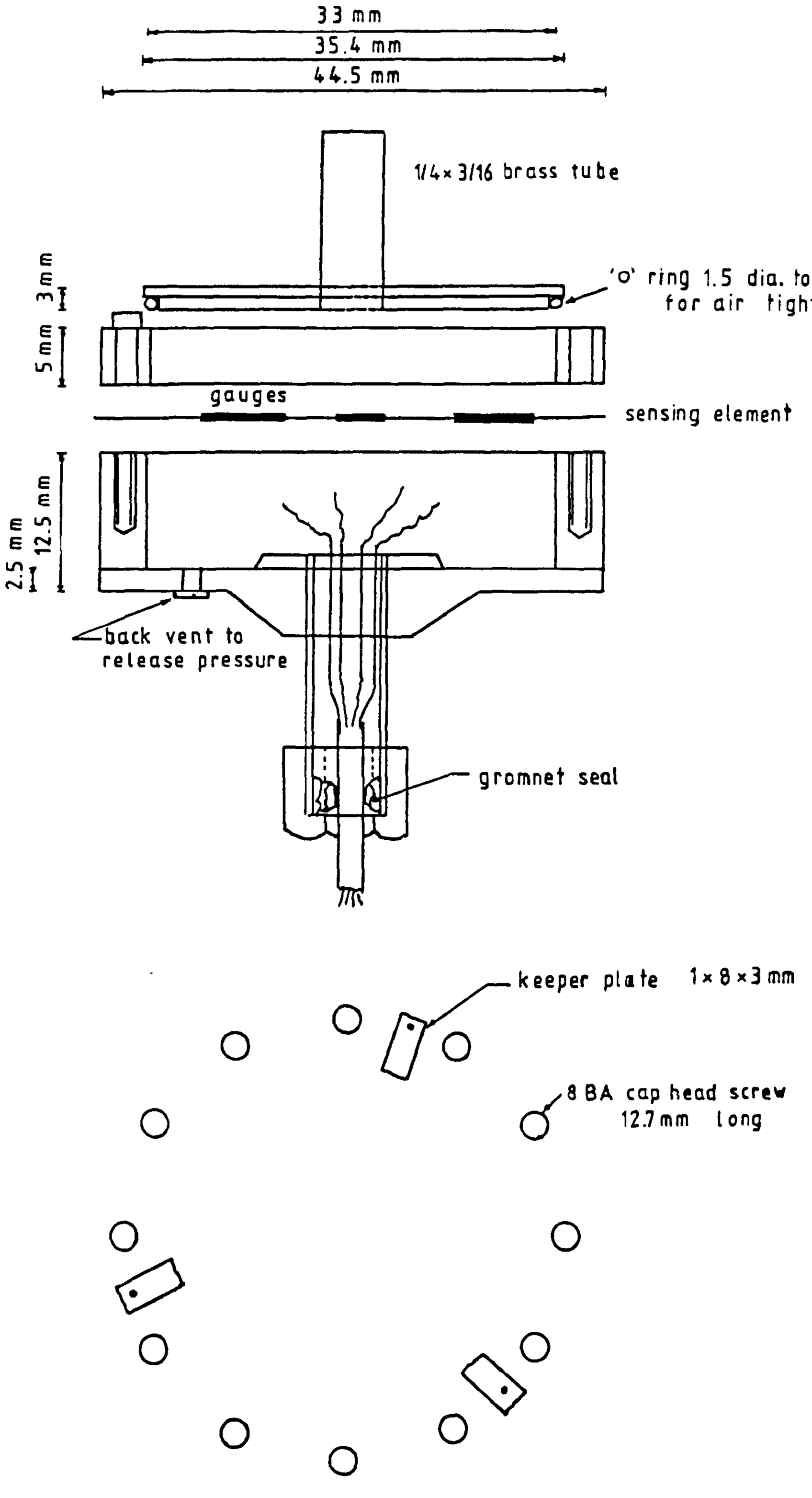


Figure (2A) - The Transducer Housing.

A P P E N D I X B

PREAMPLIFIER

A linear amplifier produces at its output a waveform which is a perfect copy, but of greater amplitude. The most important class of linear amplifier i.c. is the operational amplifier which features high gain, high input resistance, low output resistance and narrow bandwidth extending to d.c.

The 741 is typical of the operational amplifier generally, so that the design methods, circuits and bias arrangements can be used with small modification.

Figure (1B) shows the circuits used for an inverting preamplifier. The components used in this circuit were:

R = potentiometer resistance of 200 kΩ 0·5 Watt

R_1 = 390 kΩ resistance of 0·25 Watt

R_2 = 68 kΩ resistance of 0·25 Watt

R_3 = 470 kΩ resistance of 0·25 Watt

C = Capacitors 0·1 μF

The voltage gain $= \frac{R_3}{R_2} \times \frac{R_3}{R_2}$

$= \frac{470}{68} \times \frac{470}{68} = 47 \cdot 77$

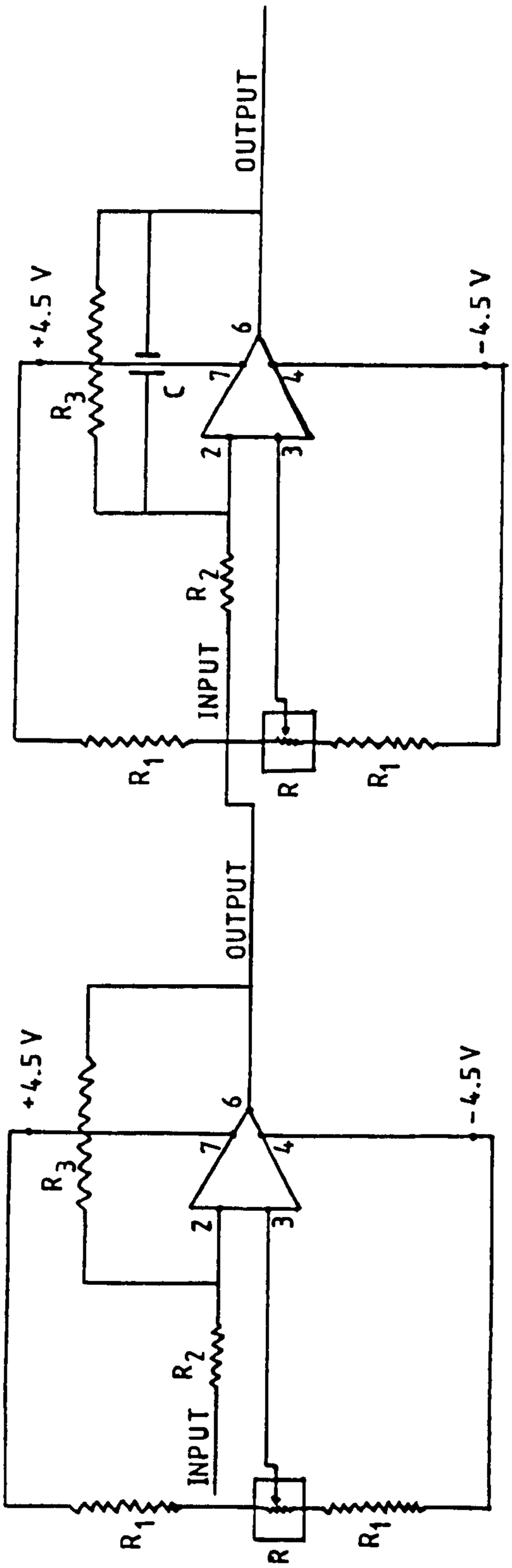
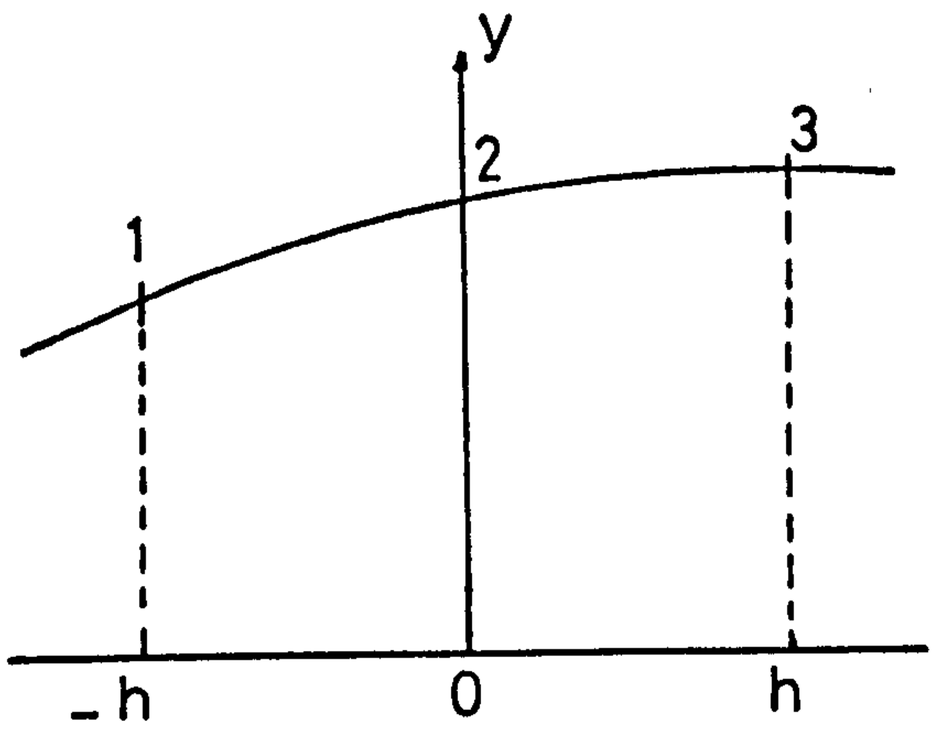


Figure (1B) - The Preamplifier Circuits.

A P P E N D I X C

SINUSOIDAL INTERPOLATION



Let $y = a \sin (bx + c)$

at point 1 $y_1 = a \sin (-bh + c)$

at point 2 $y_2 = a \sin c$

at point 3 $y_3 = a \sin (bh + c)$

$$\therefore y_1 + y_3 = a \sin (-bh + c) + a \sin (bh + c)$$
$$= 2a \sin c \cos bh$$
$$= 2y_2 \cos bh$$

$$\therefore b = \frac{1}{h} \cos^{-1} \left[\frac{y_1 + y_3}{2y_2}\right]$$

$$y_3 - y_1 = a \sin (bh + c) - a \sin (-bh + c)$$
$$= 2a \cos c \sin bh$$
$$= 2y_2 \cot c \left[1 - \left(\frac{y_1 + y_3}{2y_2}\right)^2\right]^{1/2}$$

$$\therefore\ C = \tan^{-1}\left[\frac{2y_2}{y_3 - y_1}\ \left\{1 - \left(\frac{y_1 + y_3}{2y_2}\right)^2\right\}^{1/2}\right]$$

and

$$a = \frac{y_2}{\sin c}$$

care should be taken to ensure that $y_2 \neq 0$

A P P E N D I X D

THE VARIATION OF THE RELATIVE DISPLACEMENTS X/D WITH TIME FOR ONE CYCLE OF THE WAVE

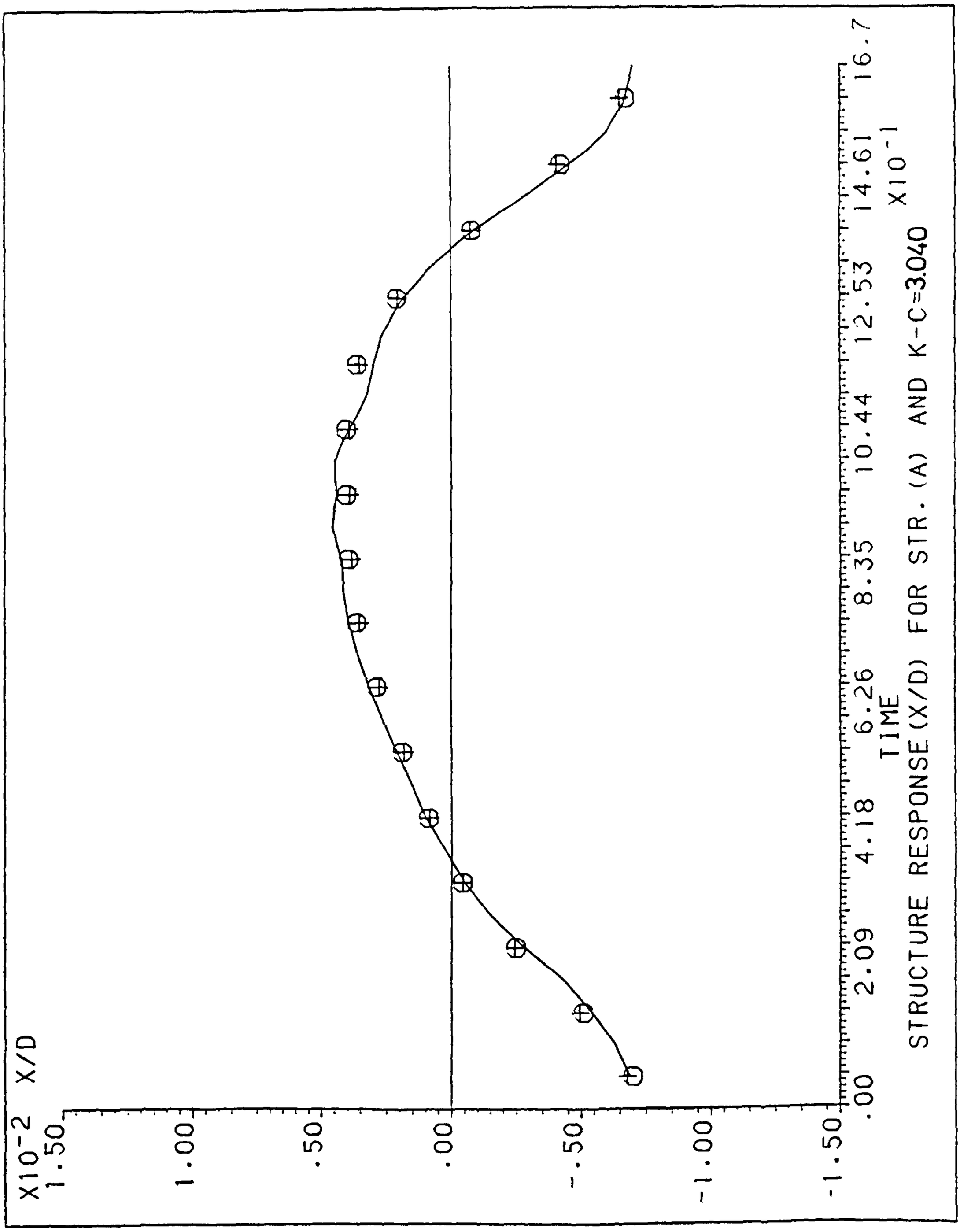


Figure (1D) - Structure's Response (X/D) during the Wave Cycle for Structure A (K-C = 3.04).

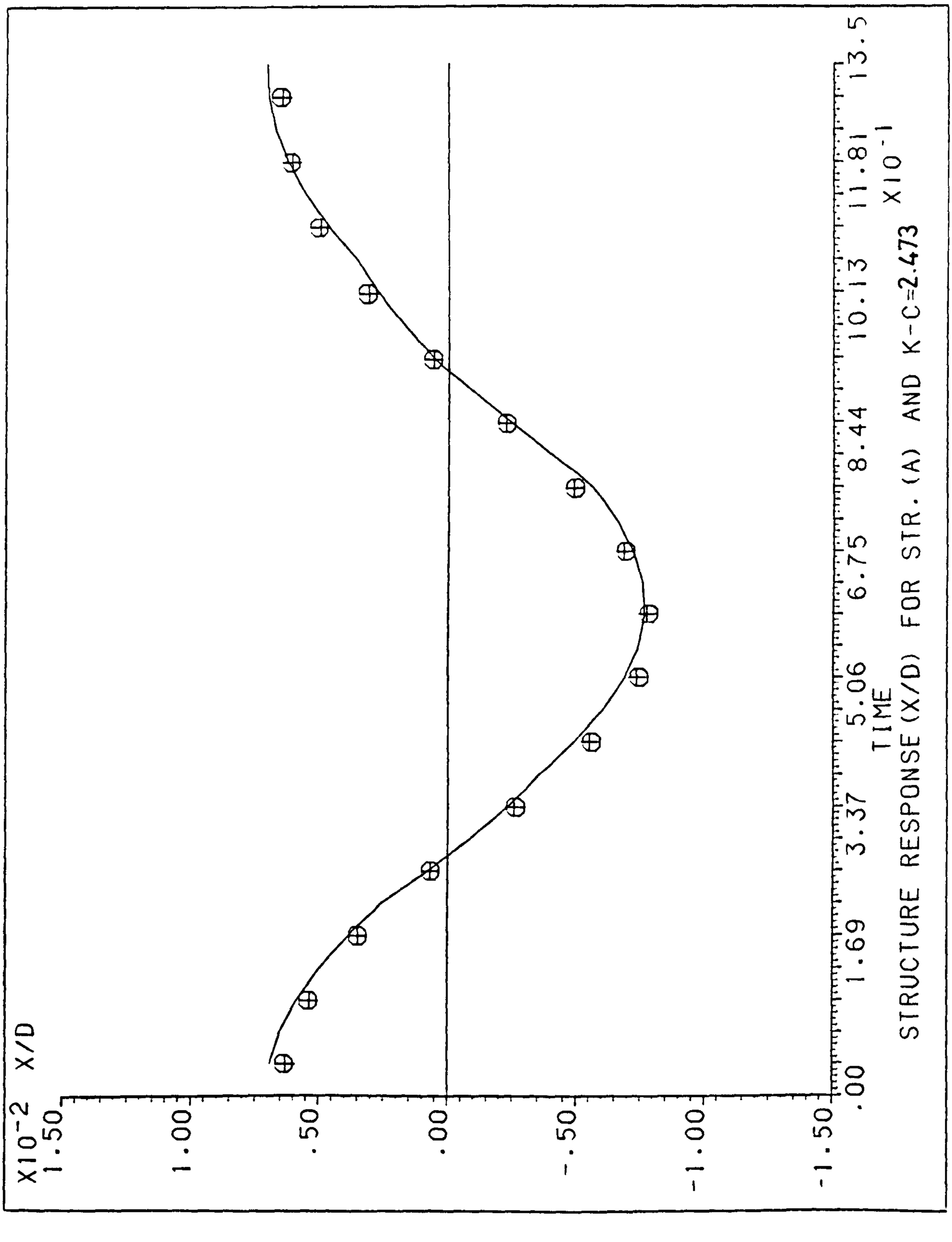


Figure (2D) - Structure's Response (X/D) during the Wave Cycle for Structure A (K-C = 2.473).

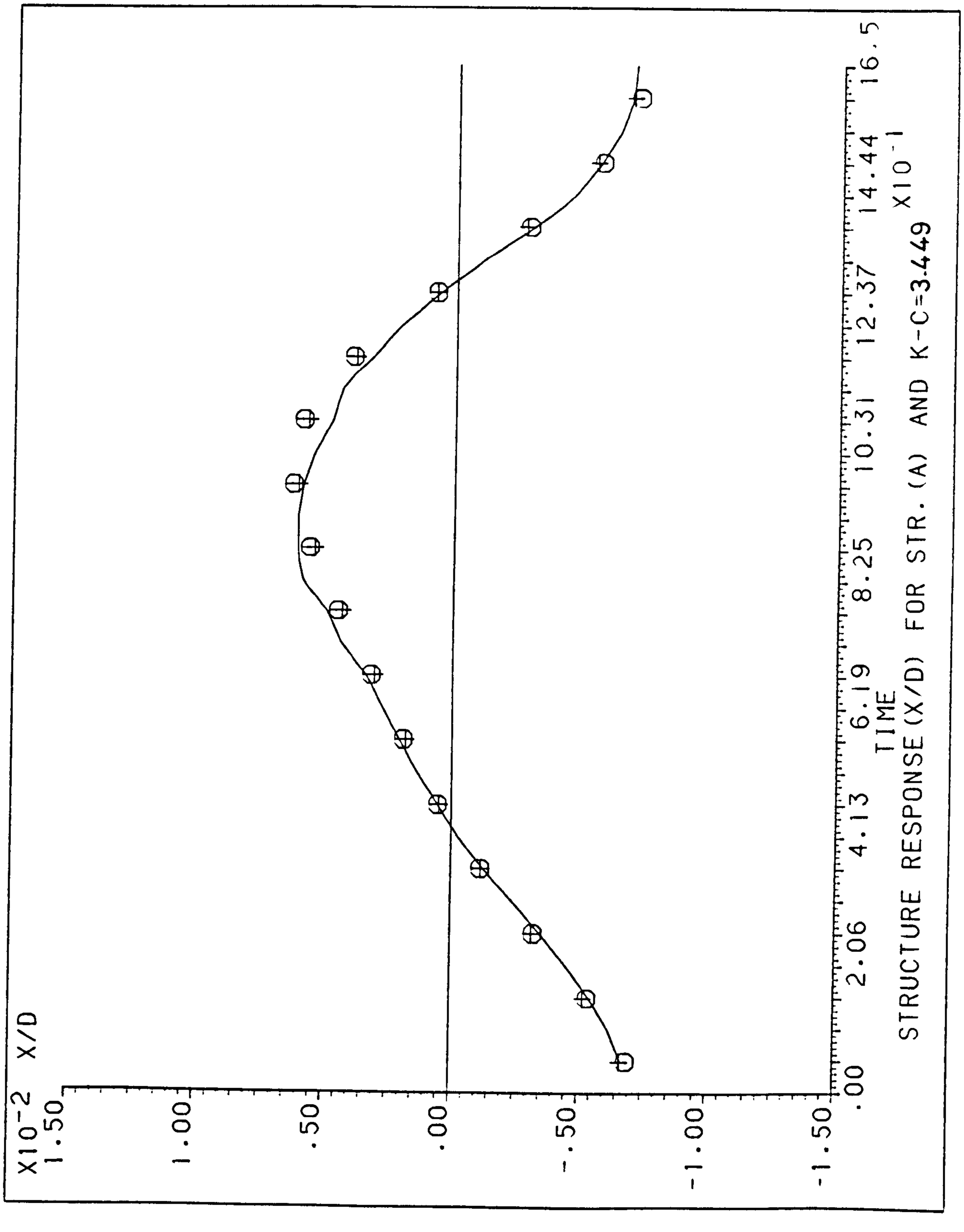


Figure (3D) - Structure's Response (X/D) during the Wave Cycle for Structure A (K-C = 3.449).

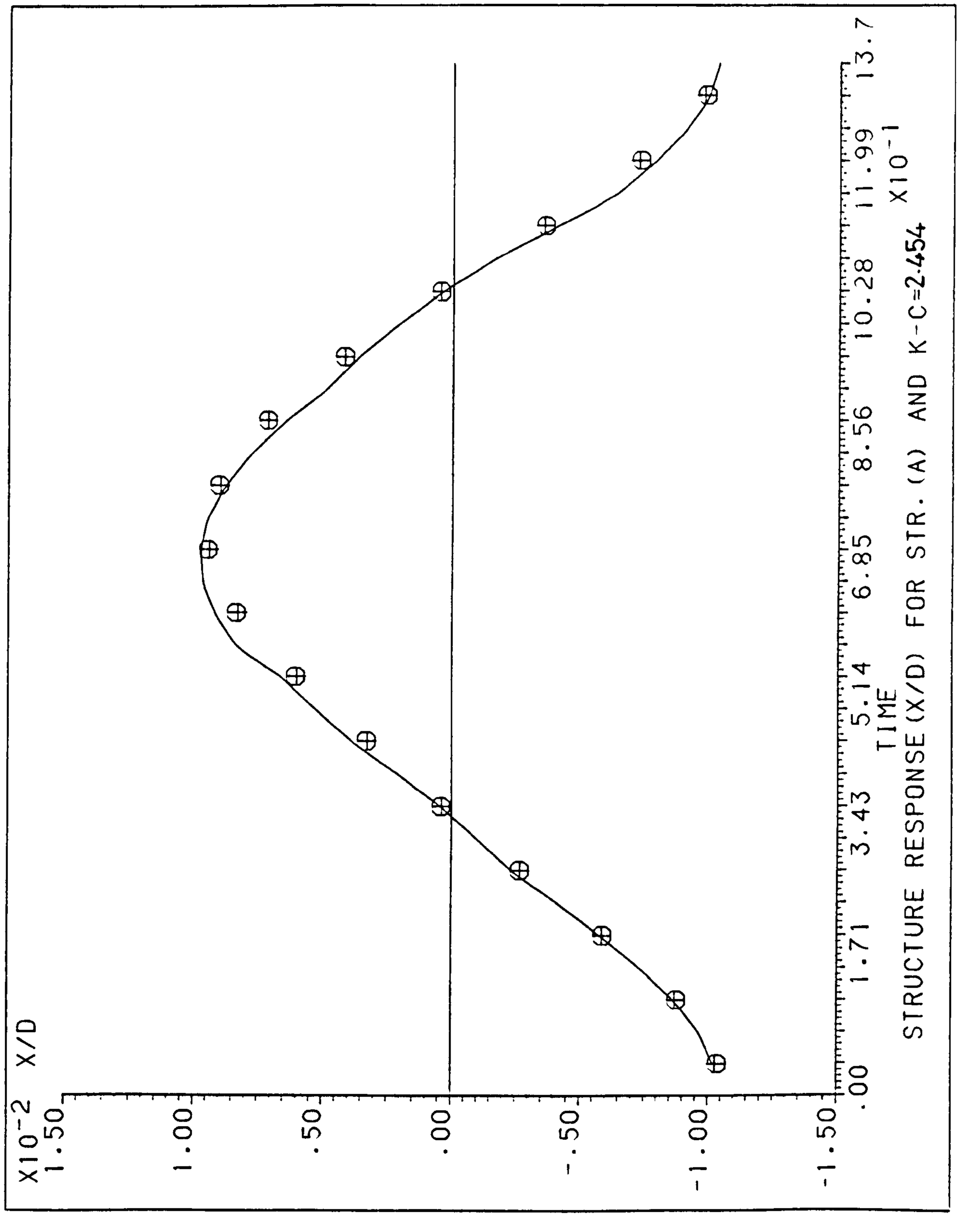


Figure (4D) - Structure's Response (X/D) during the Wave Cycle for Structure A (K-C = 2.454).

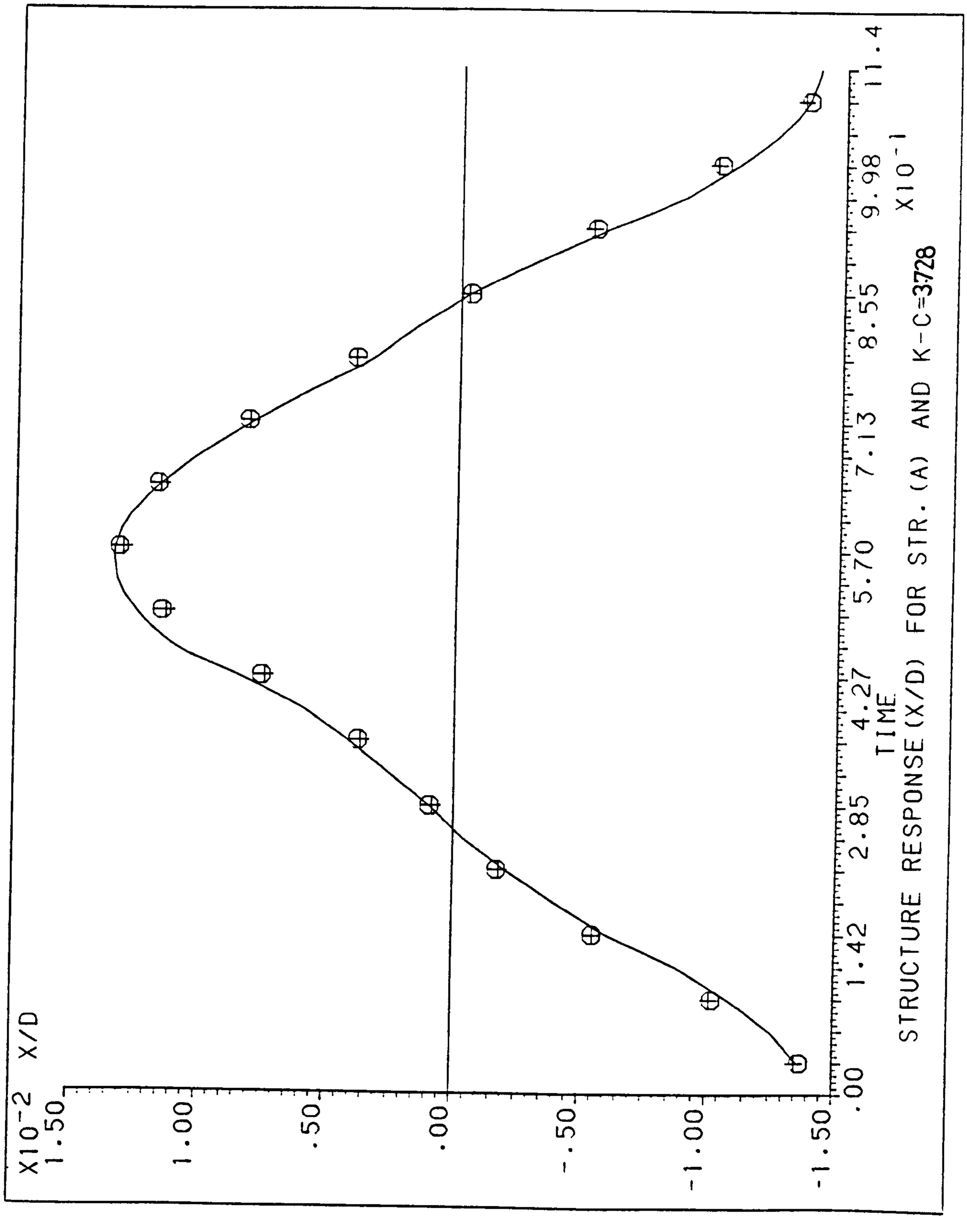


Figure (5D) - Structure's Response (X/D) during the Wave Cycle for Structure A (K-C = 3.728).

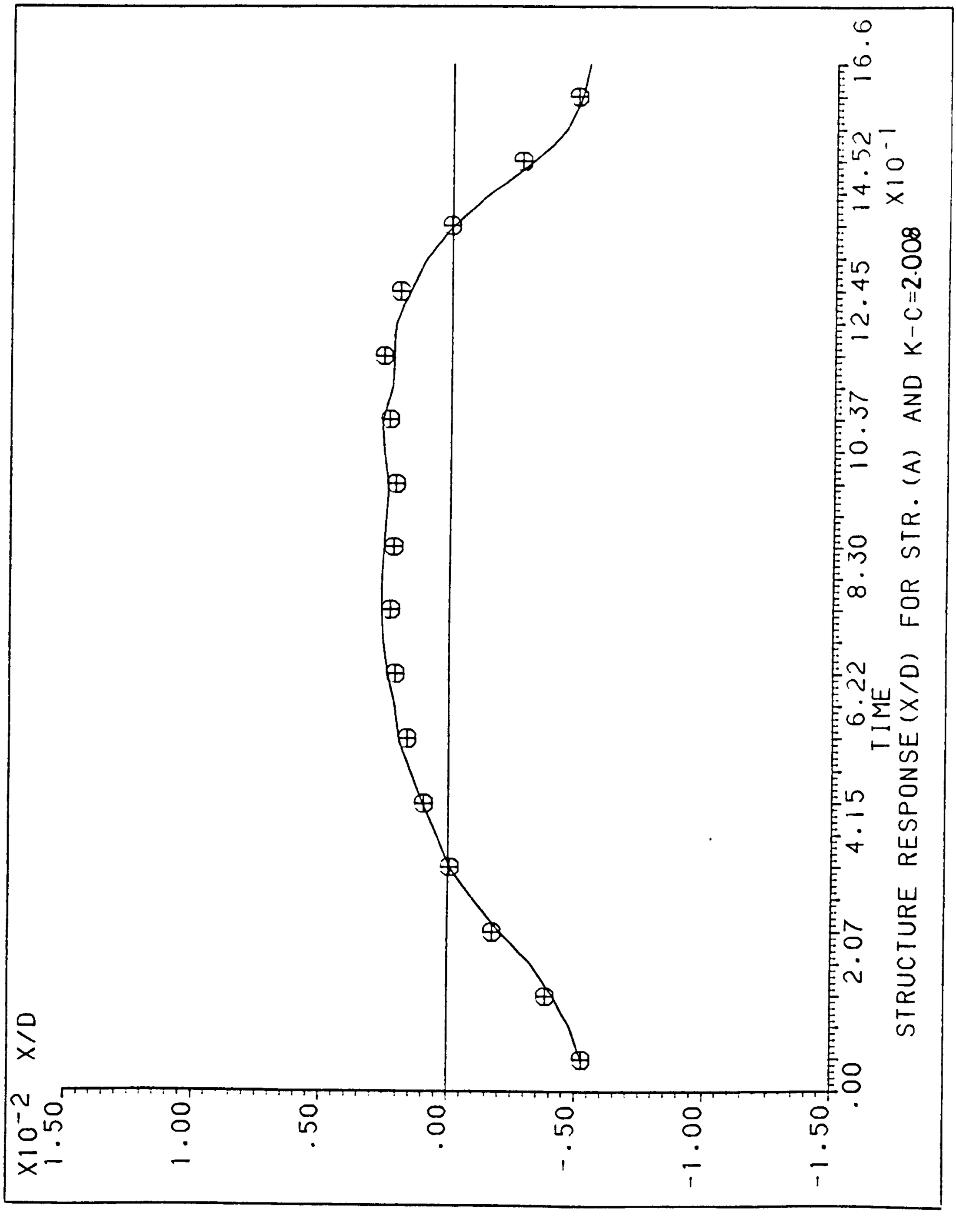


Figure (6D) - Structure's Response (X/D) during the Wave Cycle for Structure A (K-C = 2.008).

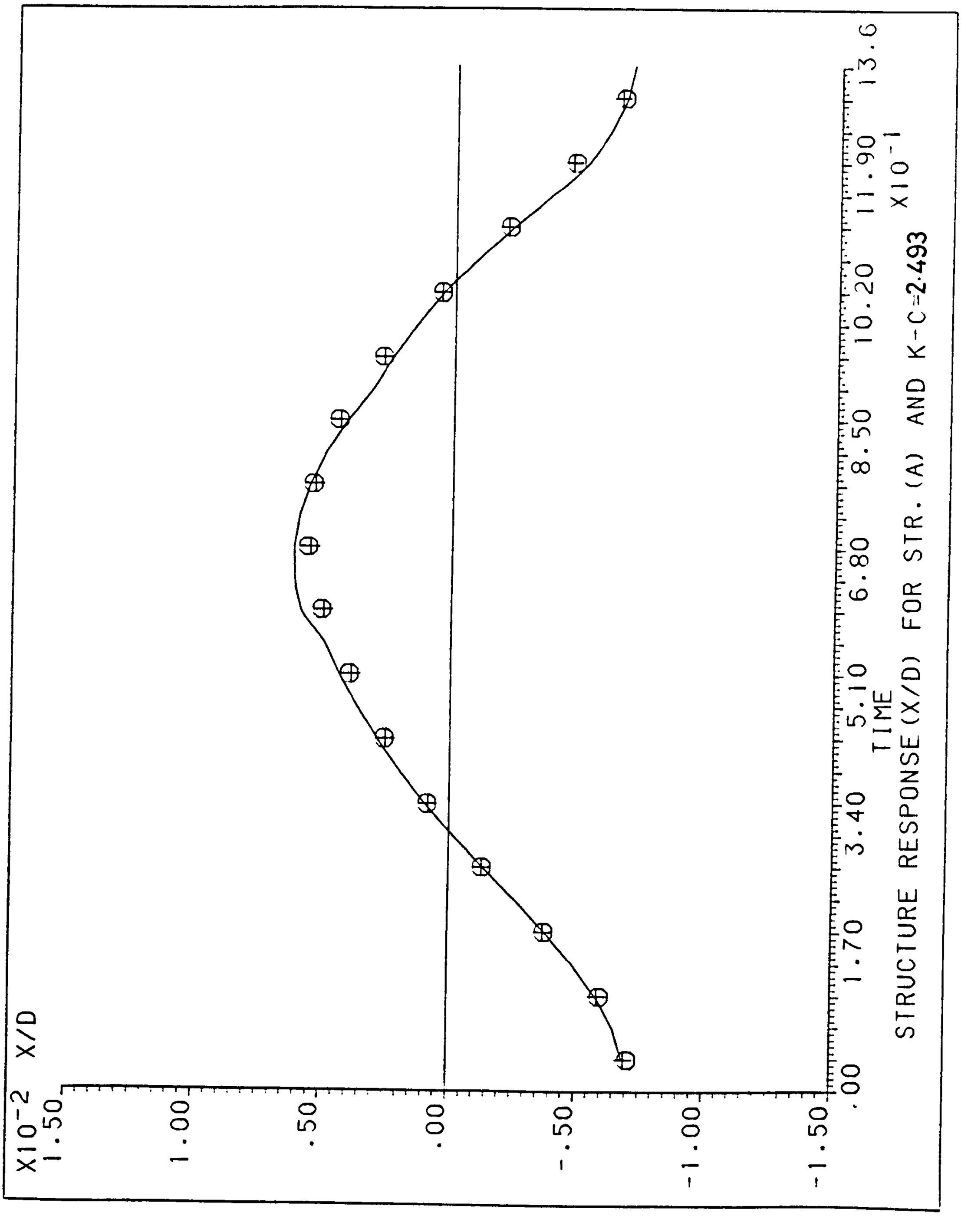


Figure (7D) - Structure's Response (X/D) during the Wave Cycle for Structure A (K-C = 2.493).

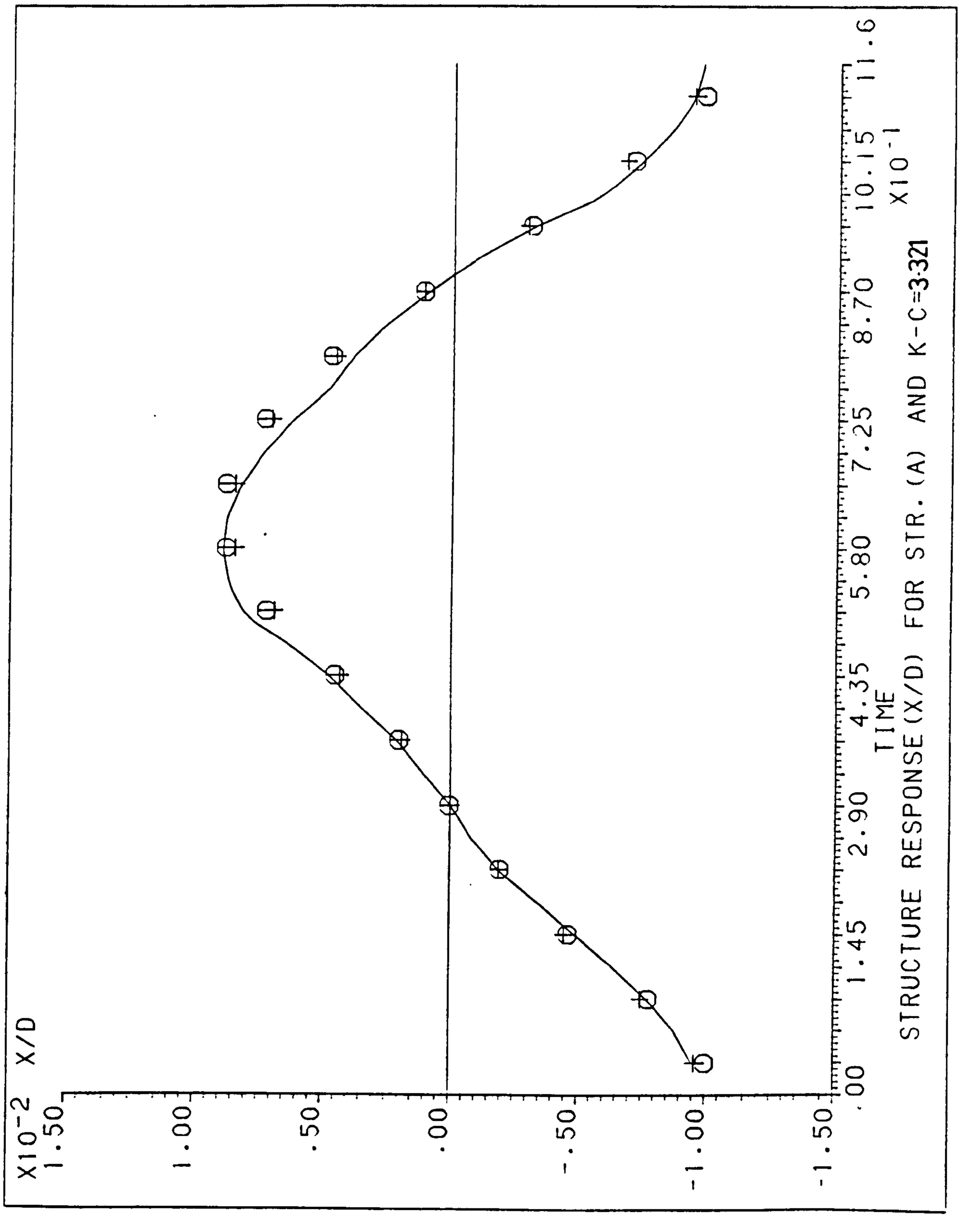


Figure (8D) - Structure's Response (X/D) during the Wave Cycle for Structure A (K-C = 3.321).

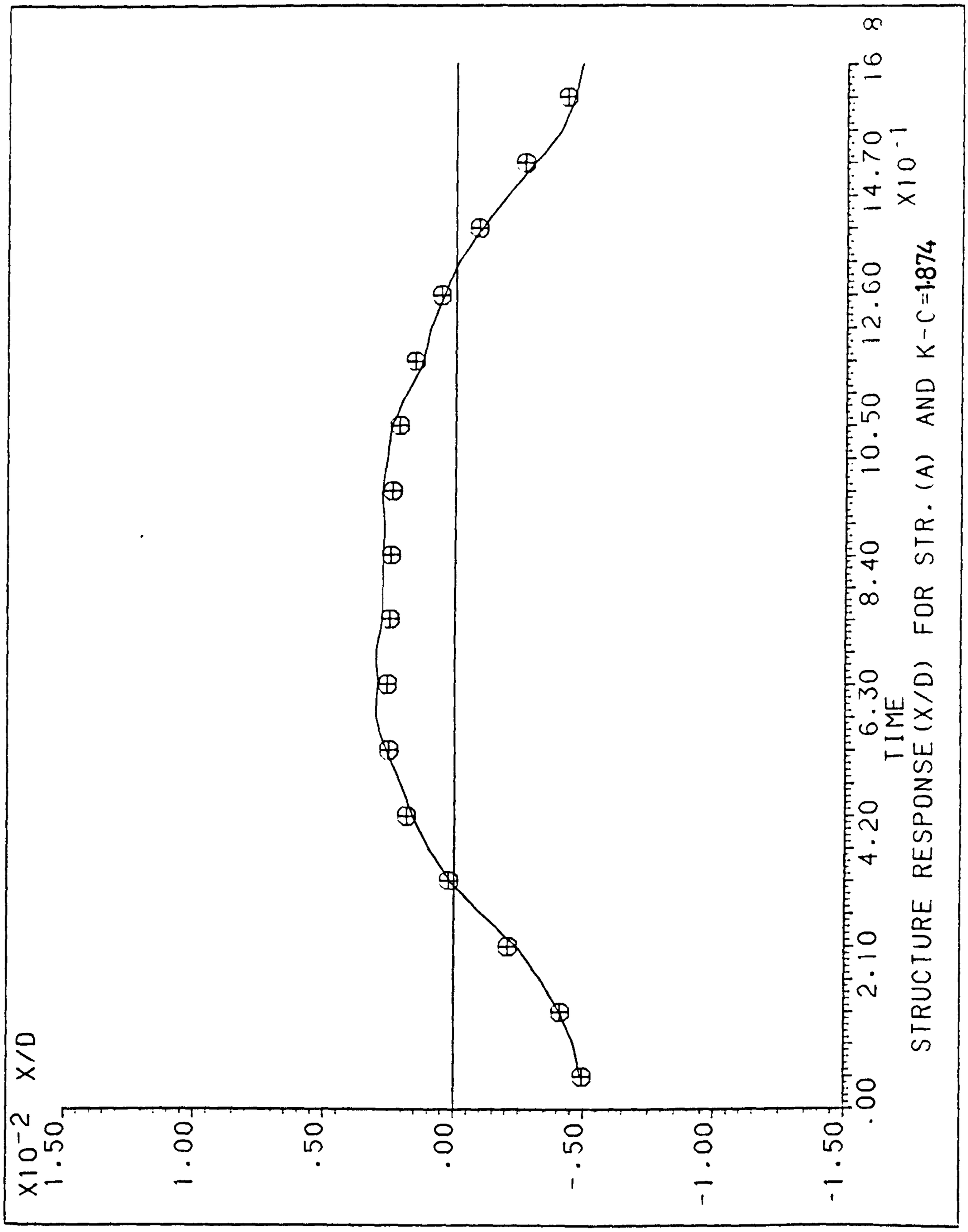


Figure (9D) - Structure's Response (X/D) during the Wave Cycle for Structure A (K-C = 1.874).

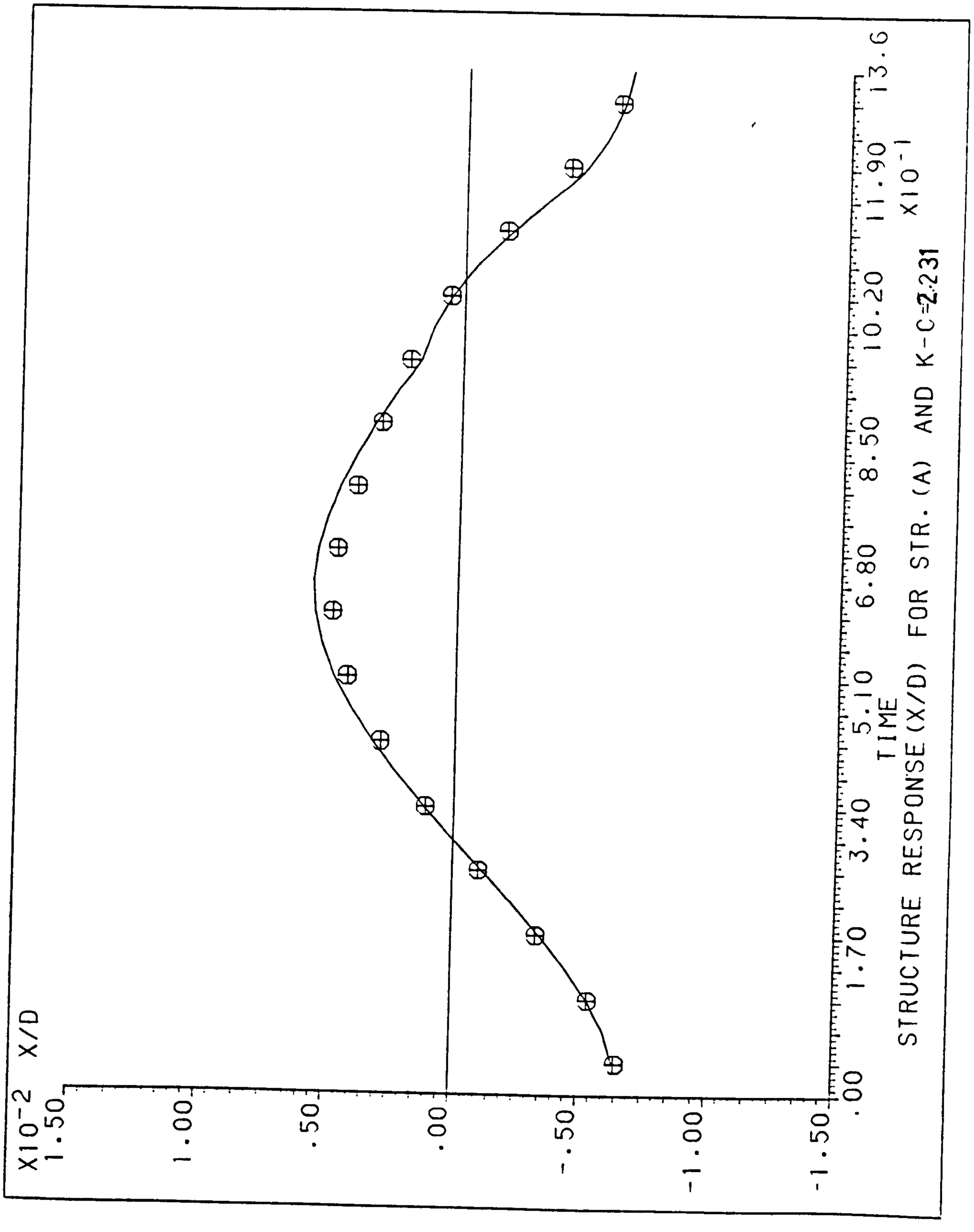


Figure (10D) - Structure's Response (X/D) during the Wave Cycle for Structure A (K-C = 2.231).

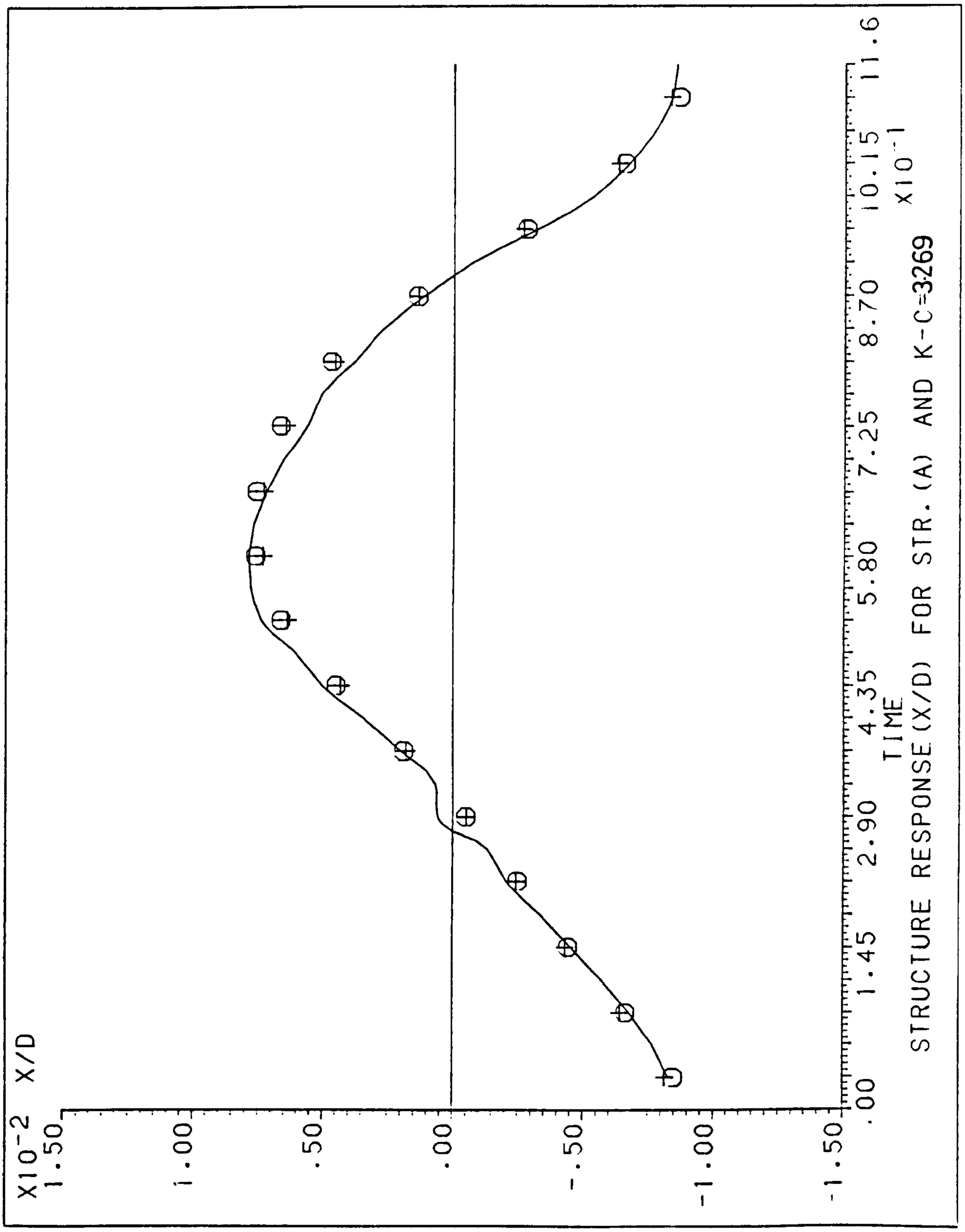


Figure (11D) - Structure's Response (X/D) during the Wave Cycle for Structure A (K-C = 3.269).

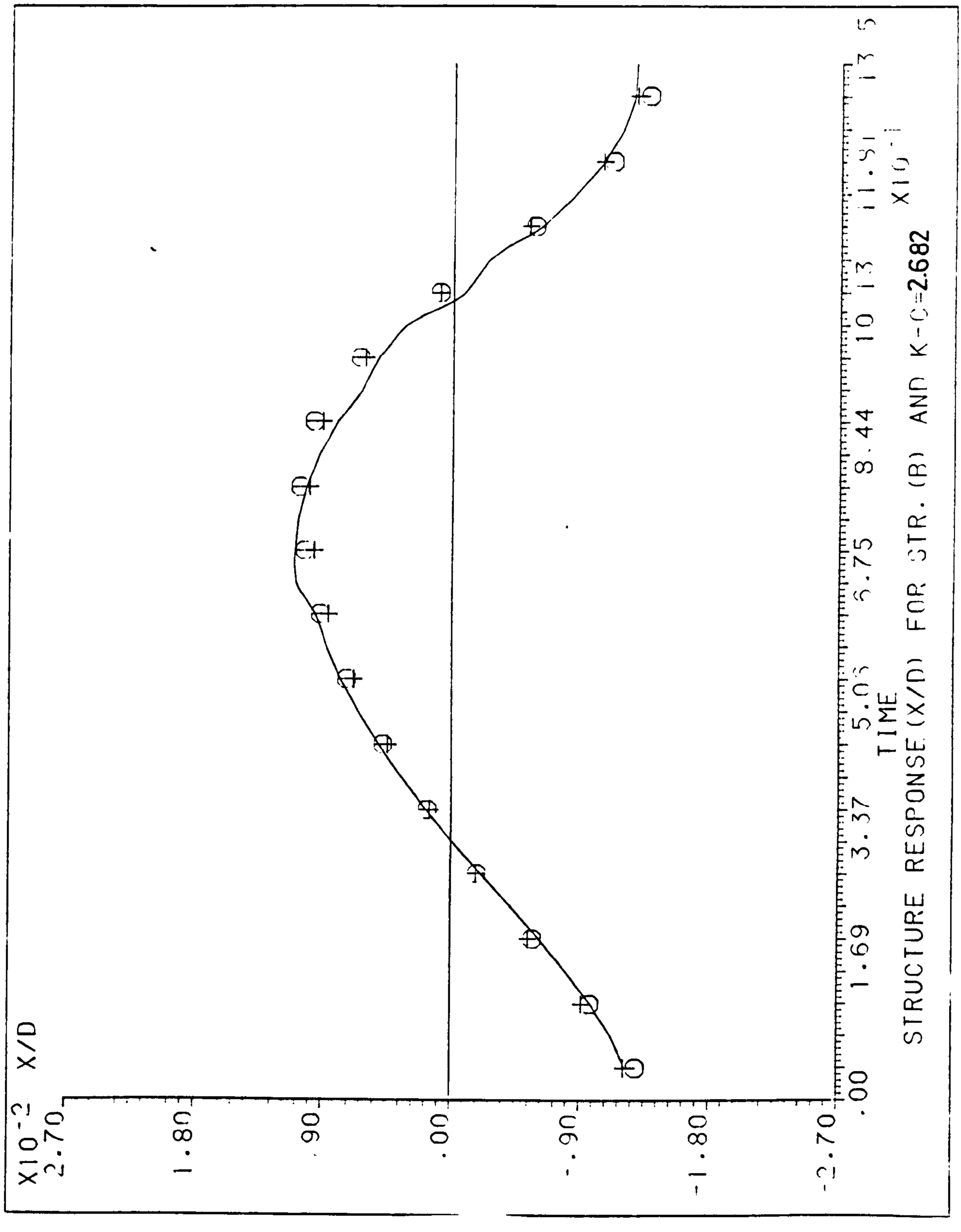


Figure (12D) - Structure's Response (X/D) during the Wave Cycle for Structure B (K-C = 2.682).

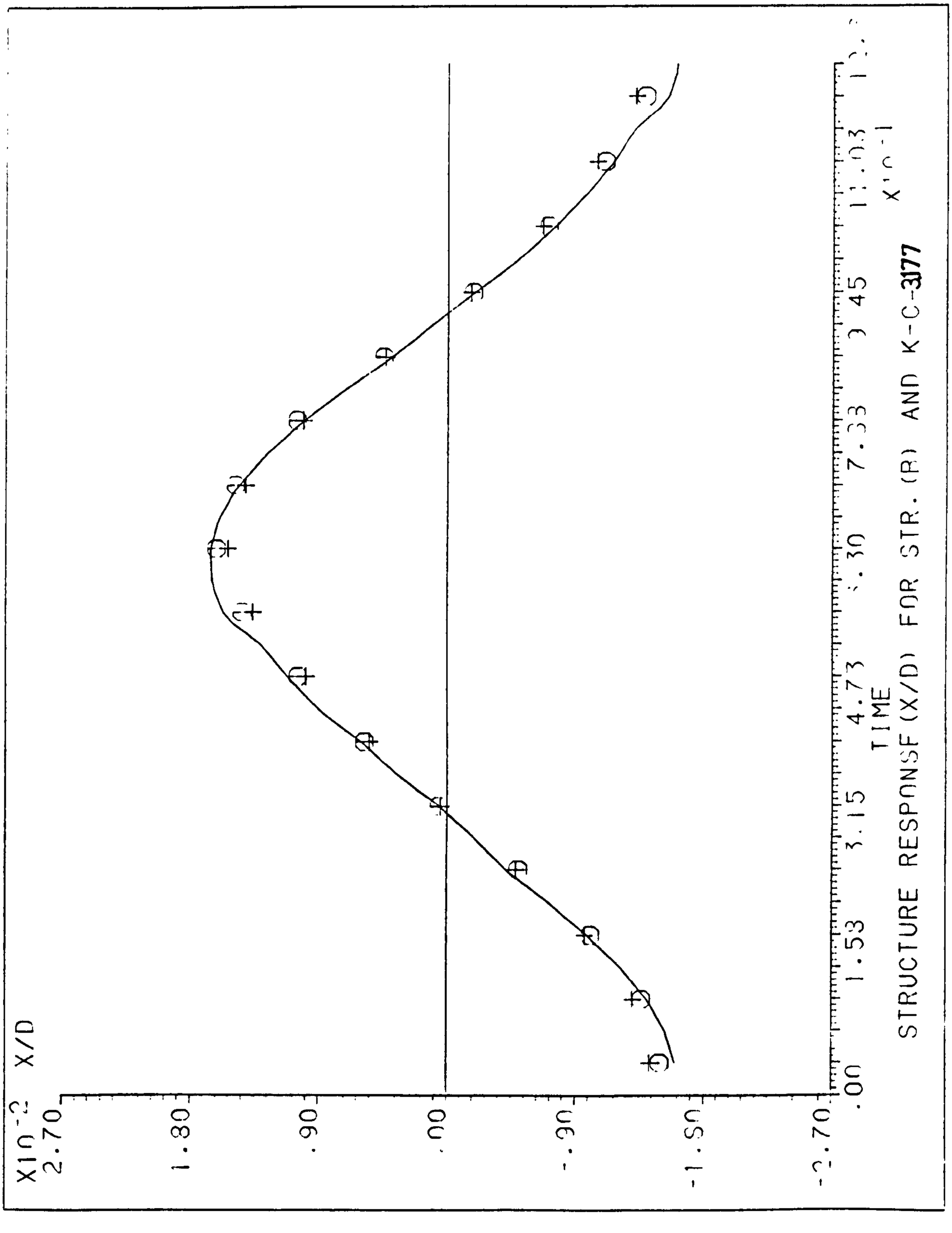


Figure (13D) - Structure's Response (X/D) during the Wave Cycle for Structure B (K-C = 3.177).

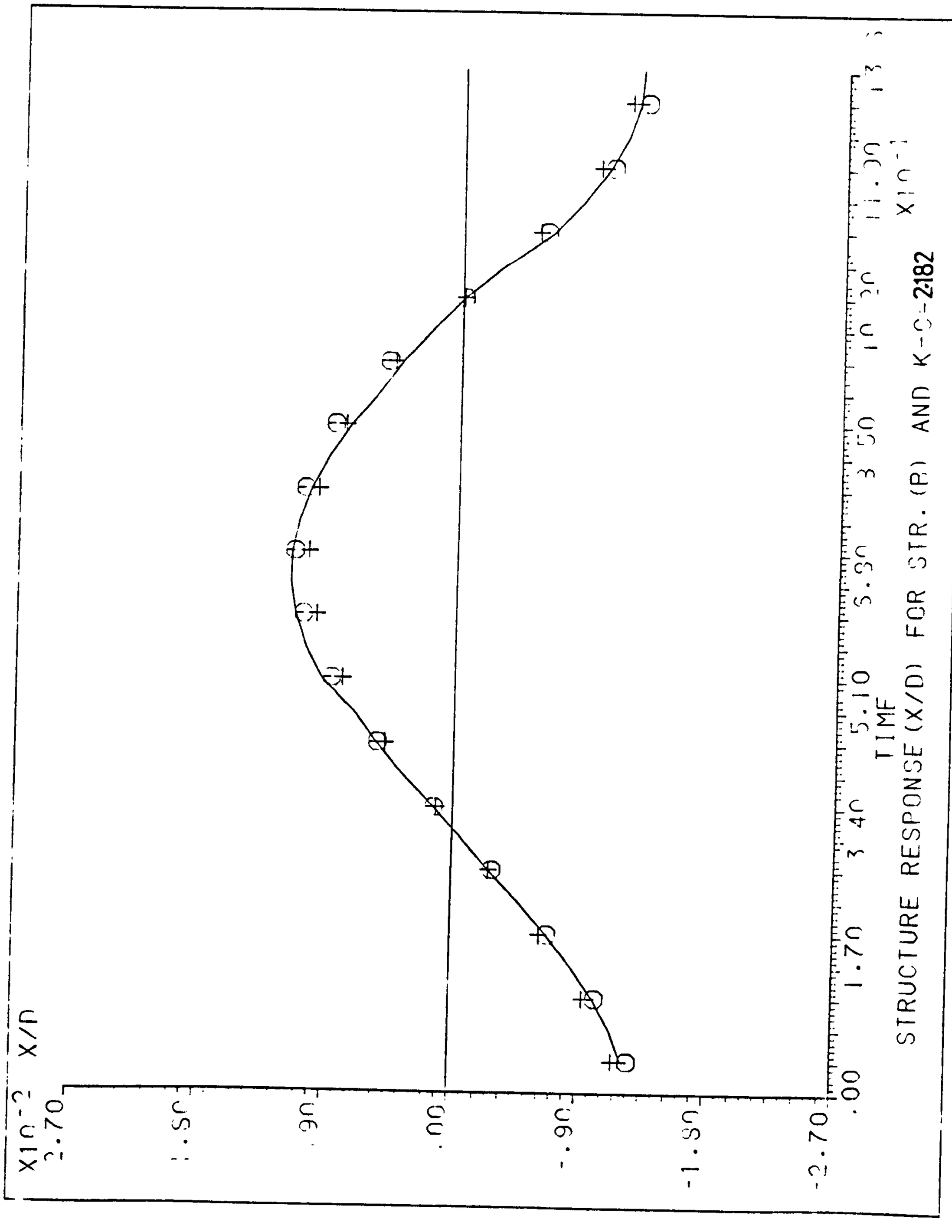


Figure (14D) - Structure's Response (X/D) during the Wave Cycle for Structure B (K-C = 2.182).

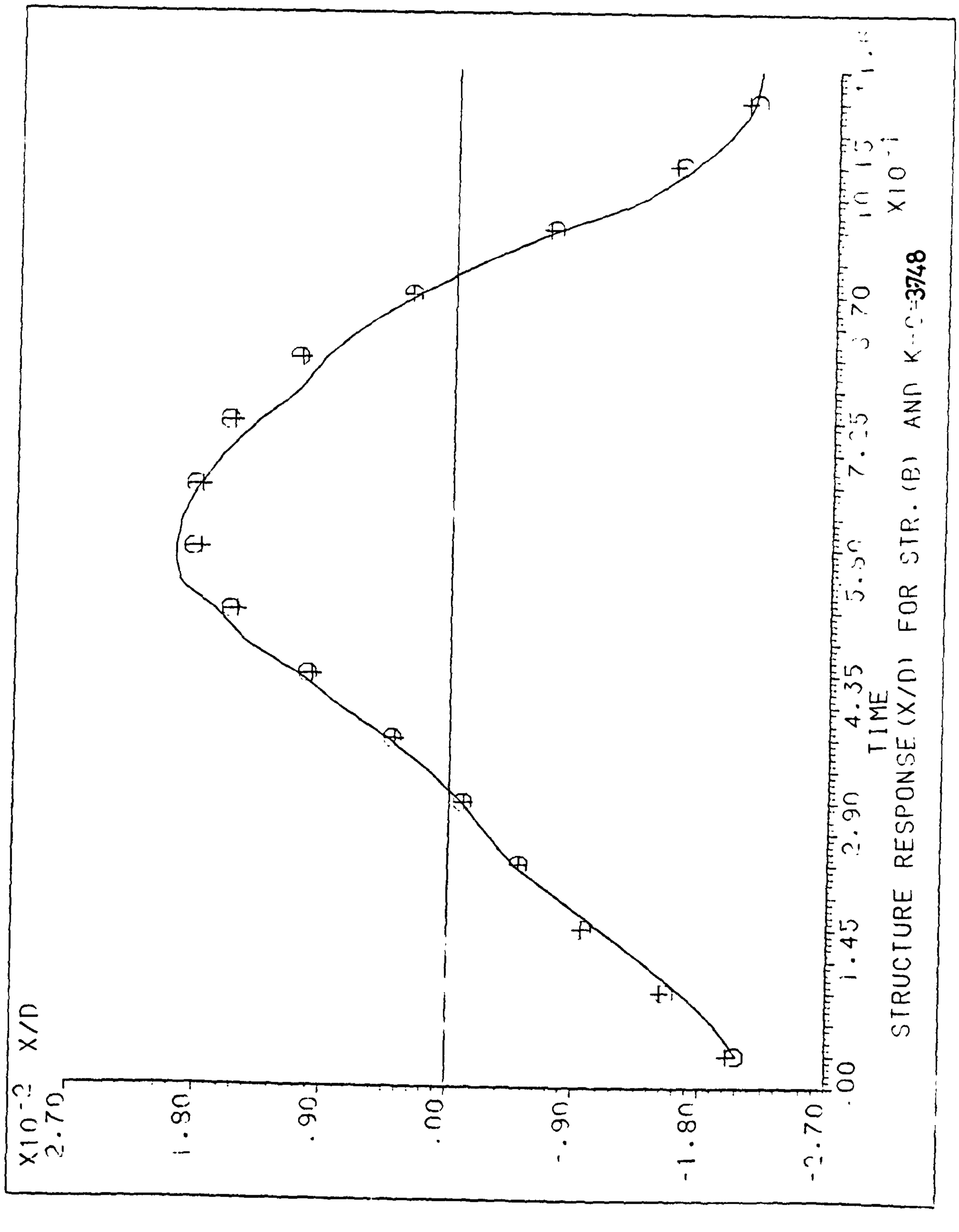


Figure (15D) - Structure's Response (X/D) during the Wave Cycle for Structure B (K-C = 3.748).

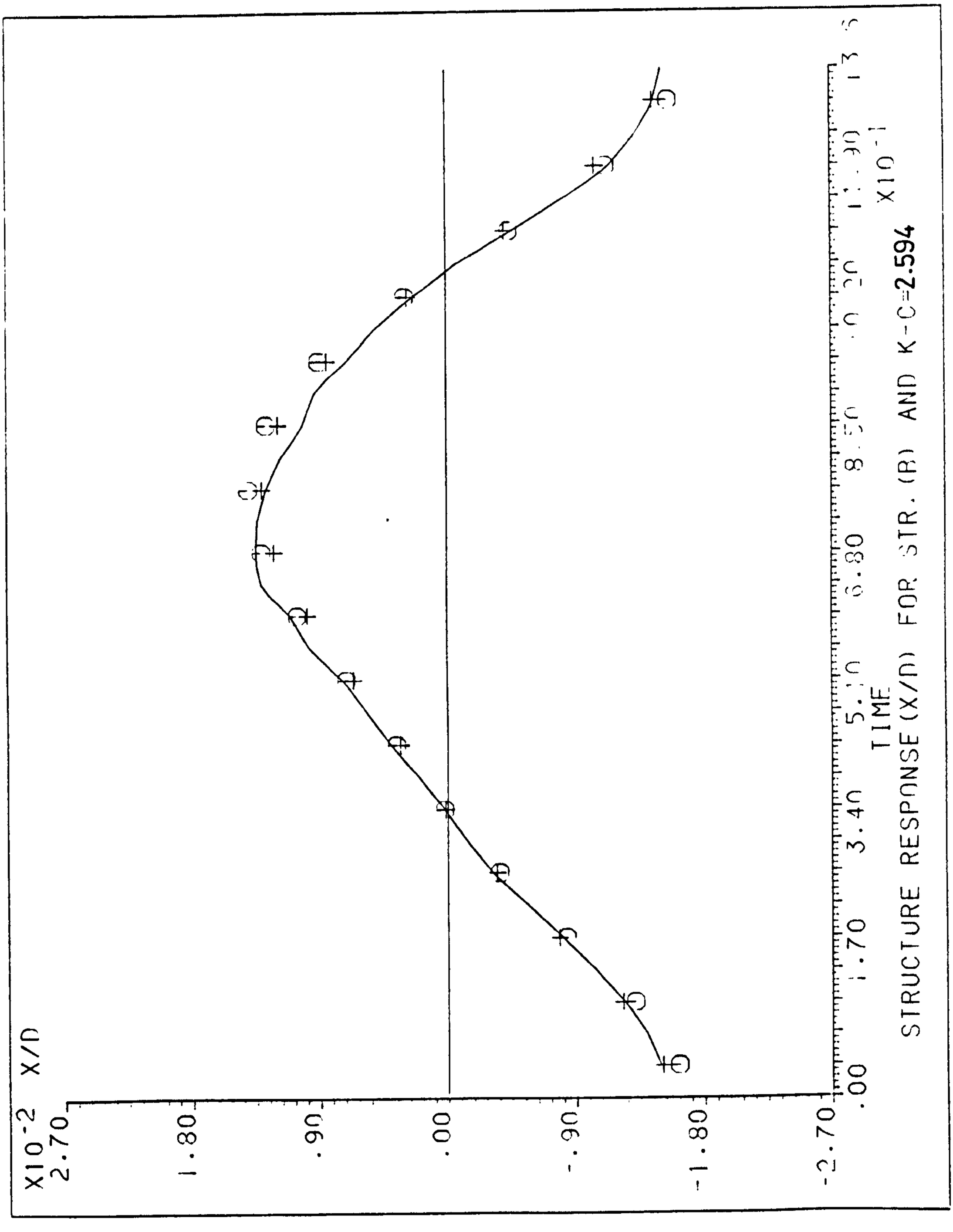


Figure (17D) - Structure's Response (X/D) during the Wave Cycle for Structure B (K-C = 2.594).

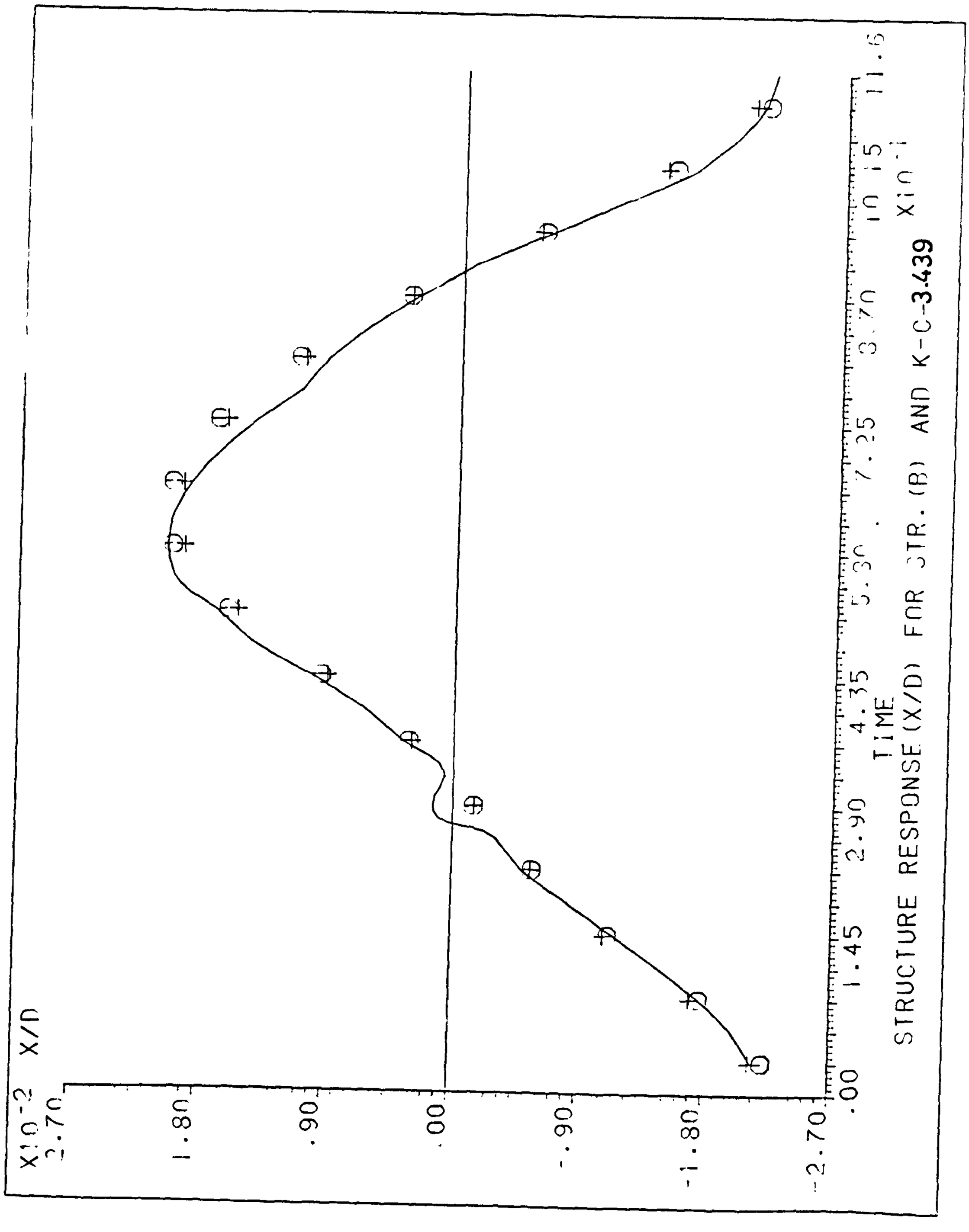


Figure (18D) - Structure's Response (X/D) during the Wave Cycle for Structure B (K-C = 3.439).

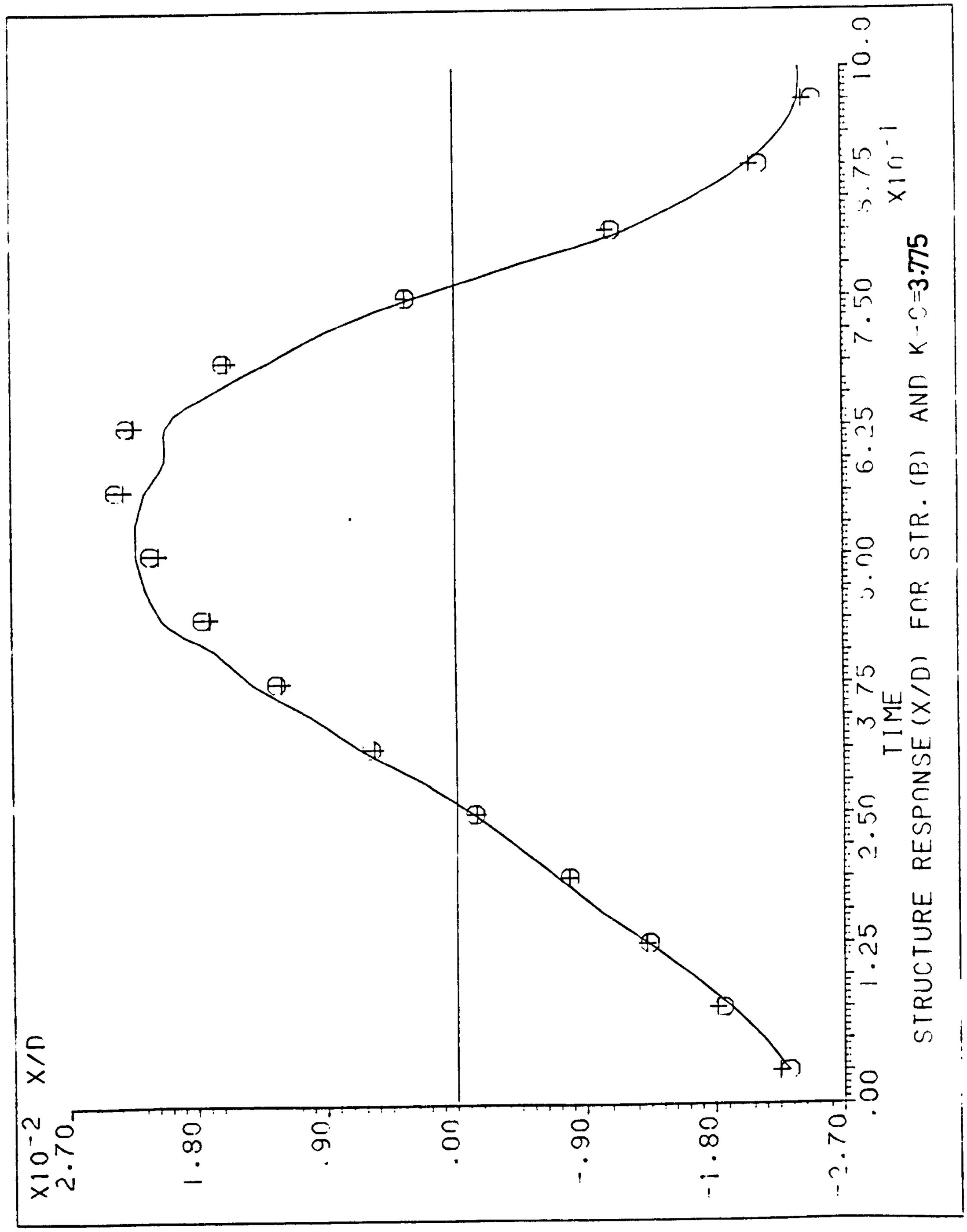


Figure (19D) - Structure's Response (X/D) during the Wave Cycle for Structure B (K-C = 3.775).

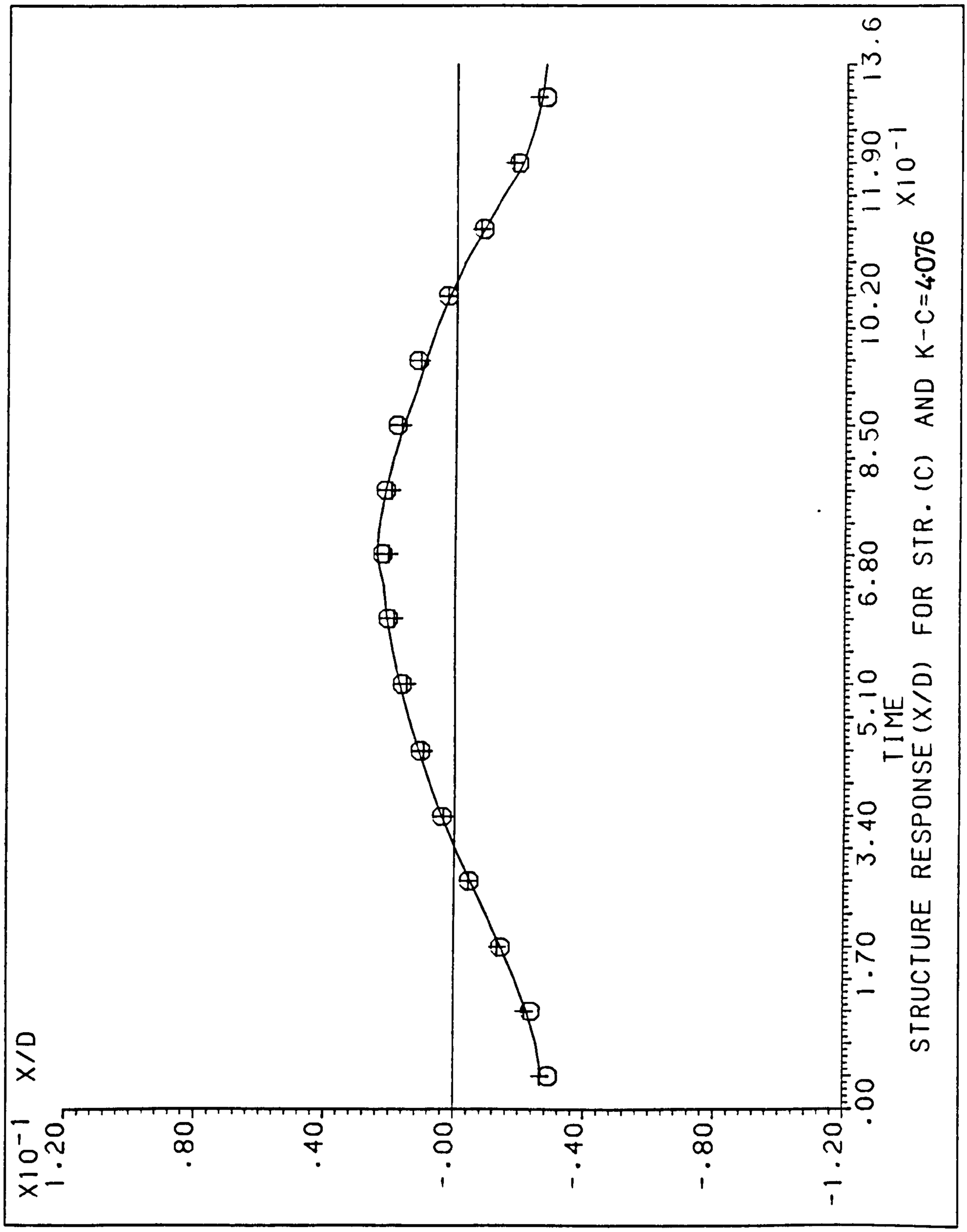


Figure (20D) - Structure's Response (X/D) during the Wave Cycle for Structure C (K-C = 4.076).

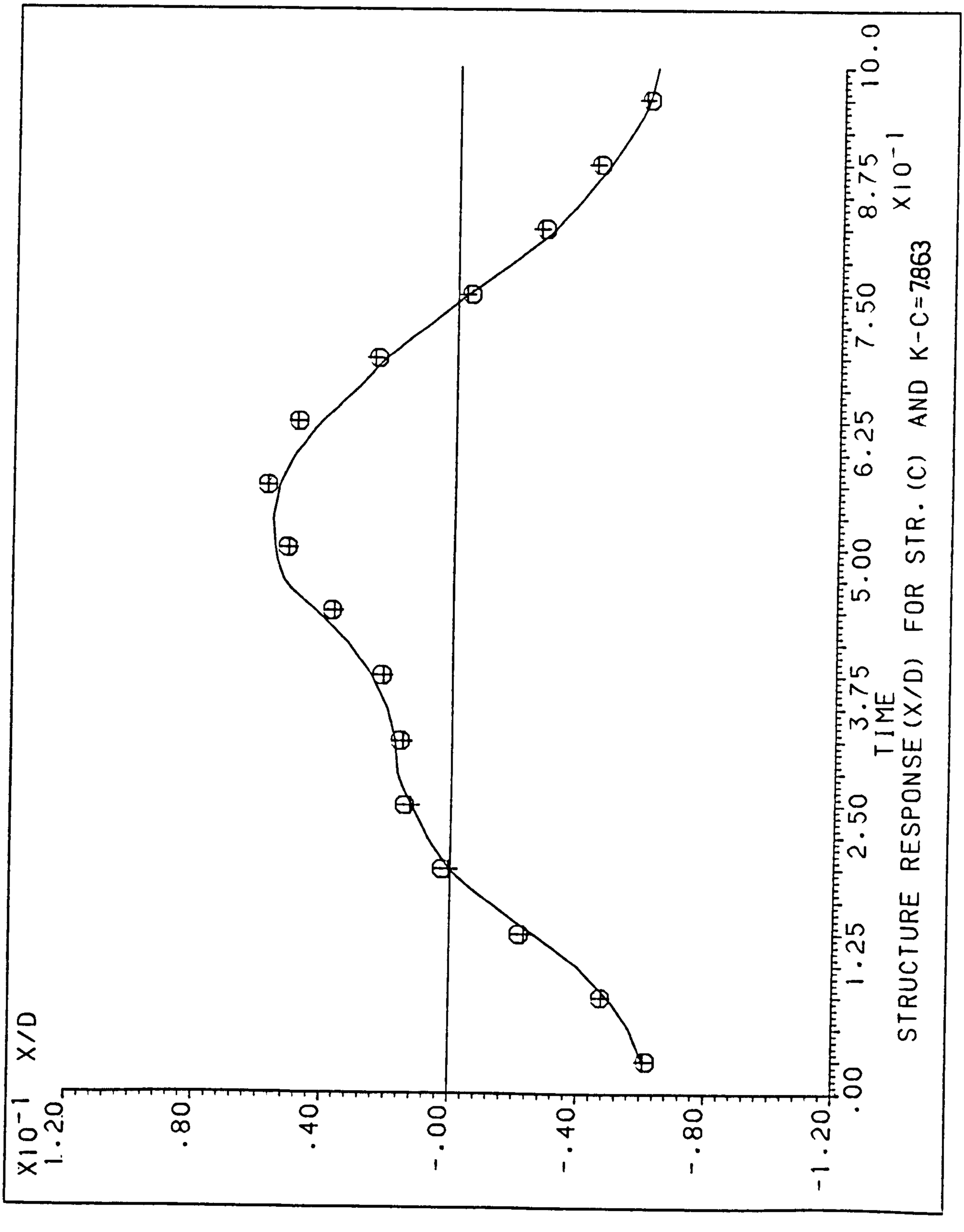


Figure (21D) - Structure's Response (X/D) during the Wave Cycle for Structure C (K-C = 7.863).

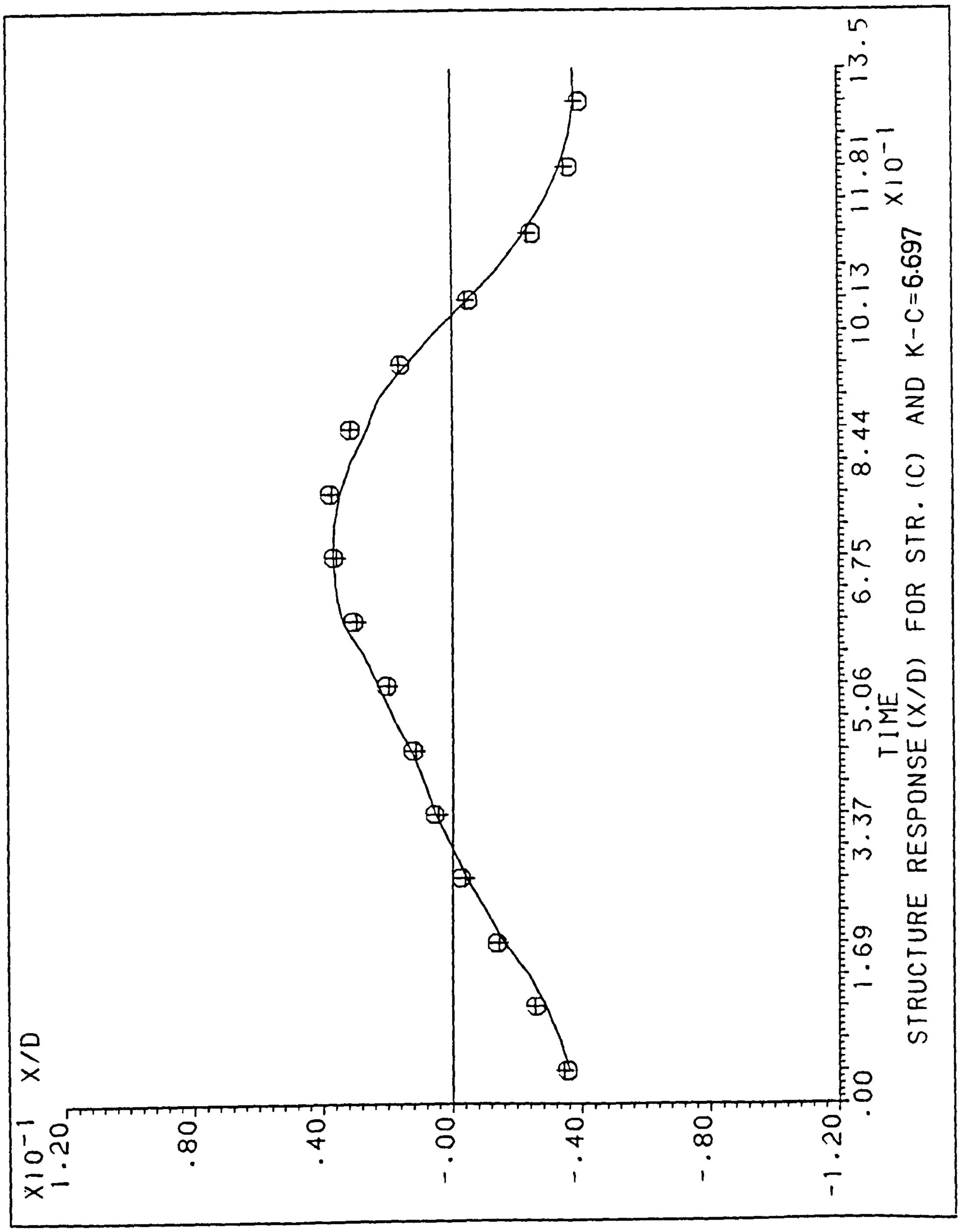


Figure (22D) - Structure's Response (X/D) during the Wave Cycle for Structure C (K-C = 6.697).

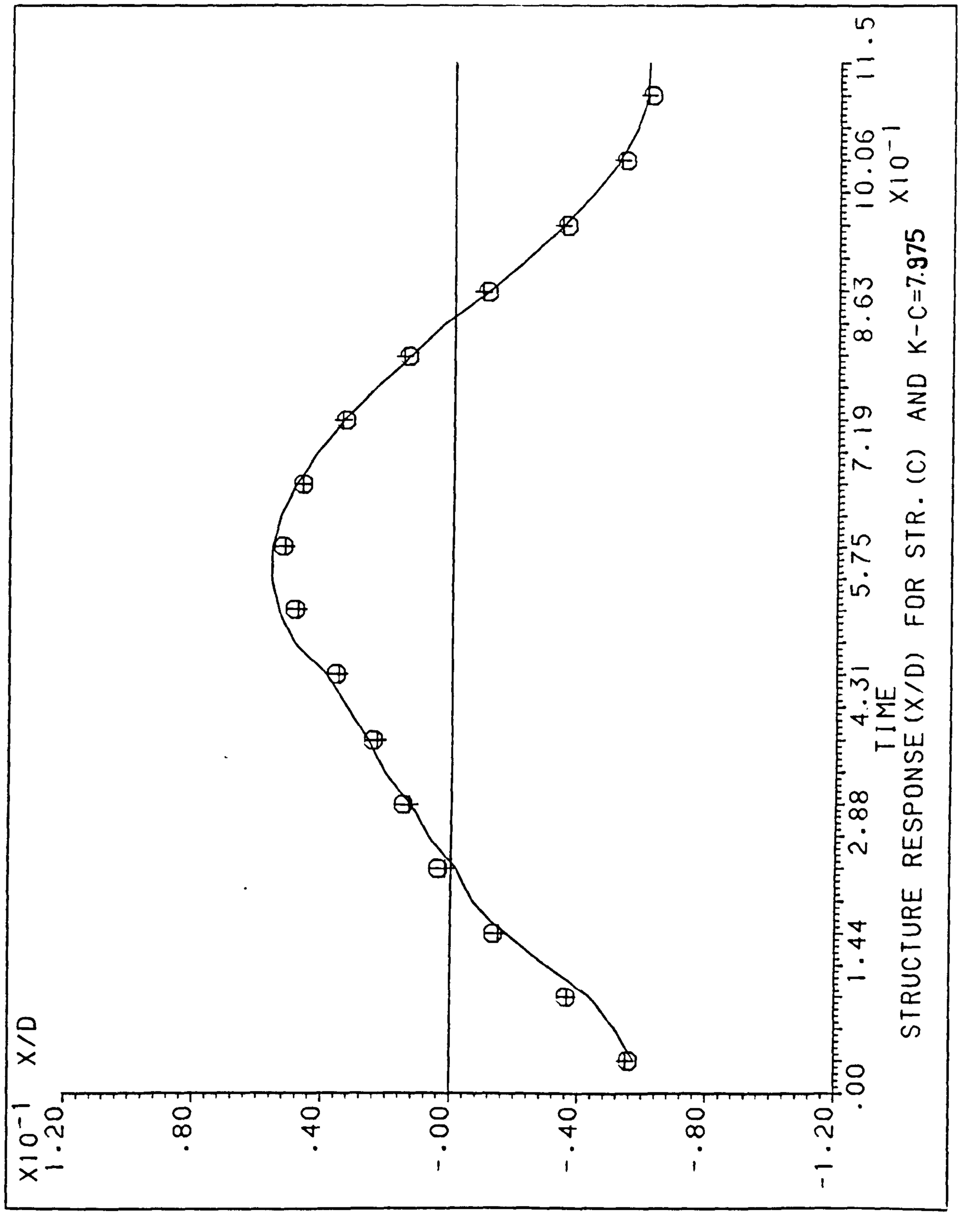


Figure (23D) - Structure's Response (X/D) during the Wave Cycle for Structure C (K-C = 7.975).

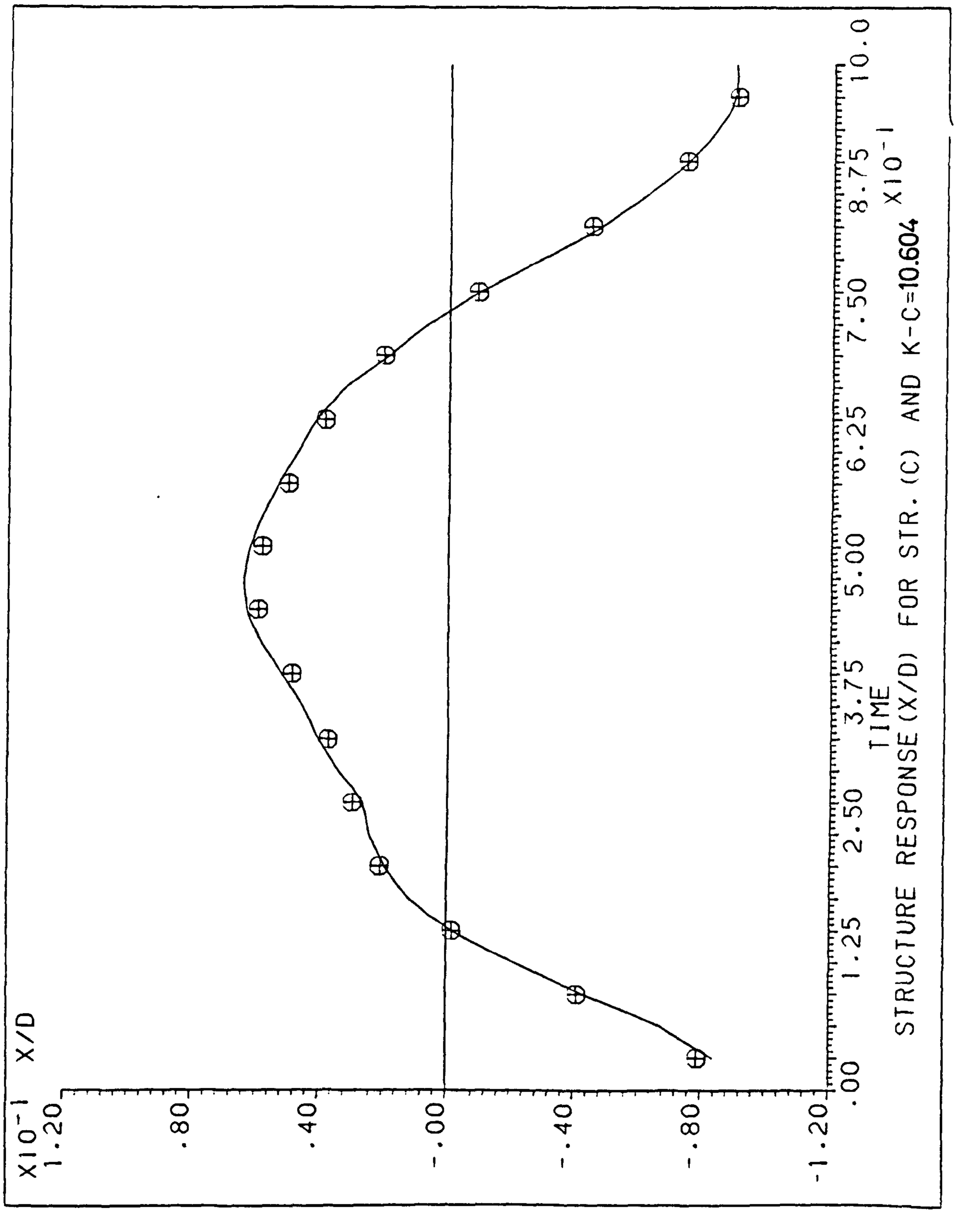


Figure (24D) - Structure's Response (X/D) during the Wave Cycle for Structure C (K-C = 10.604).

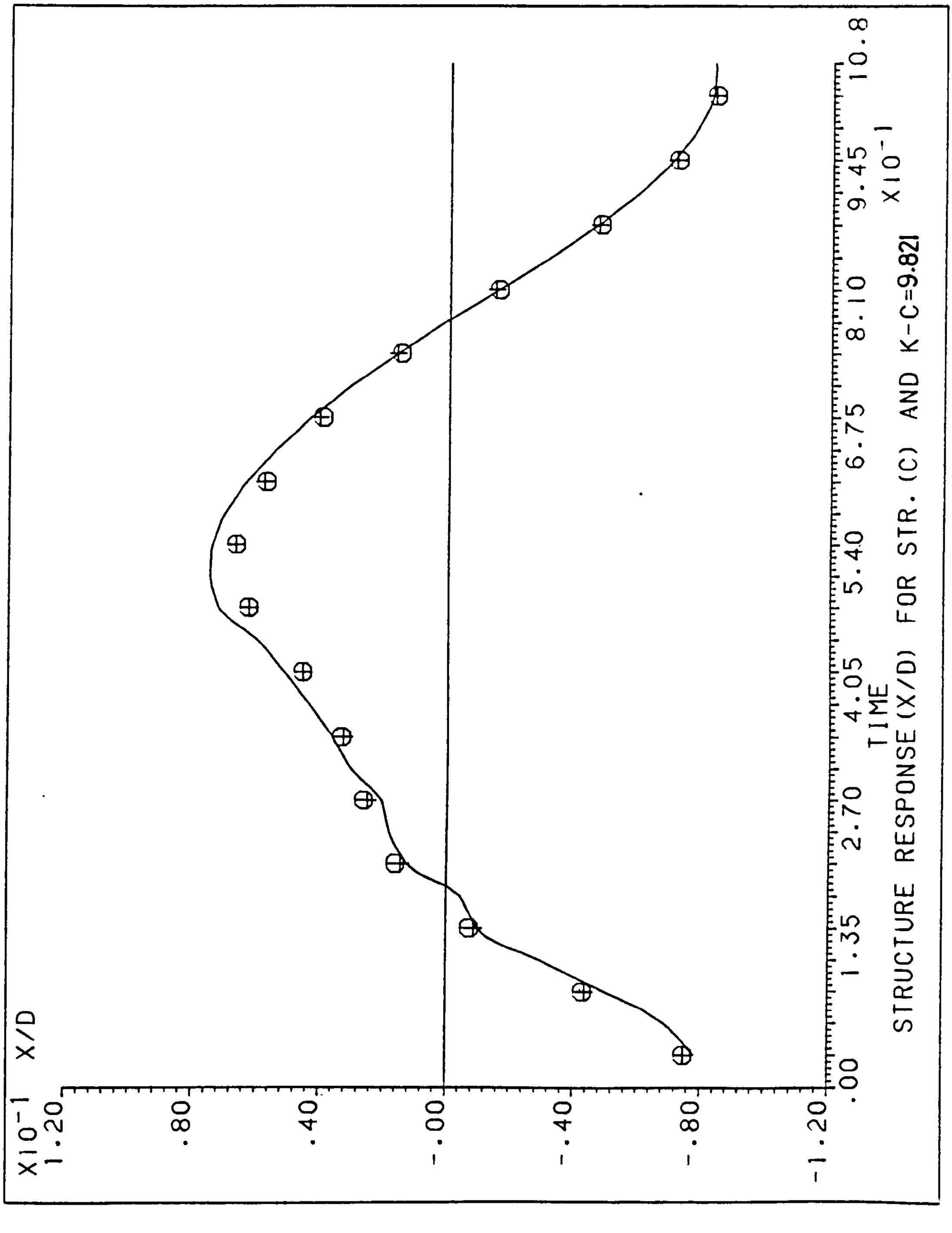


Figure (25D) - Structure's Response (X/D) during the Wave Cycle for Strucrure C (K-C = 9.821).

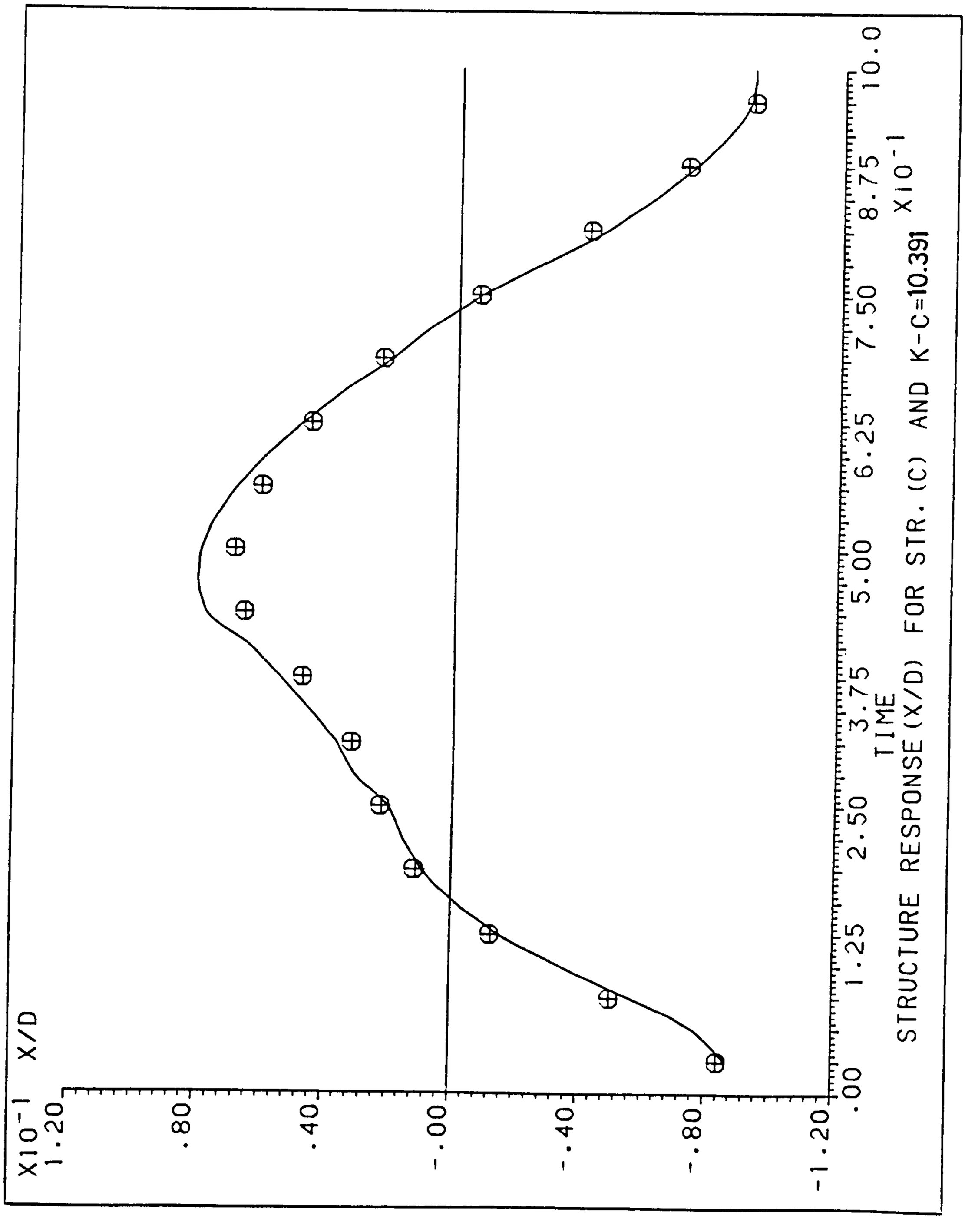


Figure (26D) - Structure's Response (X/D) during the Wave Cycle for Structure C (K-C = 10.391).

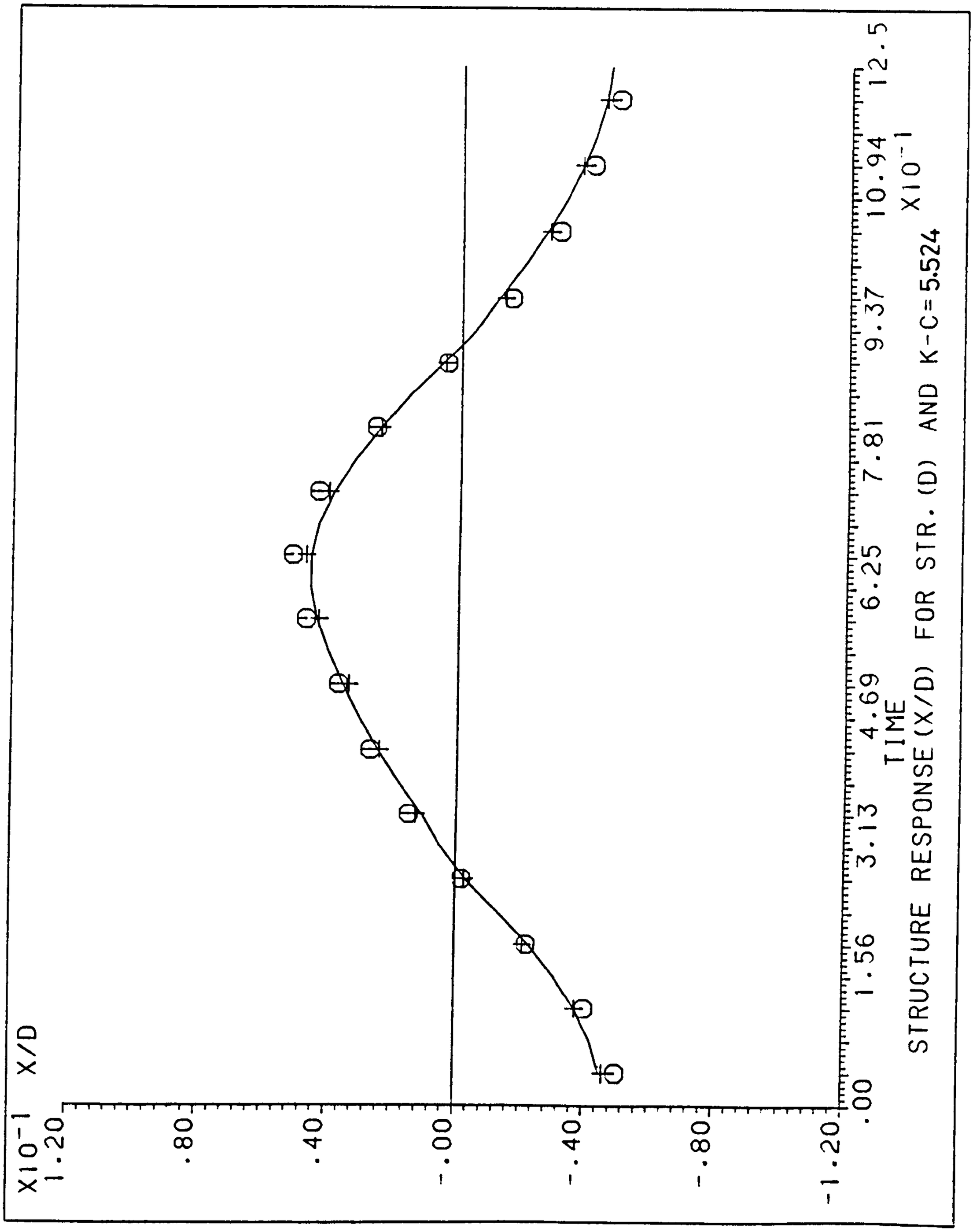


Figure (27D) - Structure's Response (X/D) during the Wave Cycle for Structure D (K-C = 5.524).

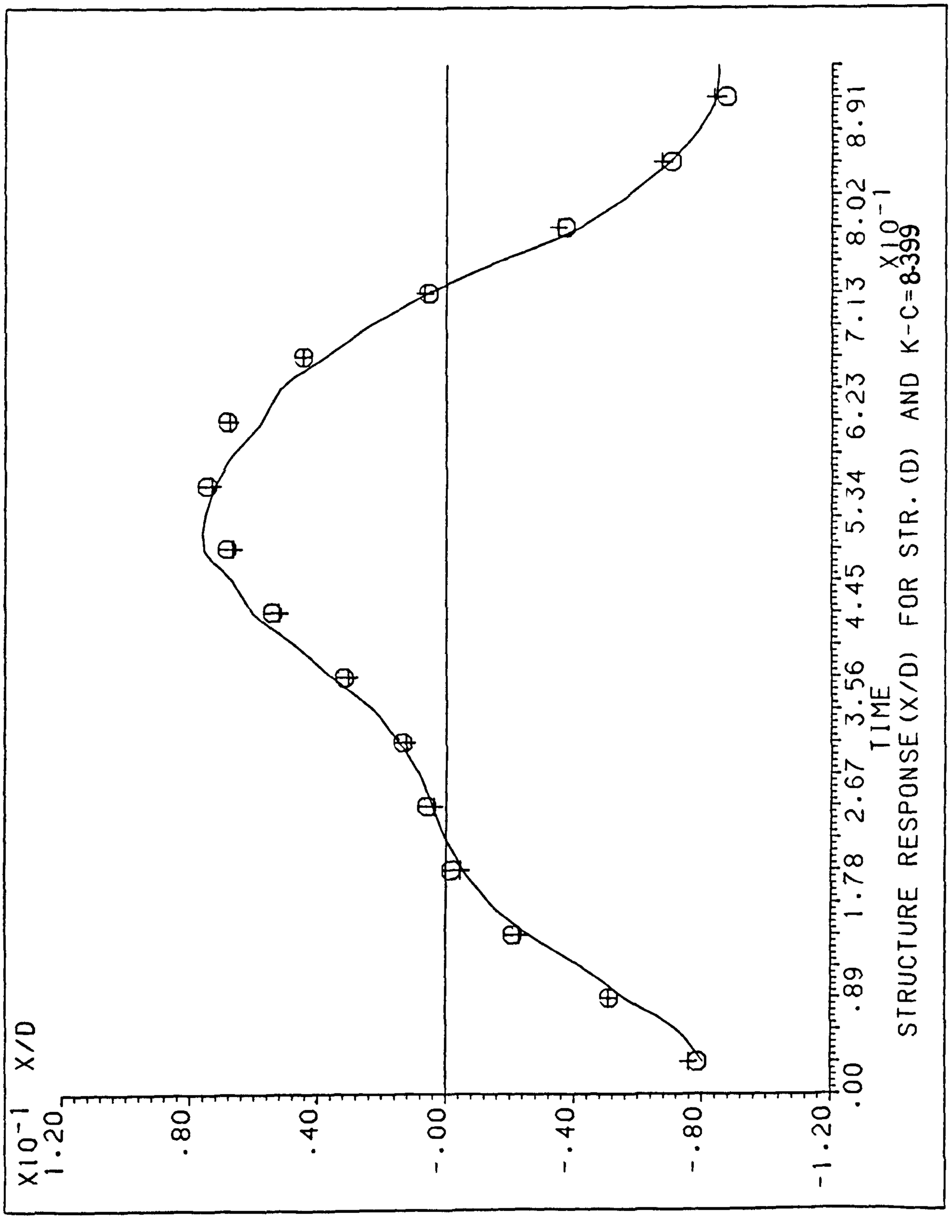


Figure (28D) - Structure's Response (X/D) during the Wave Cycle for Structure D (K-C = 8.399).

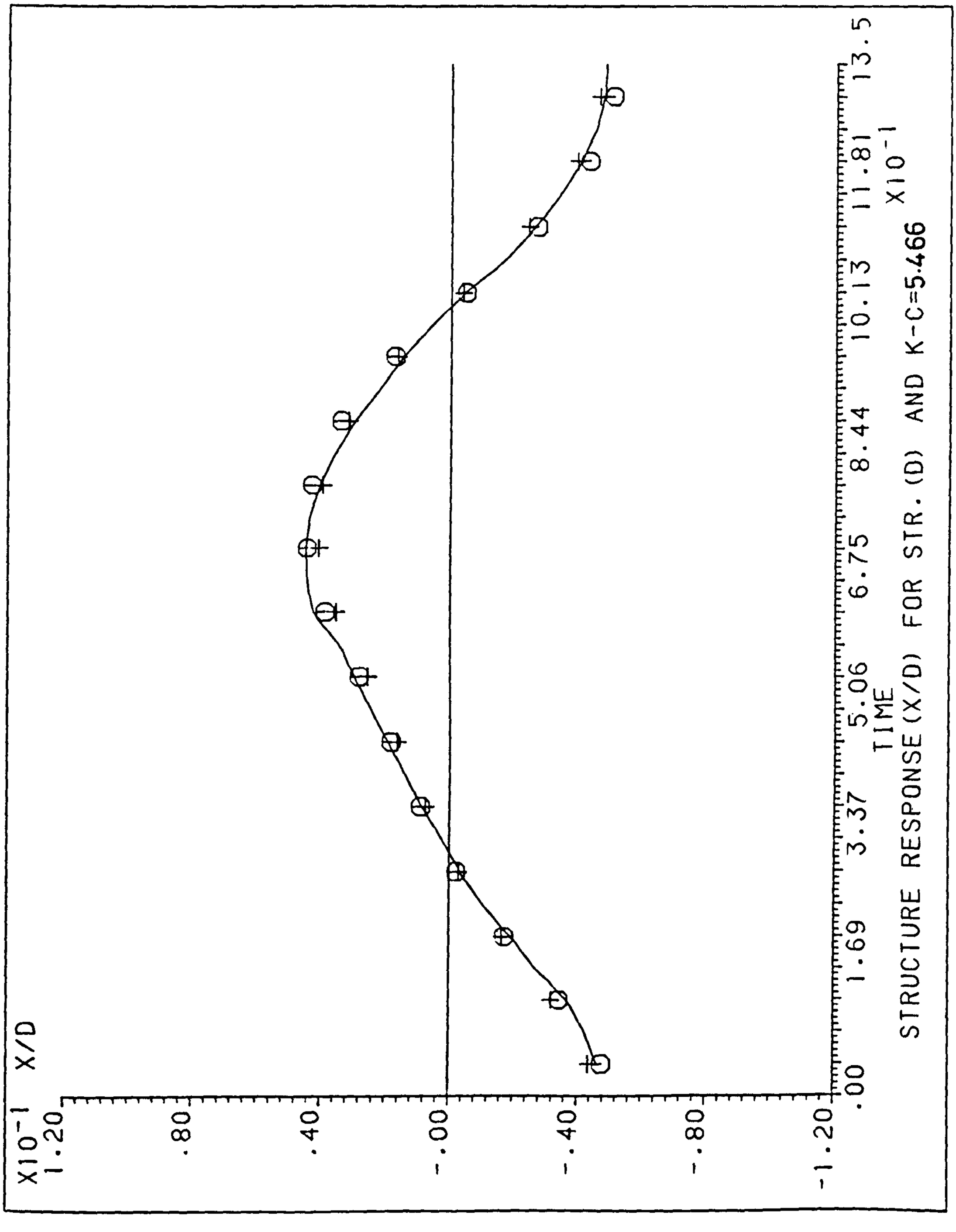


Figure (29D) - Structure's Response (X/D) during the Wave Cycle for Structure D (K-C = 5.466).

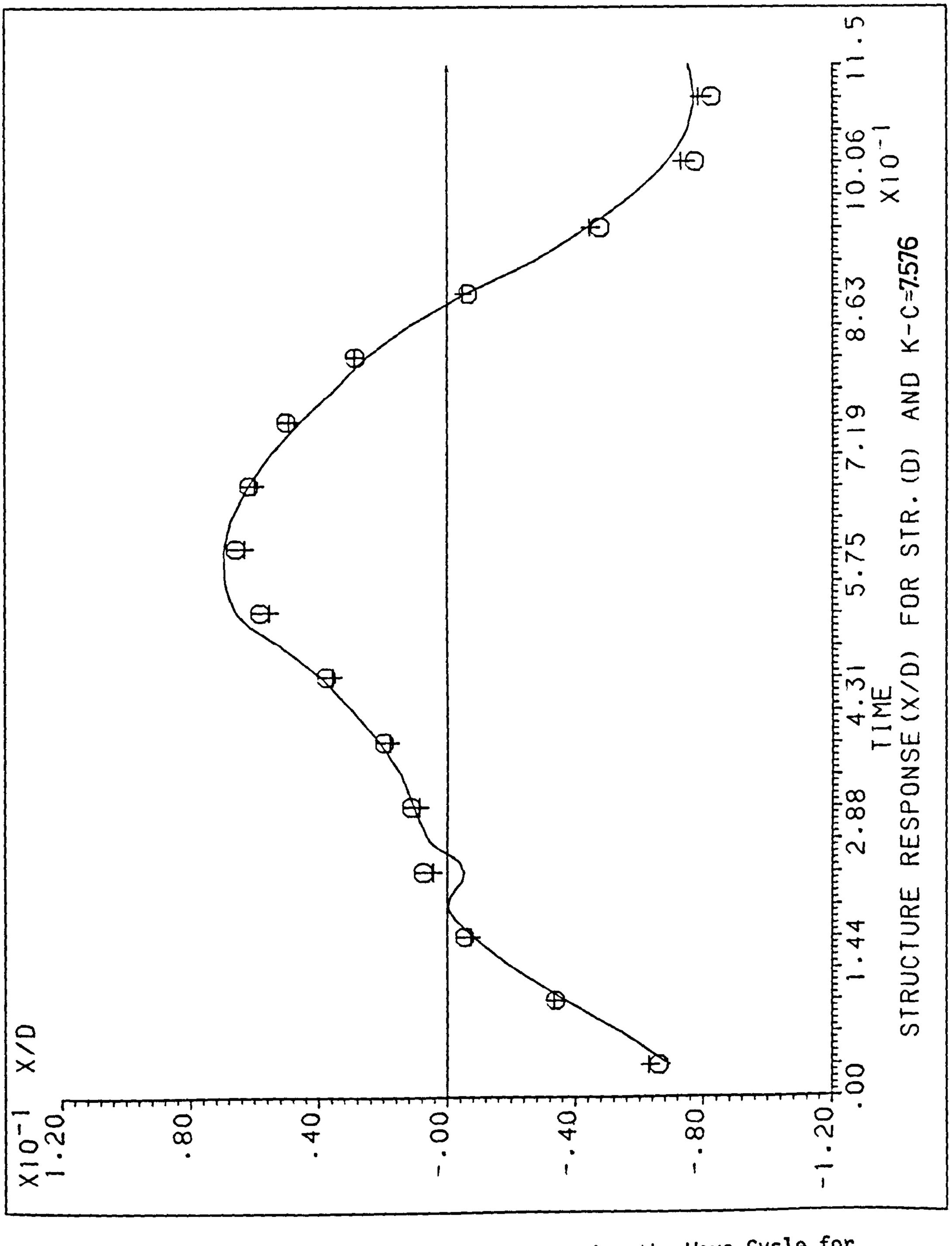


Figure (30D) - Structure's Response (X/D) during the Wave Cycle for Structure D (K-C = 7.576).

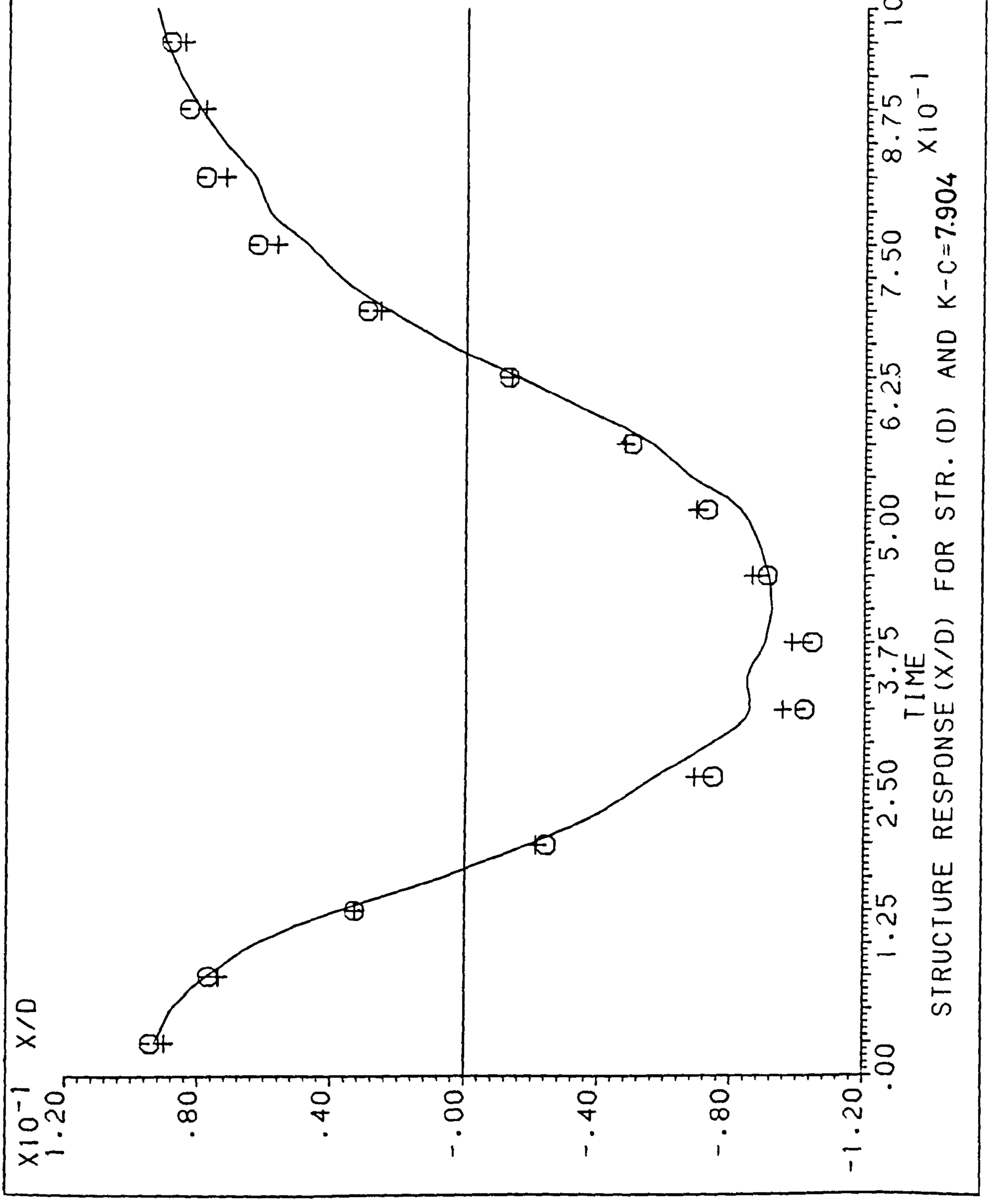


Figure (31D) - Structure's Response (X/D) during the Wave Cycle for Structure D (K-C = 7.904).

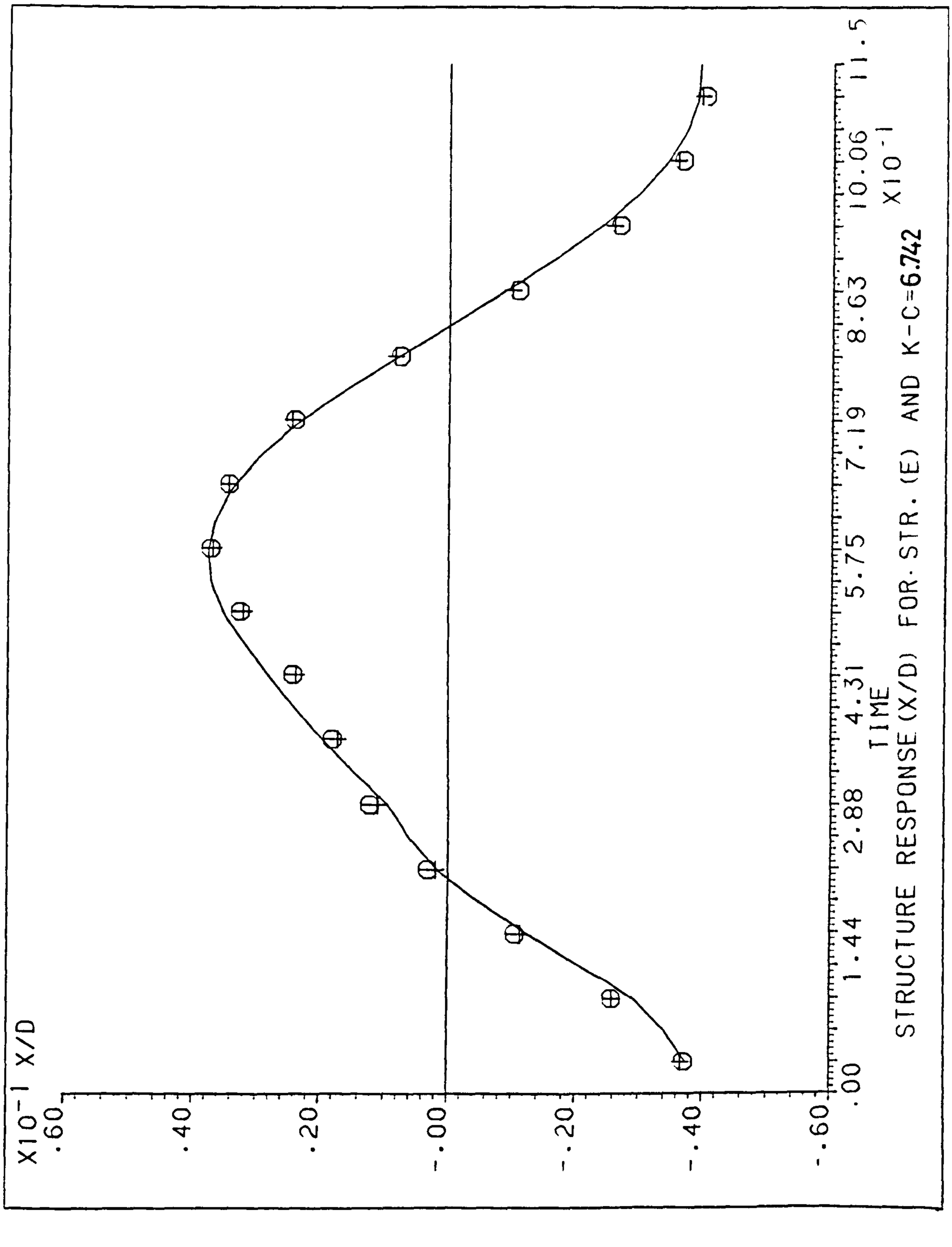


Figure (32D) - Structure's Response (X/D) during the Wave Cycle for Structure E with no Mass at Top (K-C = 6.742).

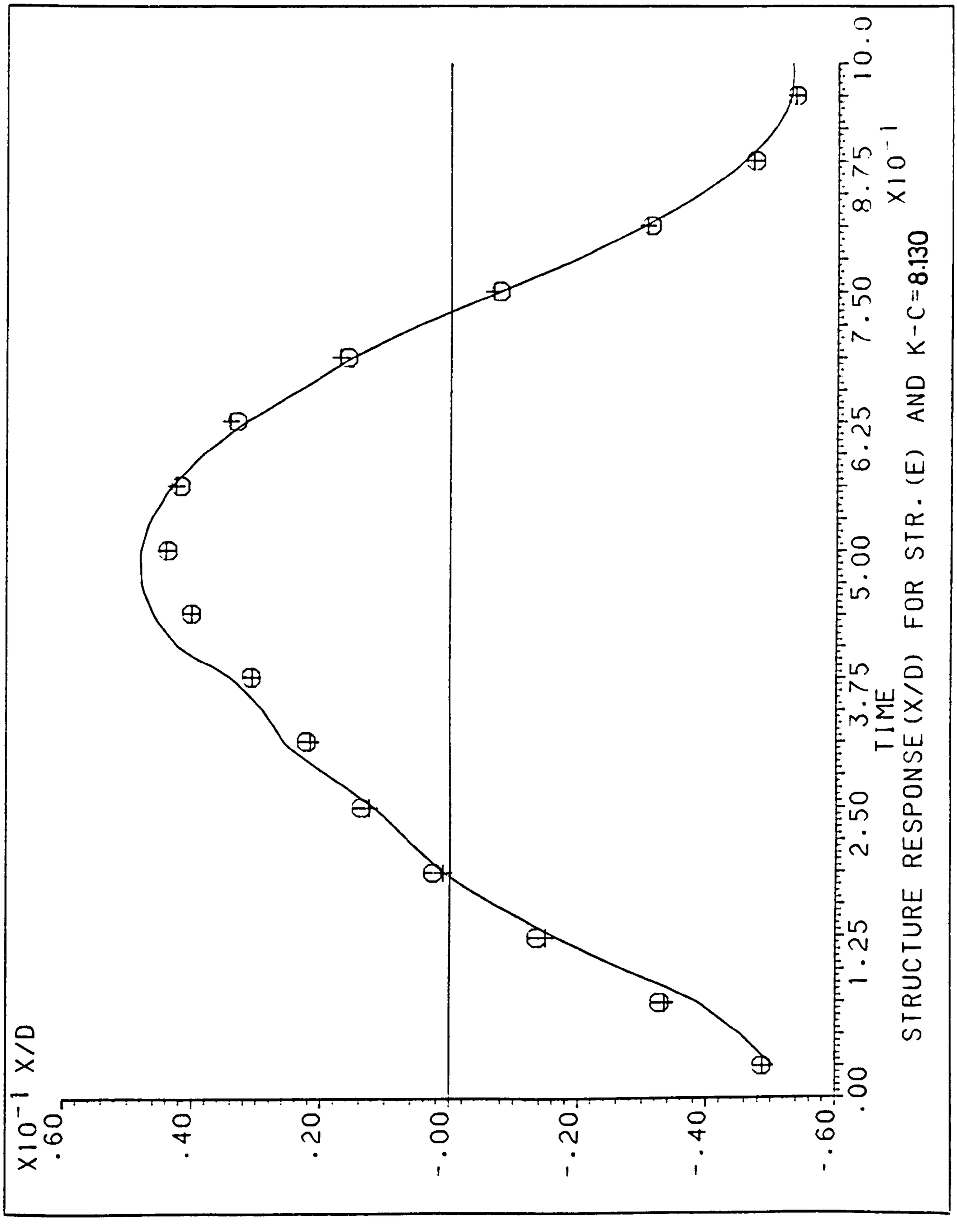


Figure (33D) - Structure's Response (X/D) during the Wave Cycle for Structure E with no Mass at Top (K-C = 8.13).

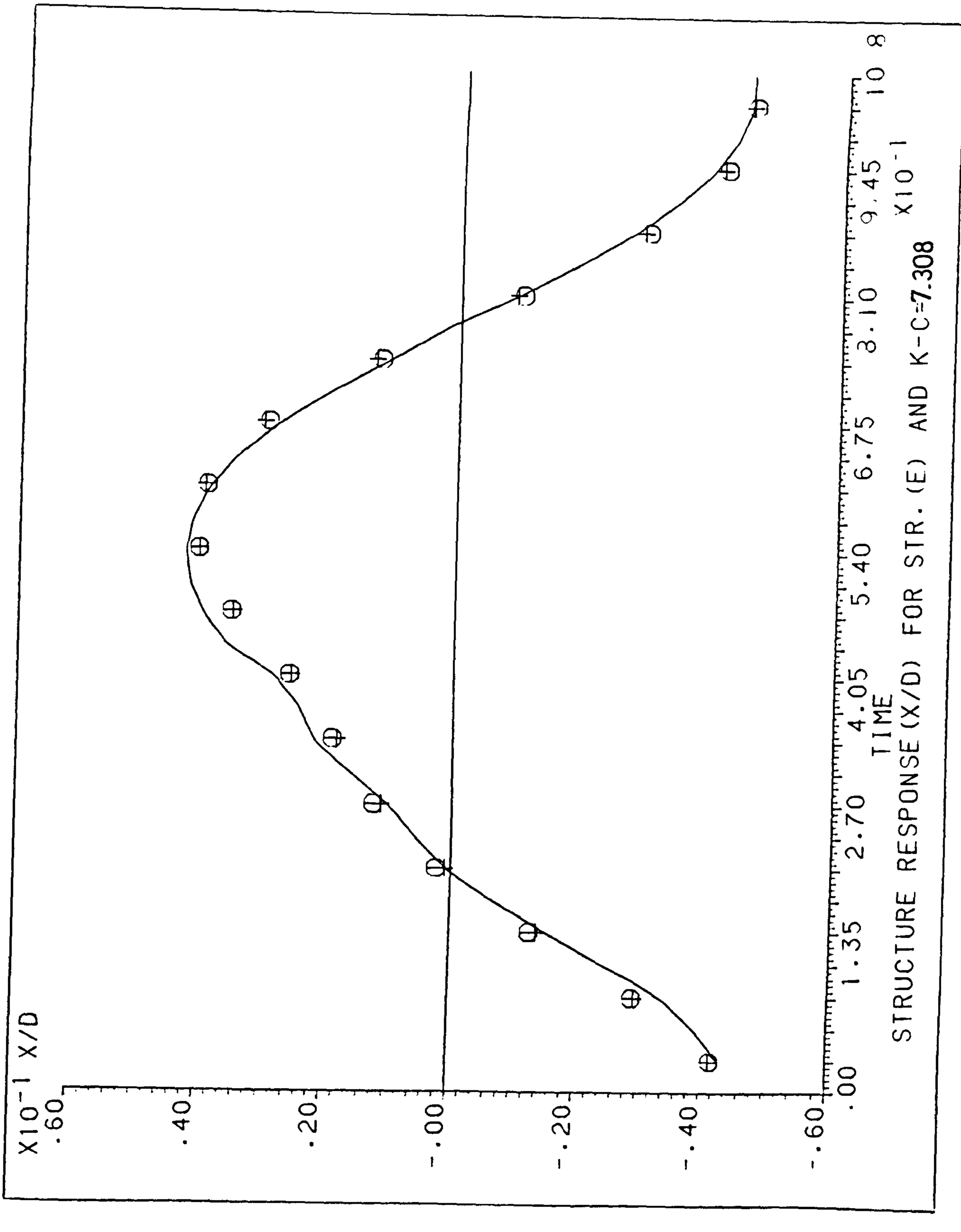


Figure (34D) - Structure's Response(X/D) during the Wave Cycle for Structure E with no Mass at Top (K-C = 7.308).

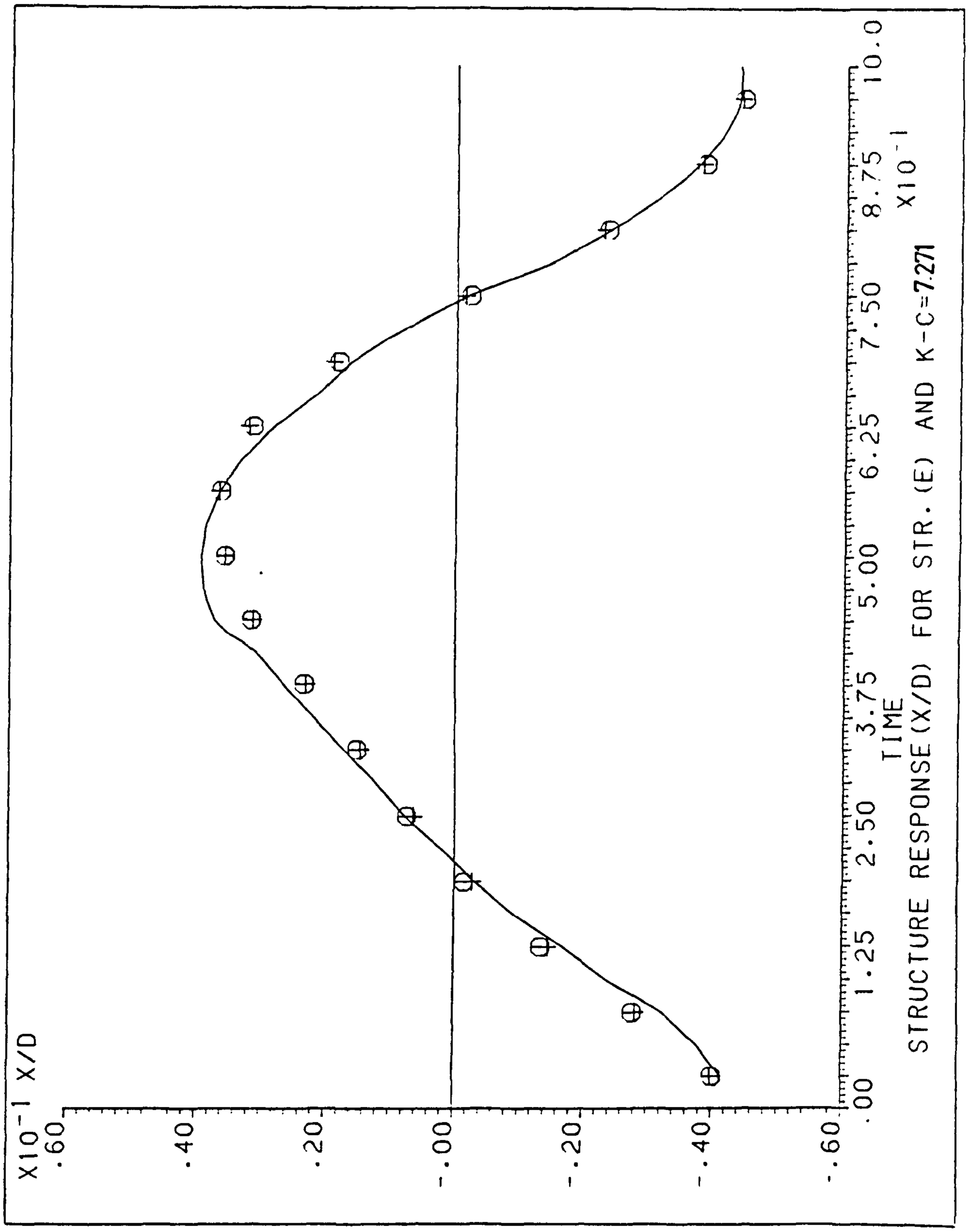


Figure (35D) - Structure's Response (X/D) during the Wave Cycle for Structure E with No Mass at Top (K-C = 7.271).

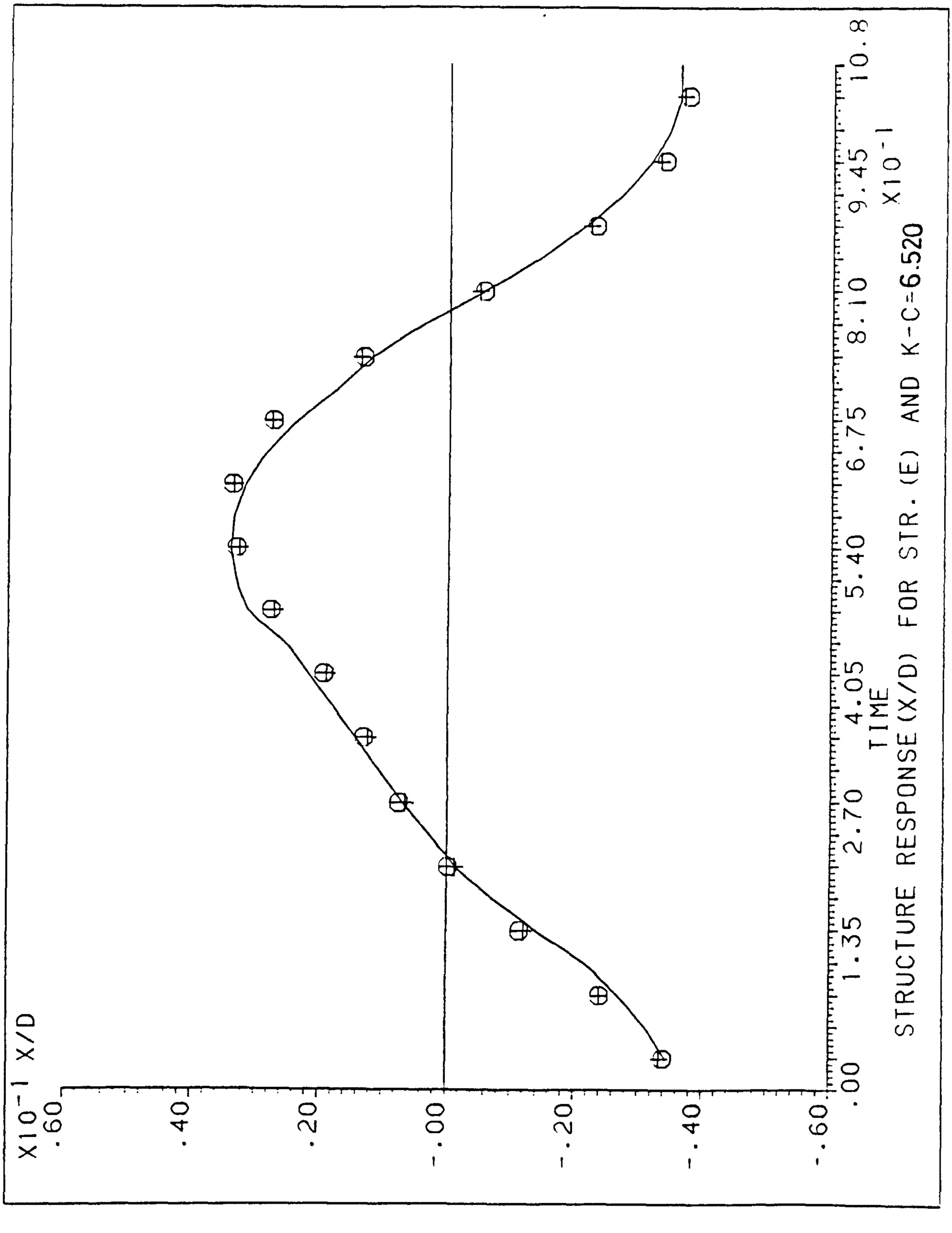


Figure (36D) - Structrure's Response (X/D) during the Wave Cycle for Structure E with no Mass at Top (K-C = 6.520).

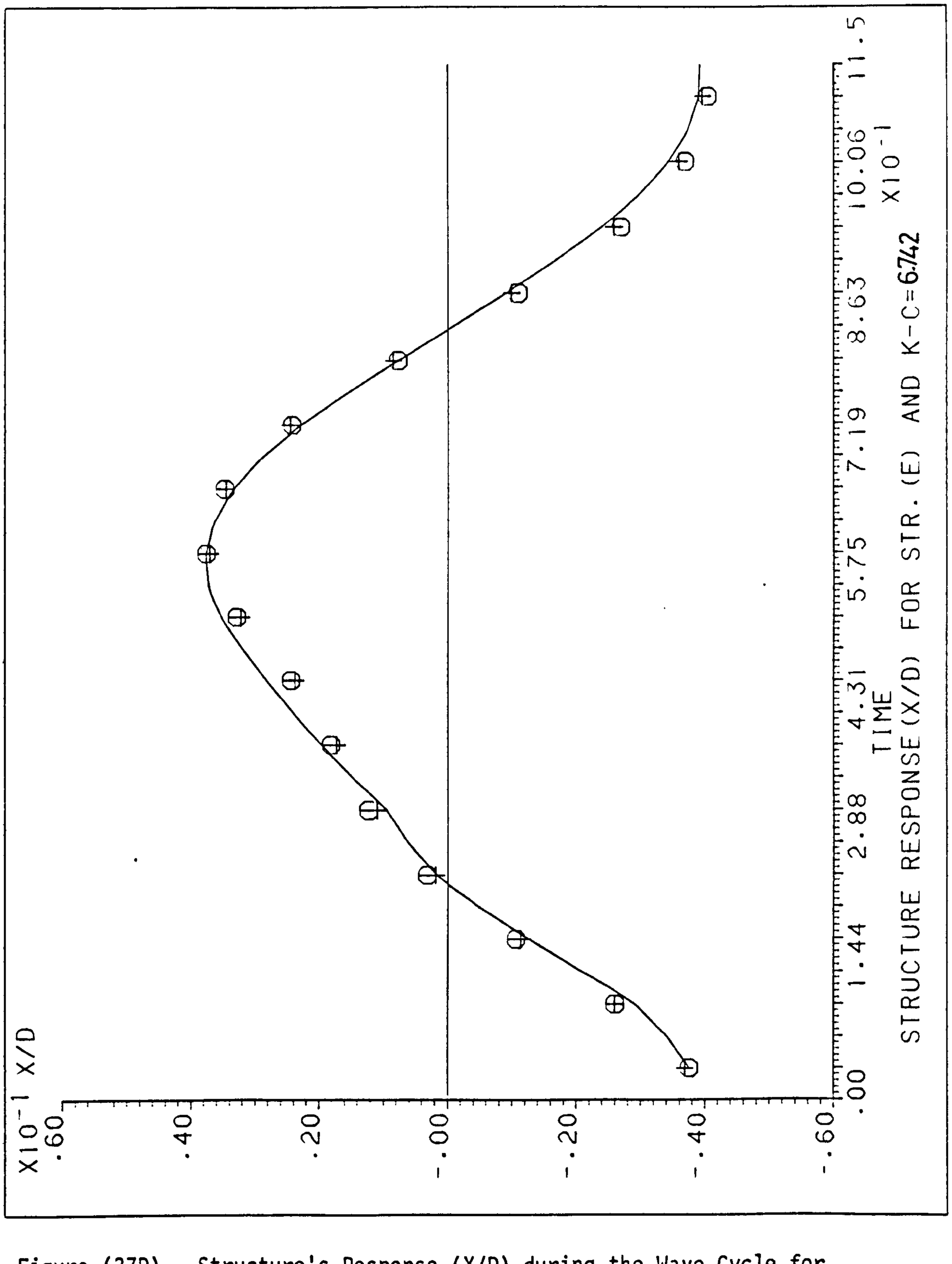


Figure (37D) - Structure's Response (X/D) during the Wave Cycle for Structure E with 1010 Grs at Top (K-C = 6.743).

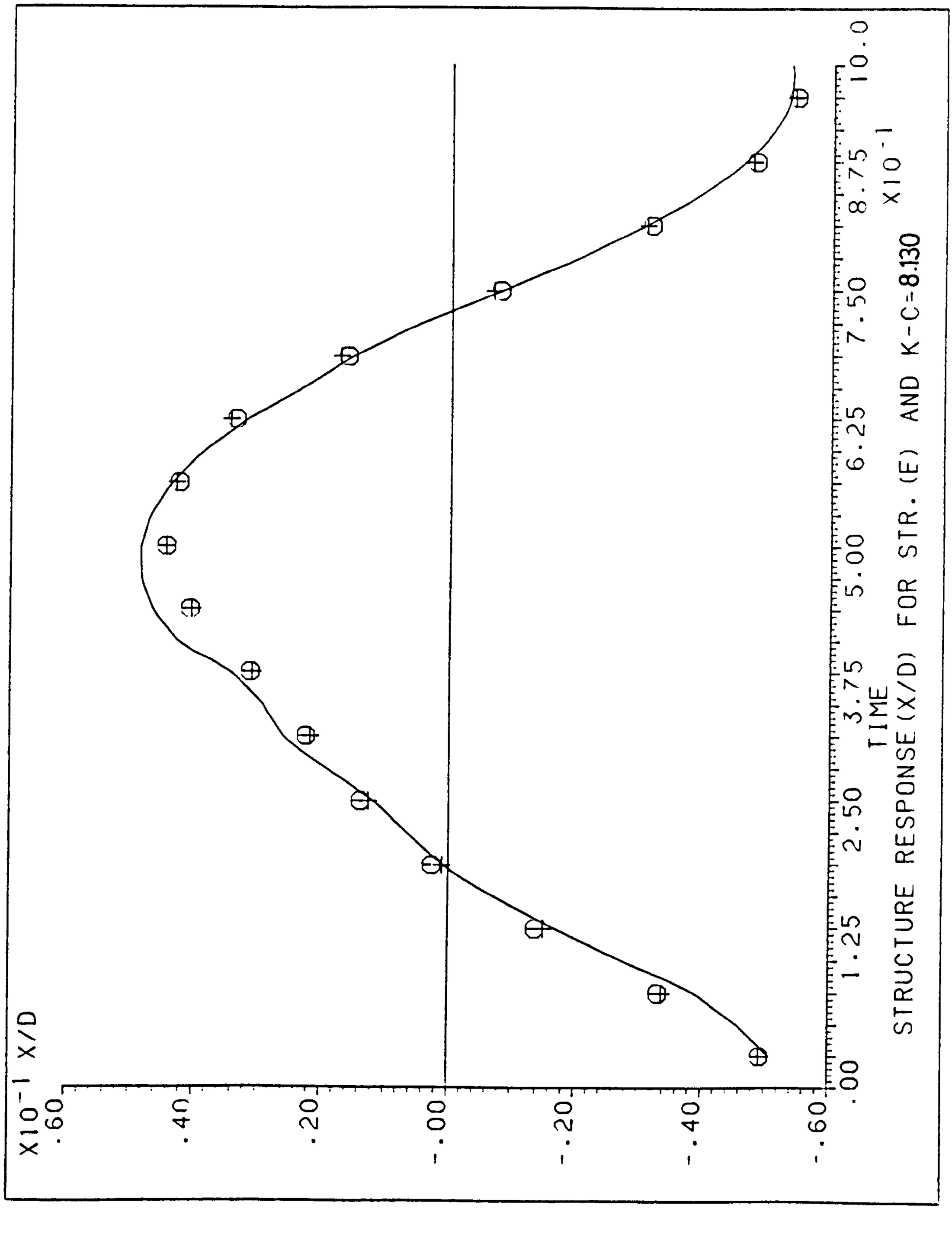


Figure (38D) - Structure's Response (X/D) during the Wave Cycle for Structure E with 1010 Grs at Top (K-C = 8.130).

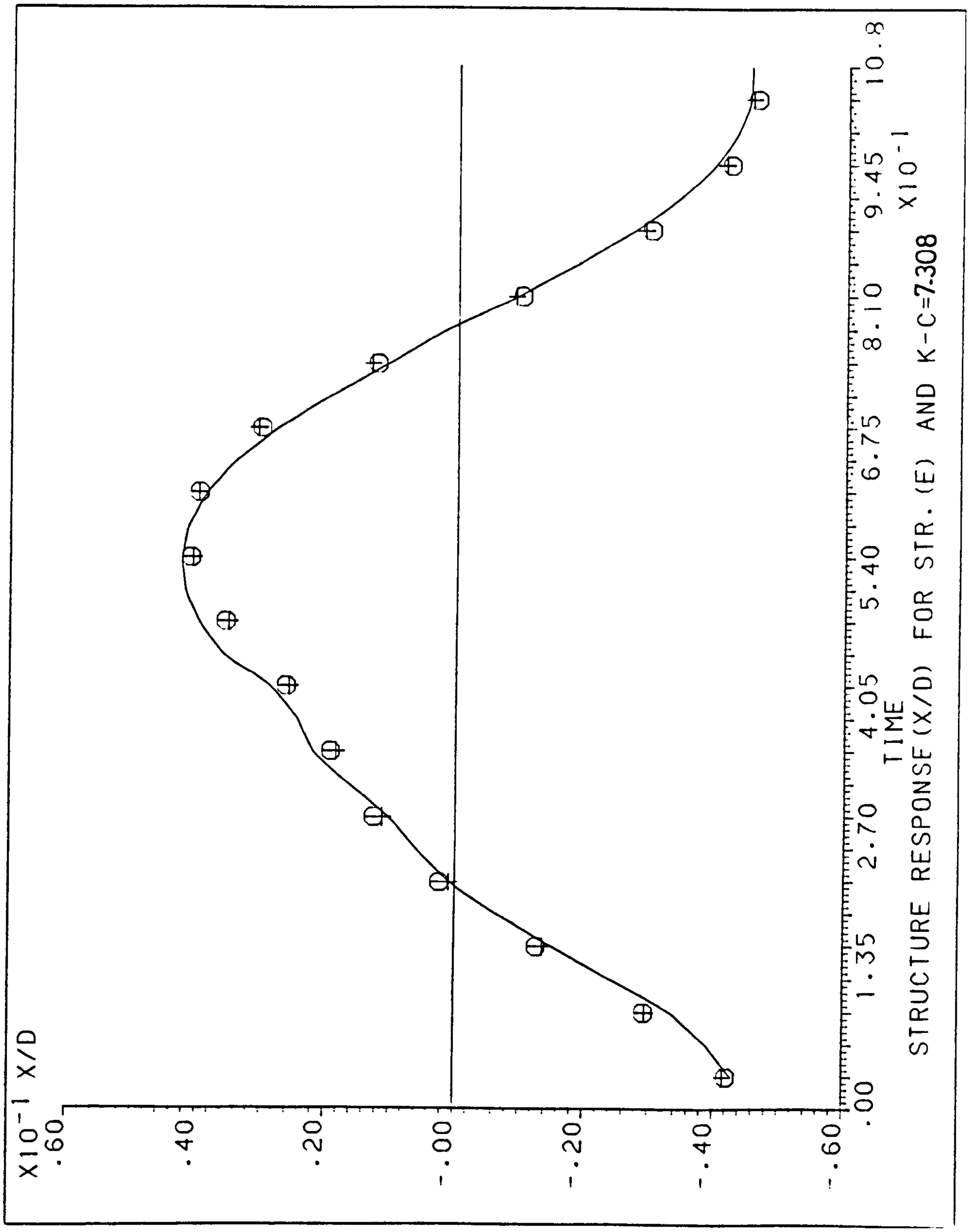


Figure (39D) - Structure's Response (X/D) during the Wave Cycle for Structure E with 1010 Grs at Top (K-C = 7.308).

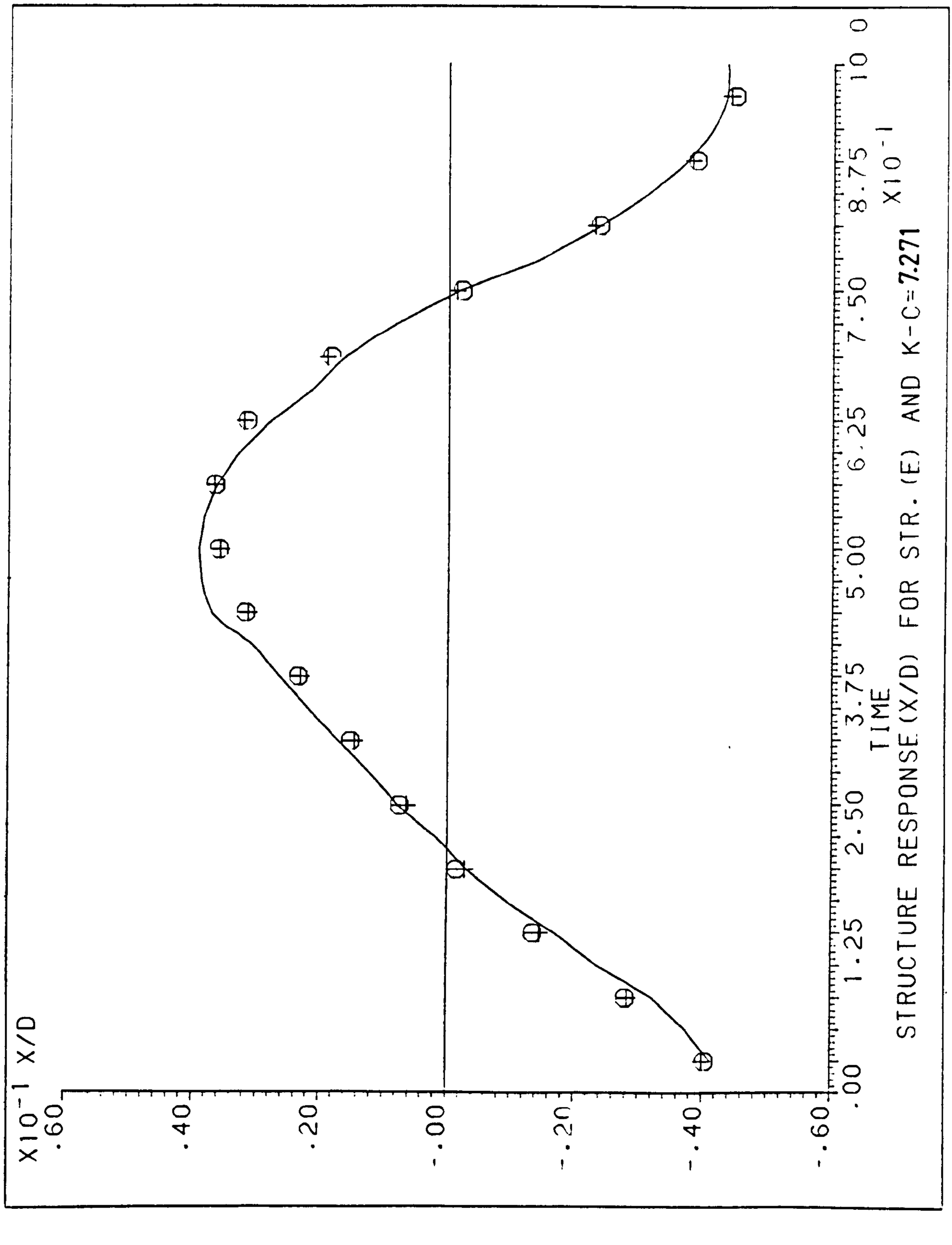


Figure (40D) - Structure's Response (C/D) during the Wave Cycle for Structure E with 1010 Grs at Top (K-C = 7.271).

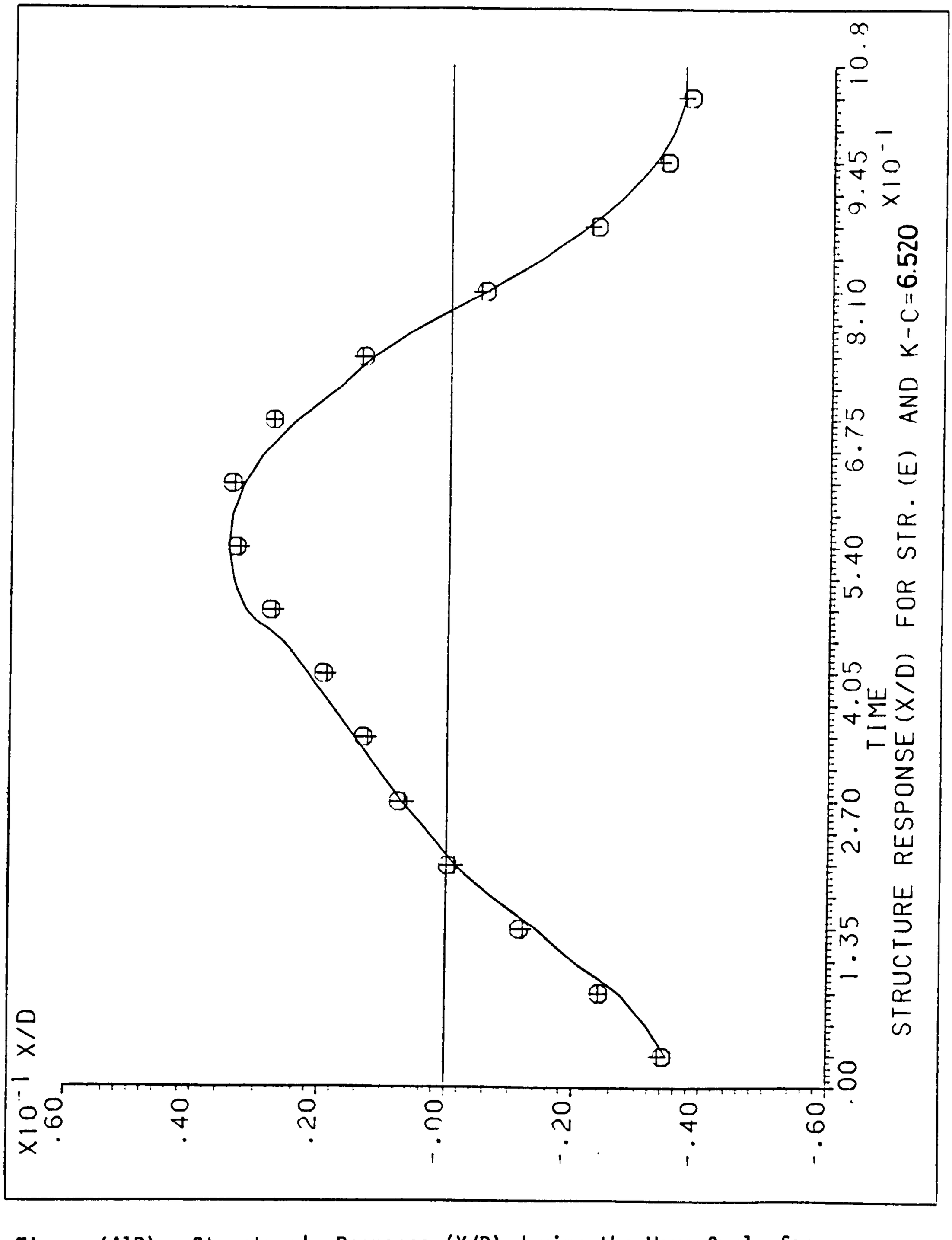


Figure (41D) - Structure's Response (X/D) during the Wave Cycle for Structure E with 1010 Grs at Top (K-C = 6.520).

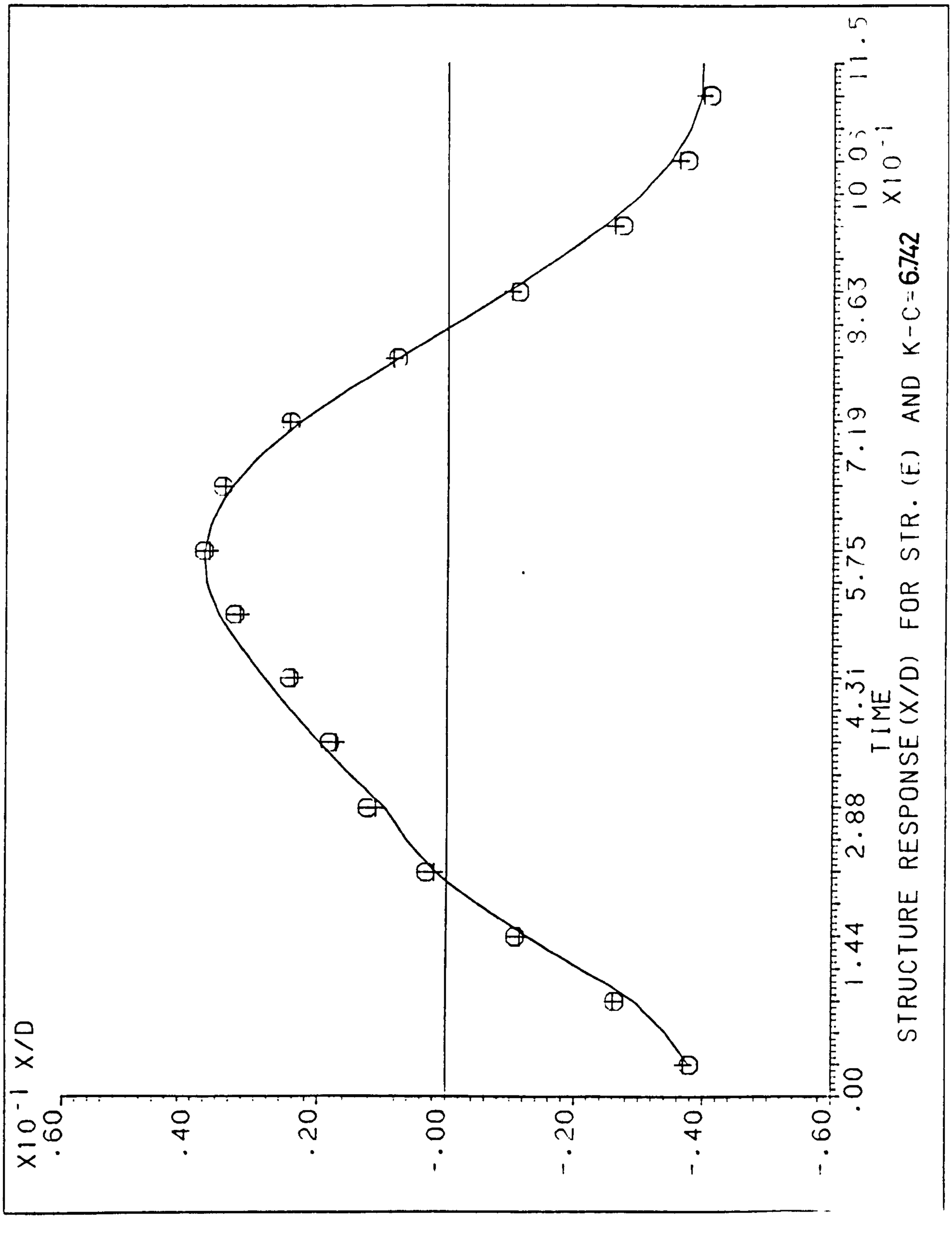


Figure (42D) - Structure's Response (X/D) during the Wave Cycle for Structure E with 2020 Grs at Top (K-C = 6.742).

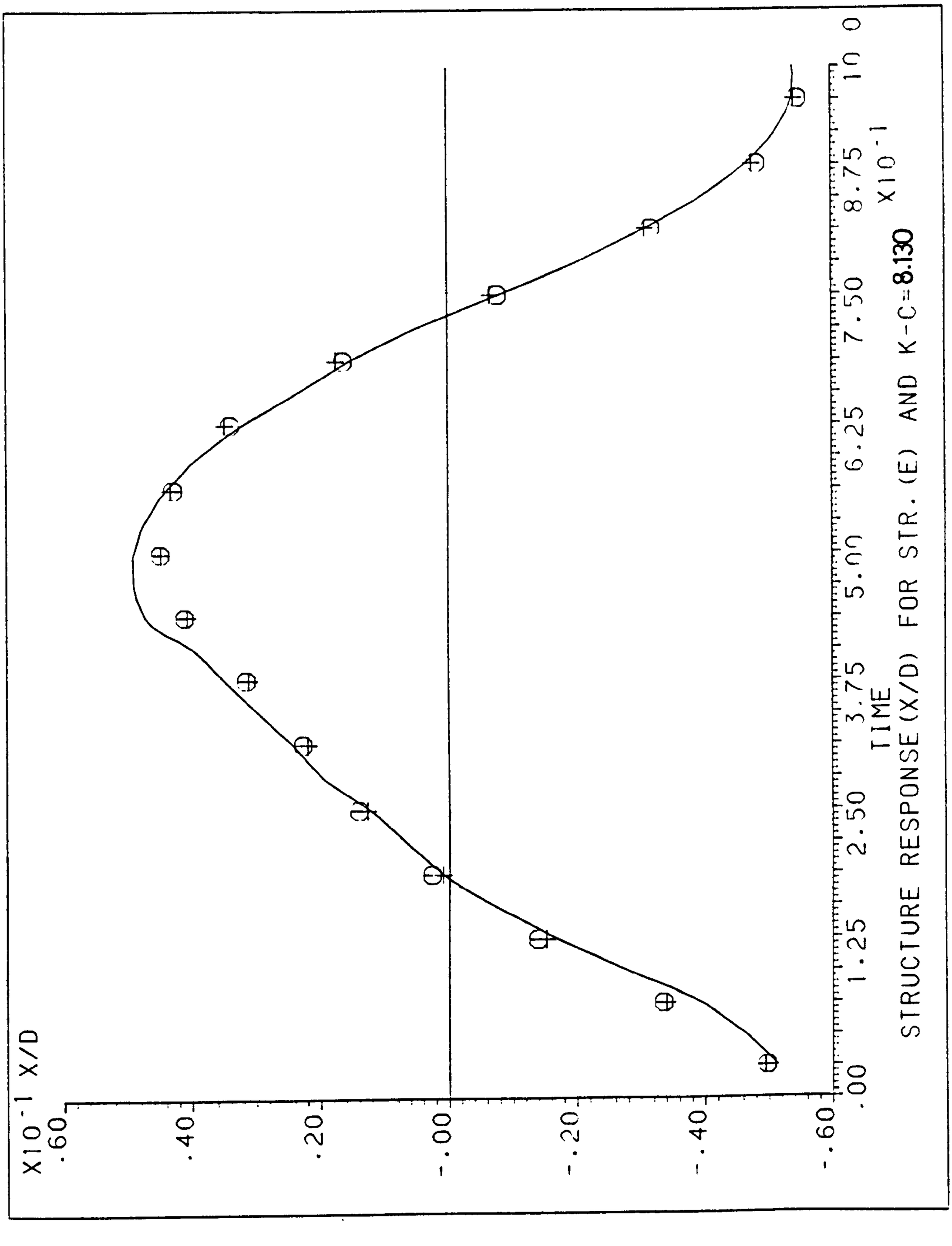


Figure (43D) - Structure's Response (X/D) during the Wave Cycle for Structure E with 2020 Grs at Top (K-C 8.130).

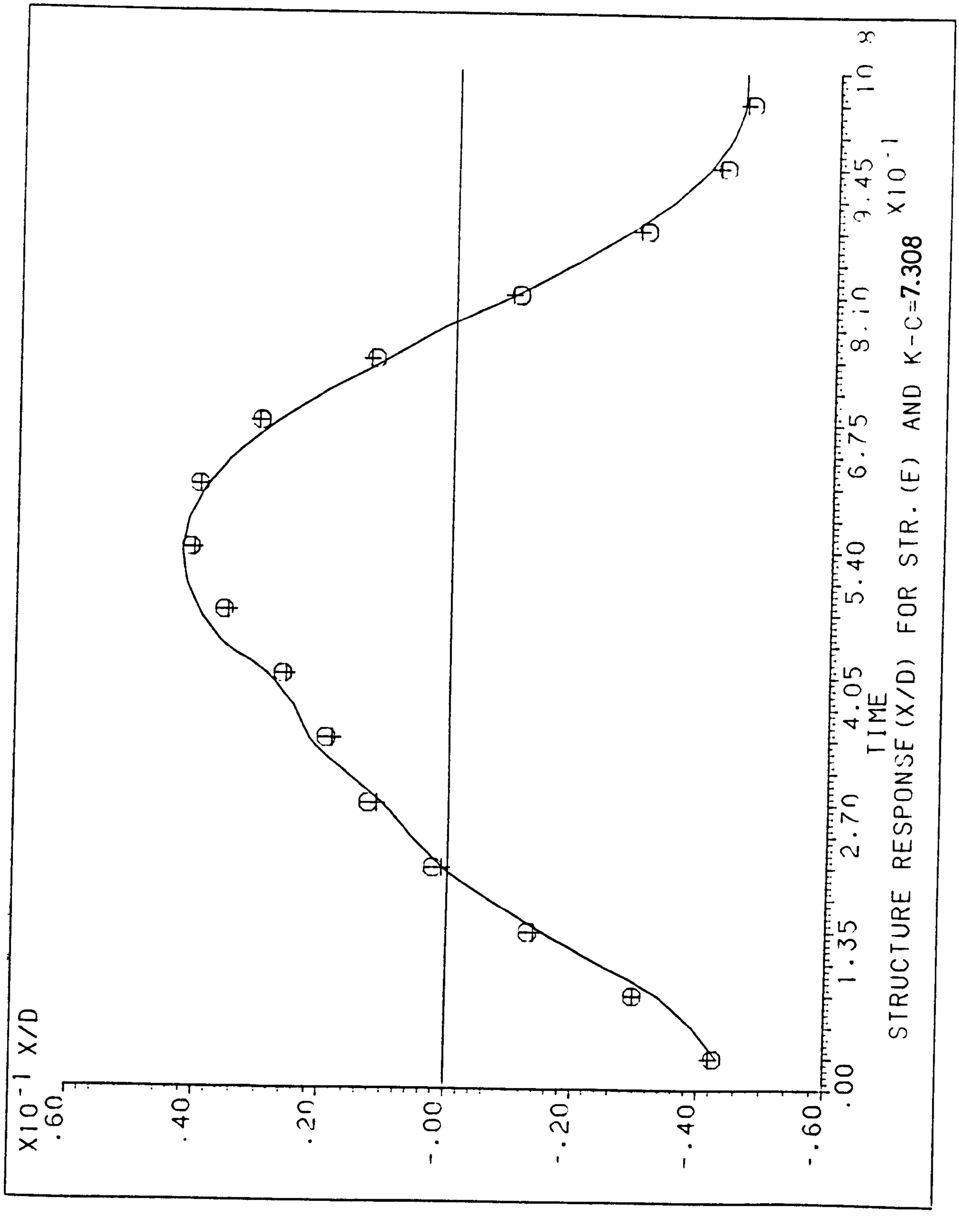


Figure (44D) - Structure's Response (X/D) during the Wave Cycle for Structure E with 2020 Grs at Top (K-C = 7.308).

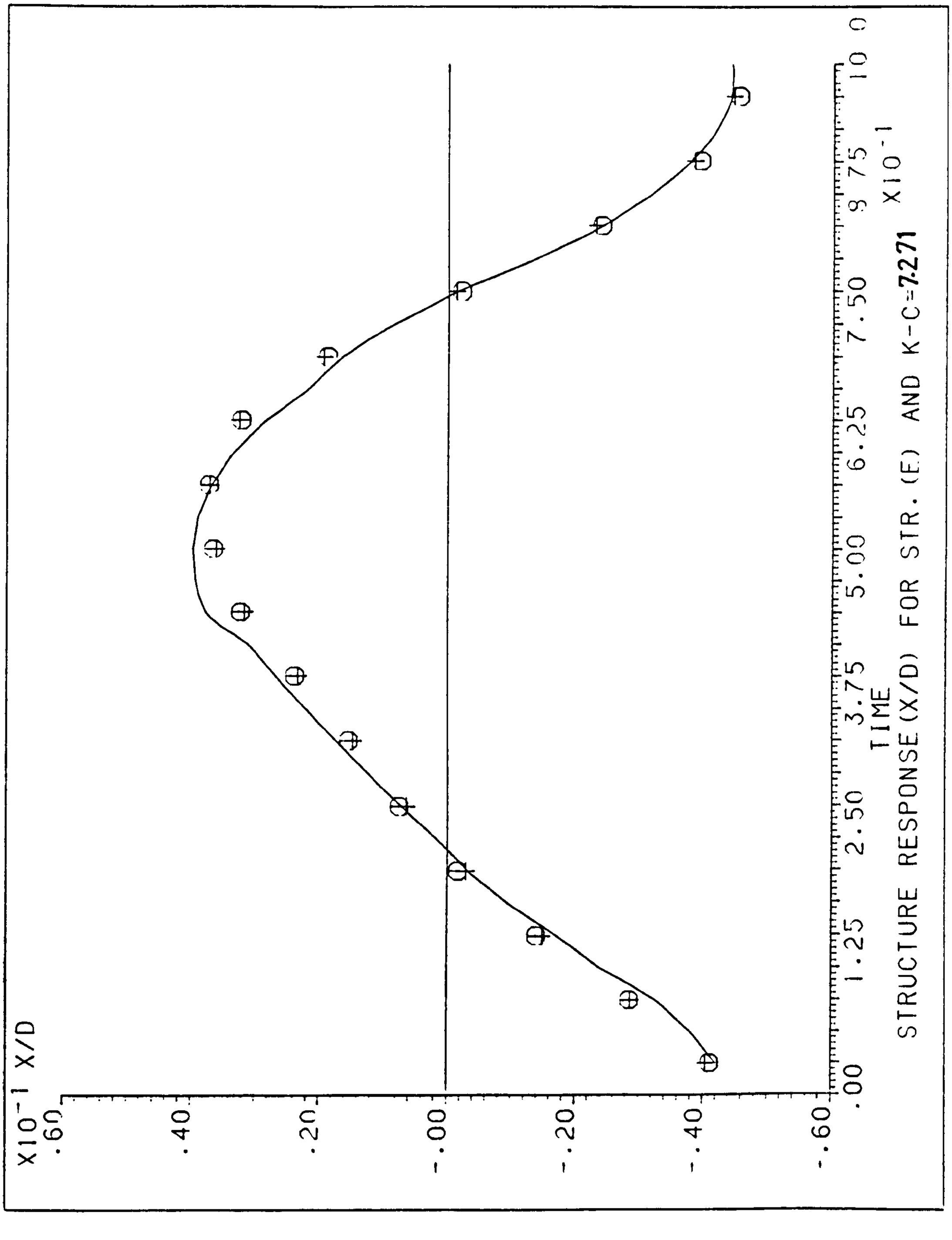


Figure (45D) - Structure's Response (X/D) during the Wave Cycle for Structure E with 2020 Grs at Top (K-C = 7.271).

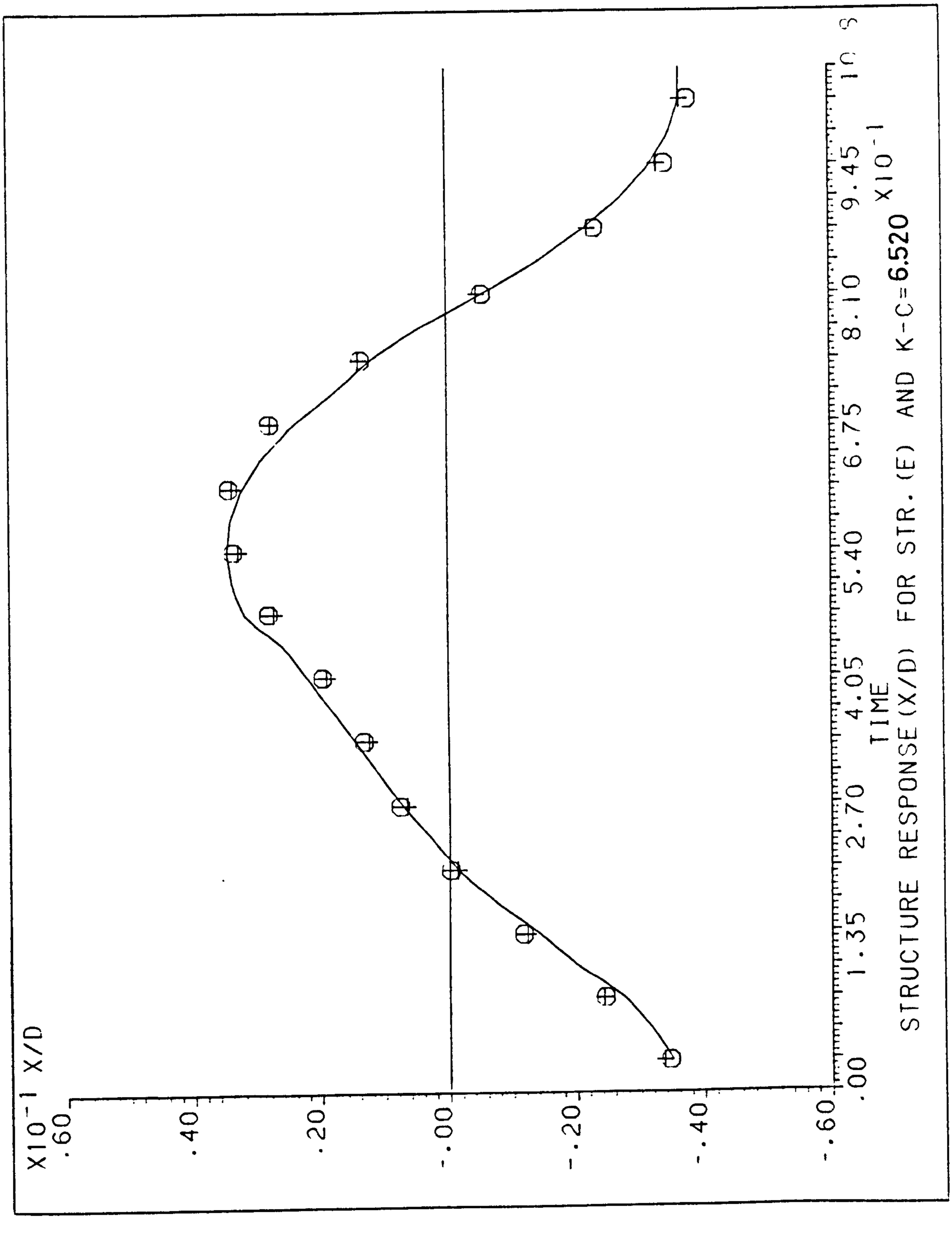


Figure (46D) - Structure's Response (X/D) during the Wave Cycle for Structure E with 2020 Grs at Top (K-C = 6.520).

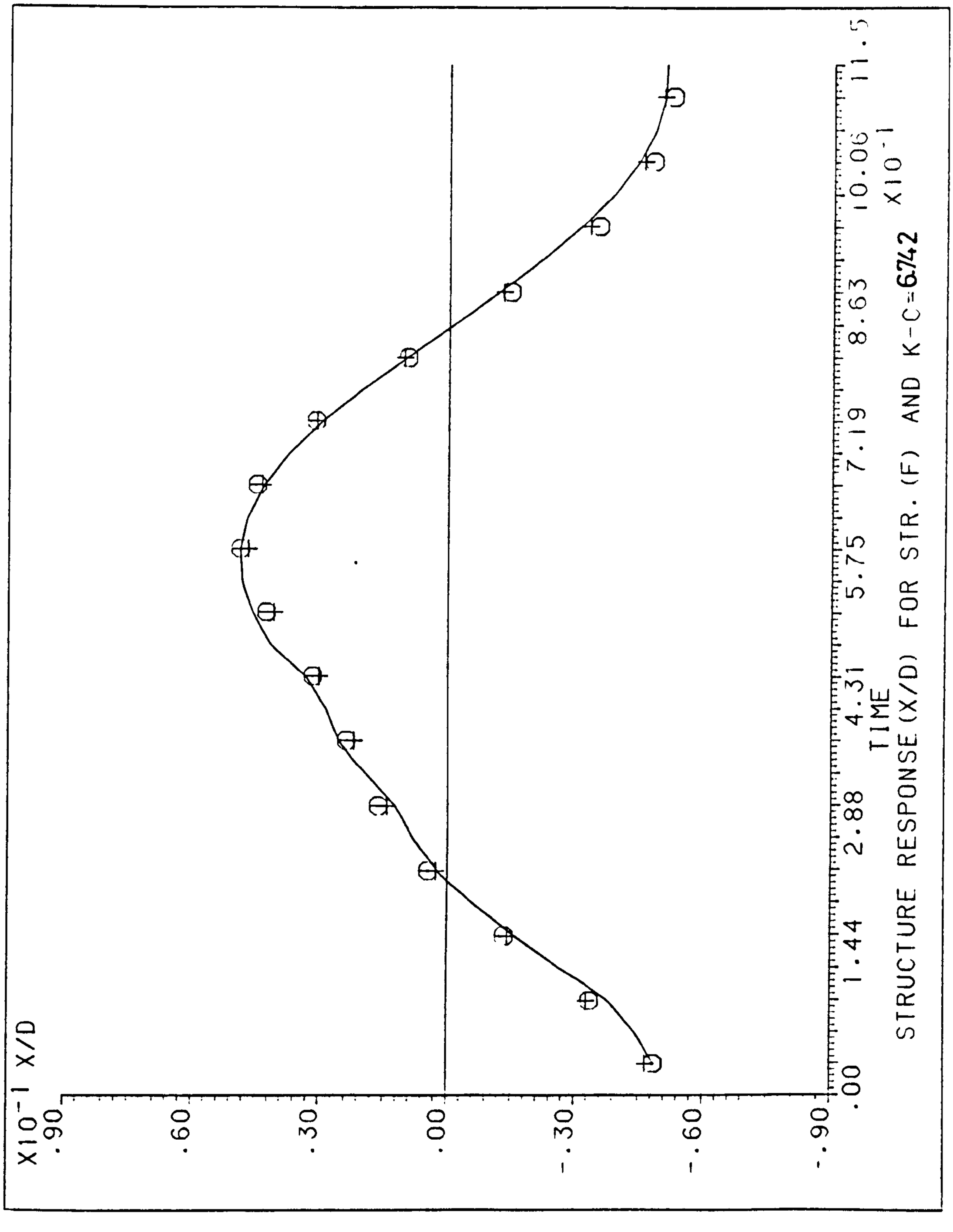


Figure (47D) - Structure's Response (X/D) during the Wave Cycle for Structure F with no Mass at Top (K-C = 6.742).

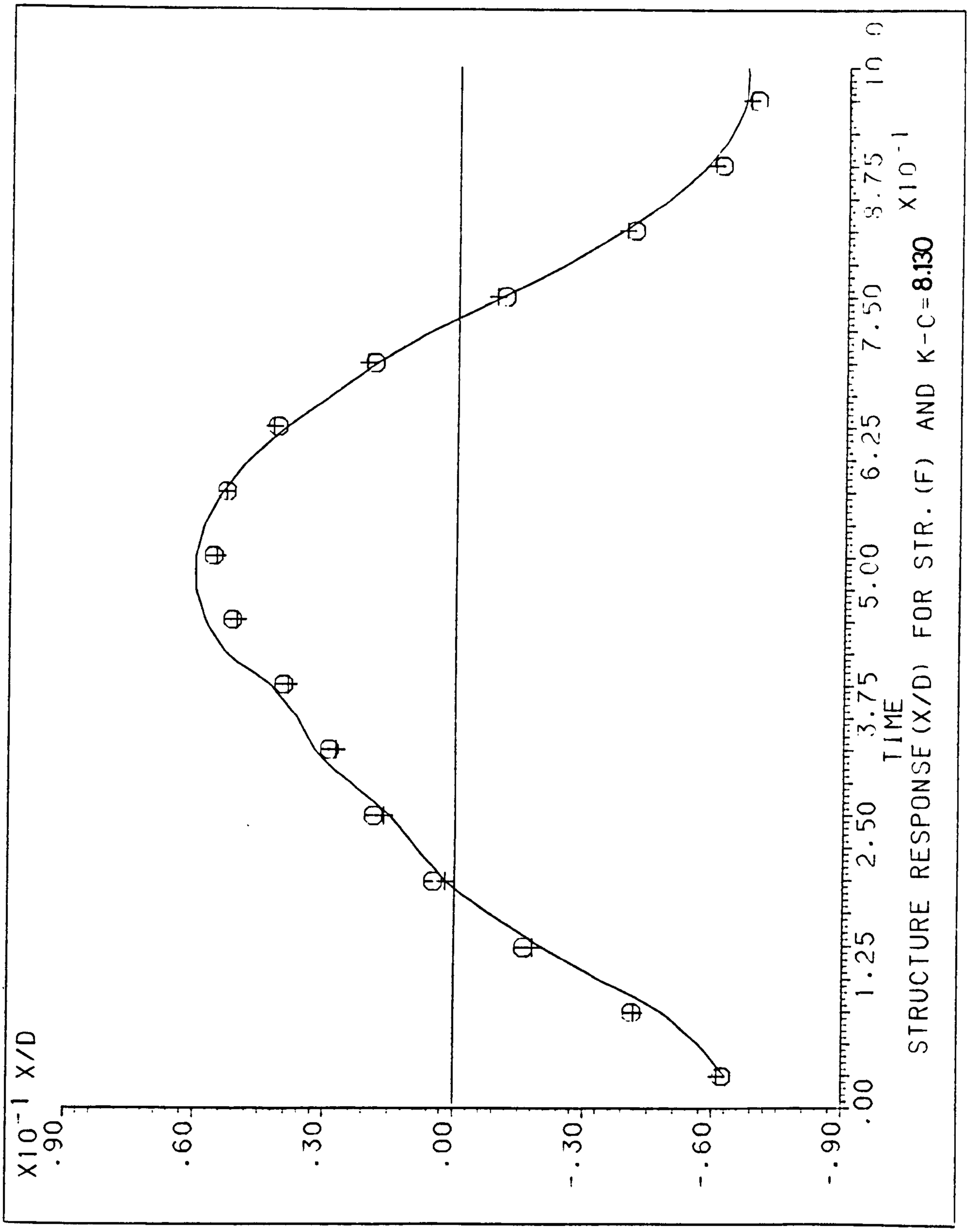


Figure (48D) - Structure's Response (X/D) during the Wave Cycle for Structure F with no Mass at Top (K-C = 8.130).

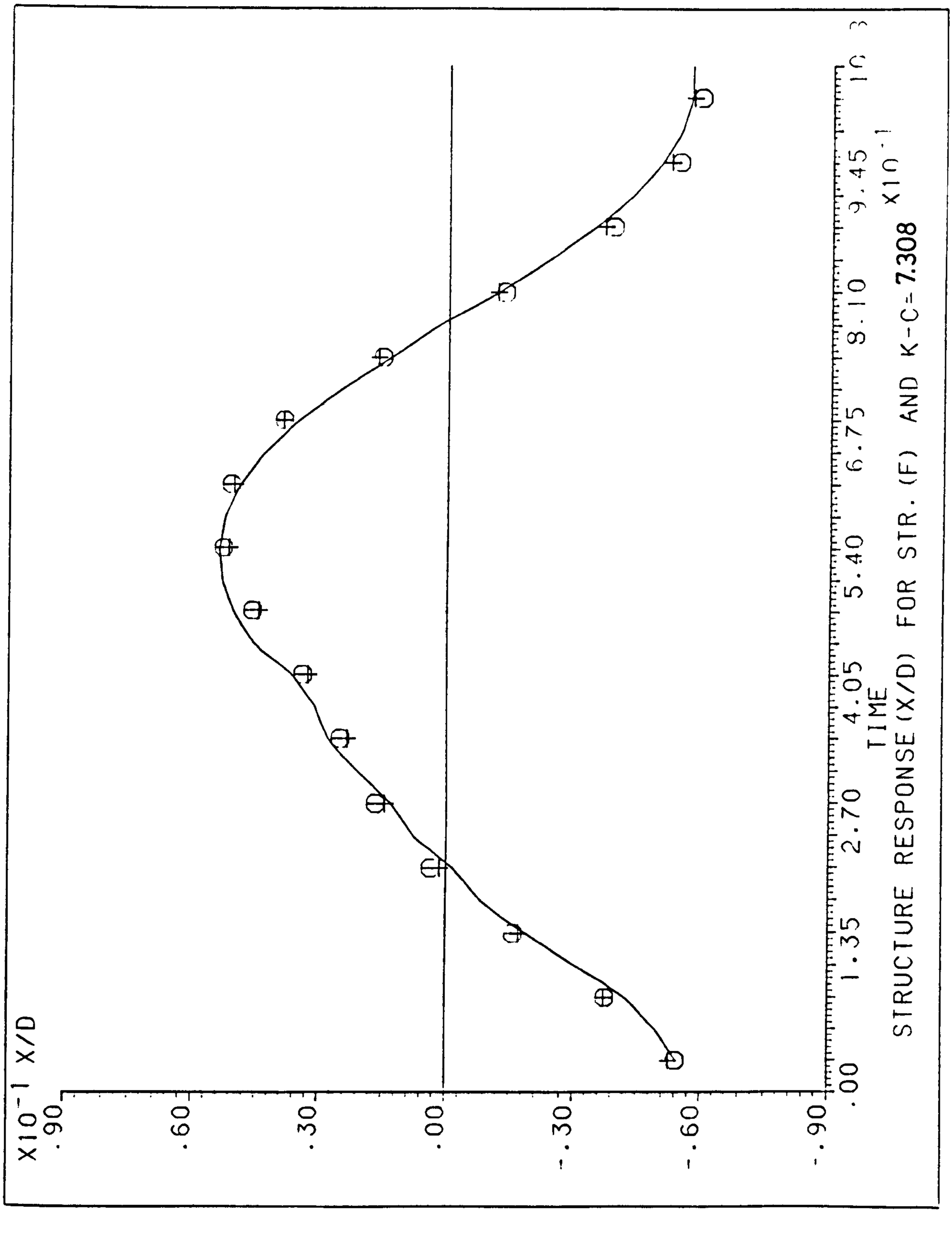


Figure (49D) - Structure's Response (X/D) during the Wave Cycle for Structure F with no Mass at Top (K-C = 7.308).

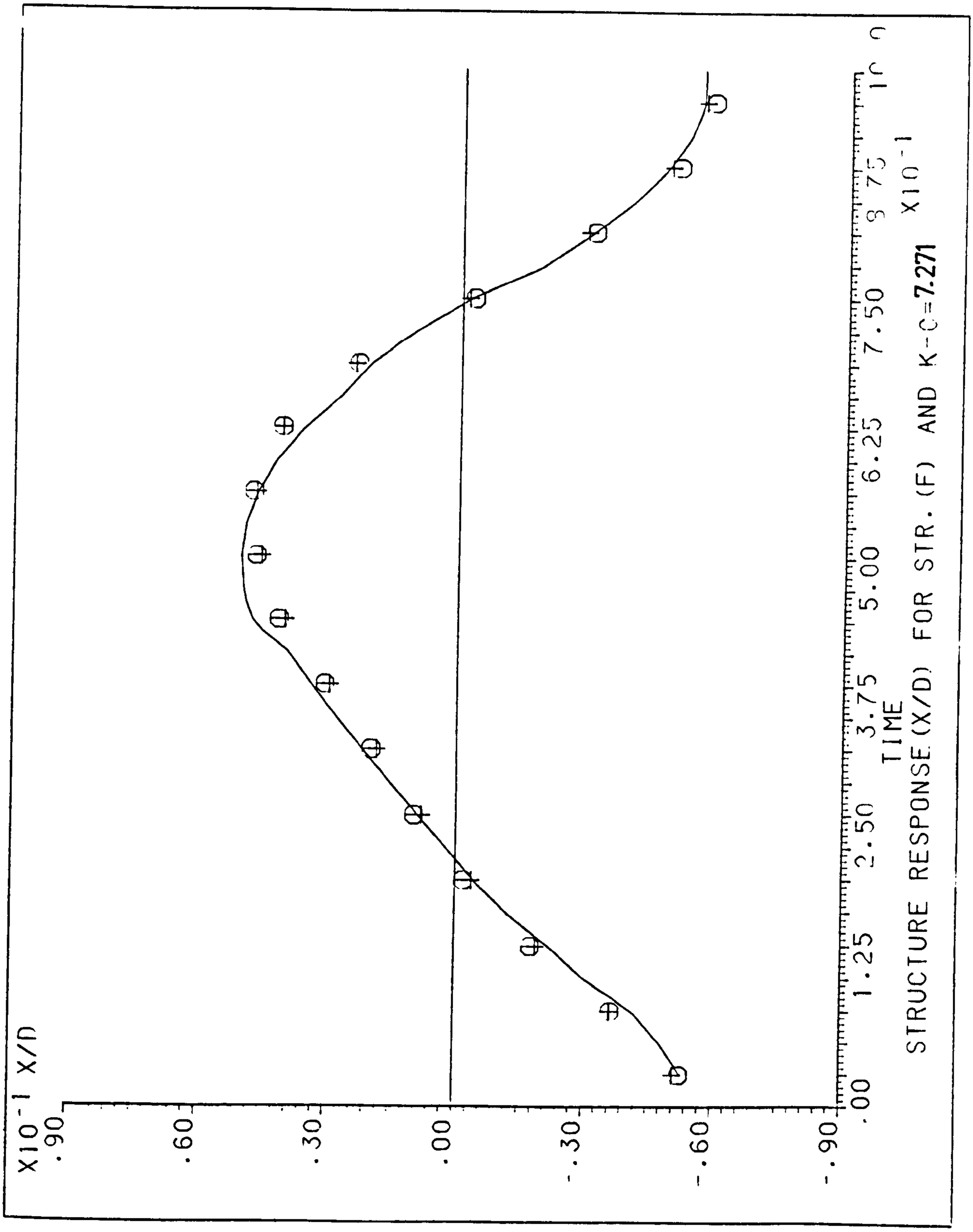


Figure (50D) - Structure's Response (X/D) during the Wave Cycle for Structure F with no Mass at Top (K-C = 7.271).

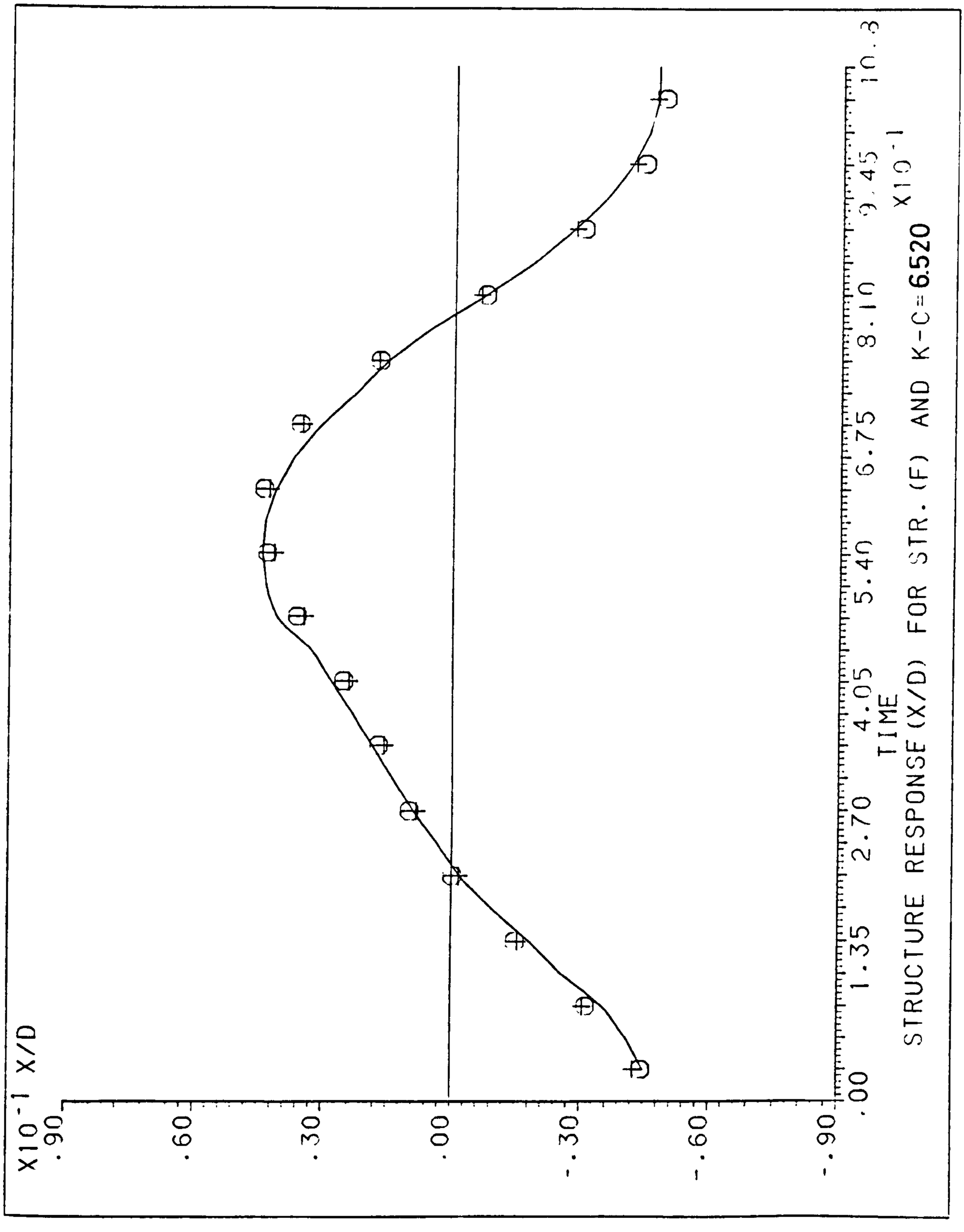


Figure (51D) - Structure's Response (X/D) during the Wave Cycle for Structure F with no Mass at Top (K-C = 6.520).

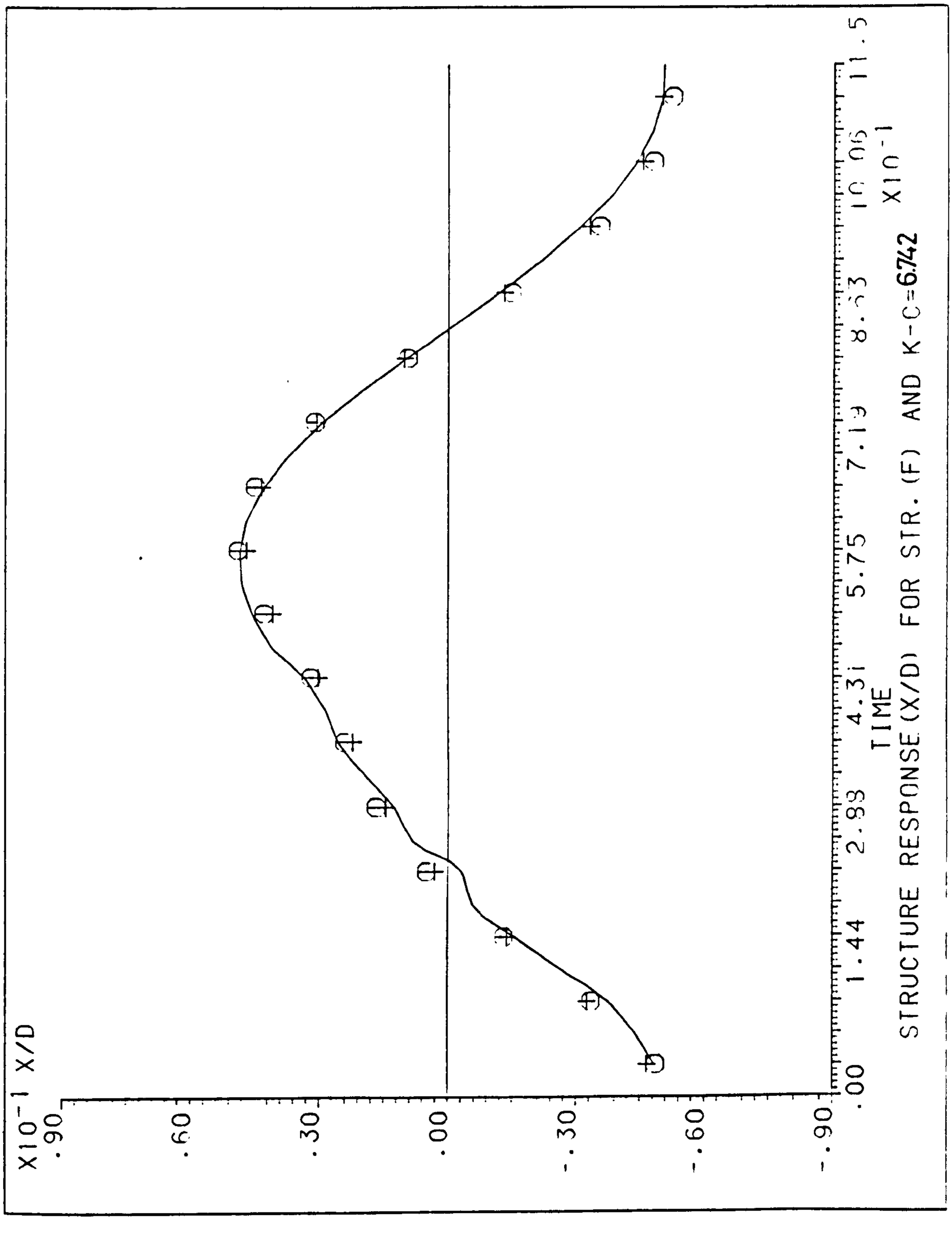


Figure (52D) - Structure's Response (X/D) during the Wave Cycle for Structure F with 1010 Grs at Top (K-C = 6.742).

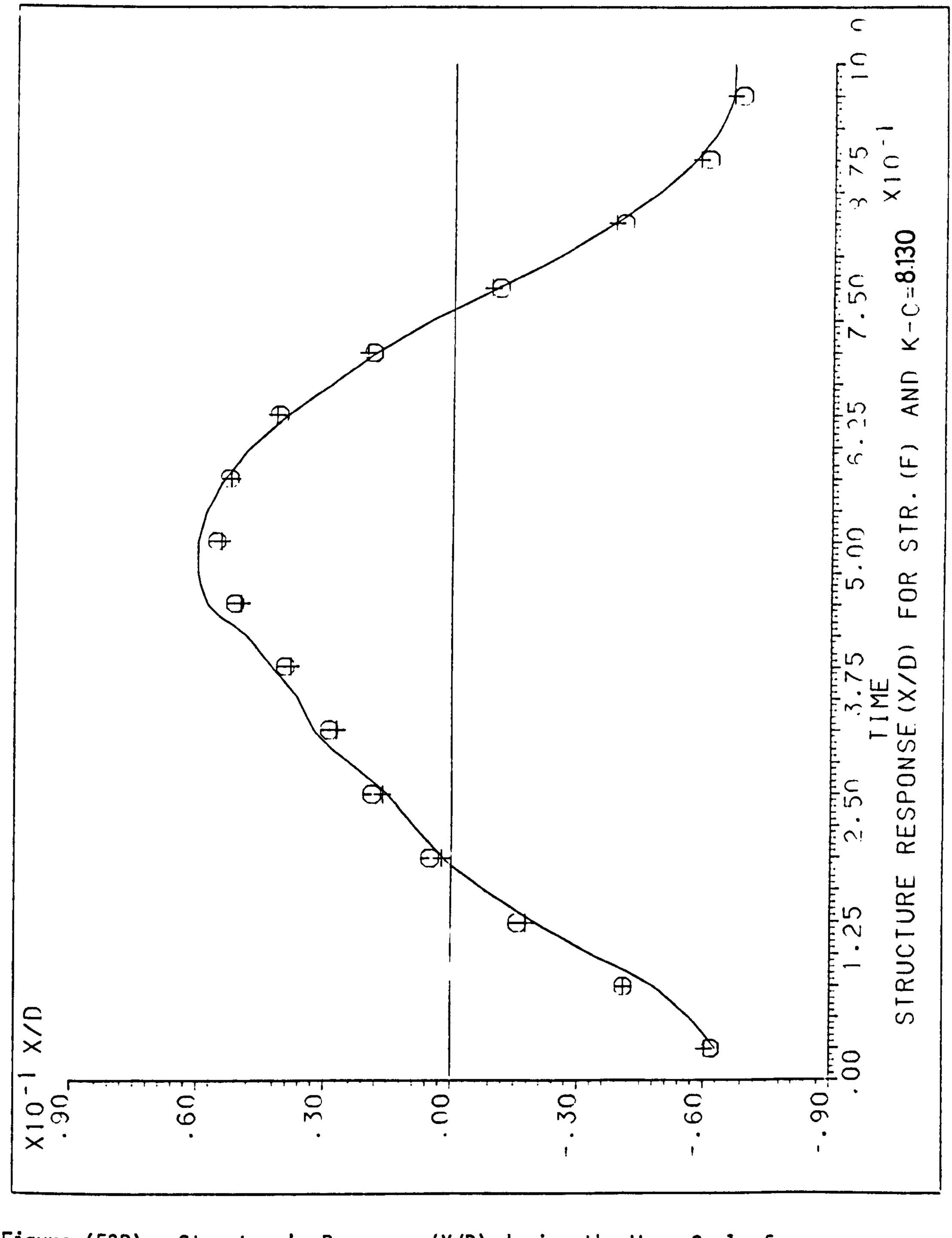


Figure (53D) - Structure's Response (X/D) during the Wave Cycle for Structure F with 1010 Grs at Top (K-C = 8.130).

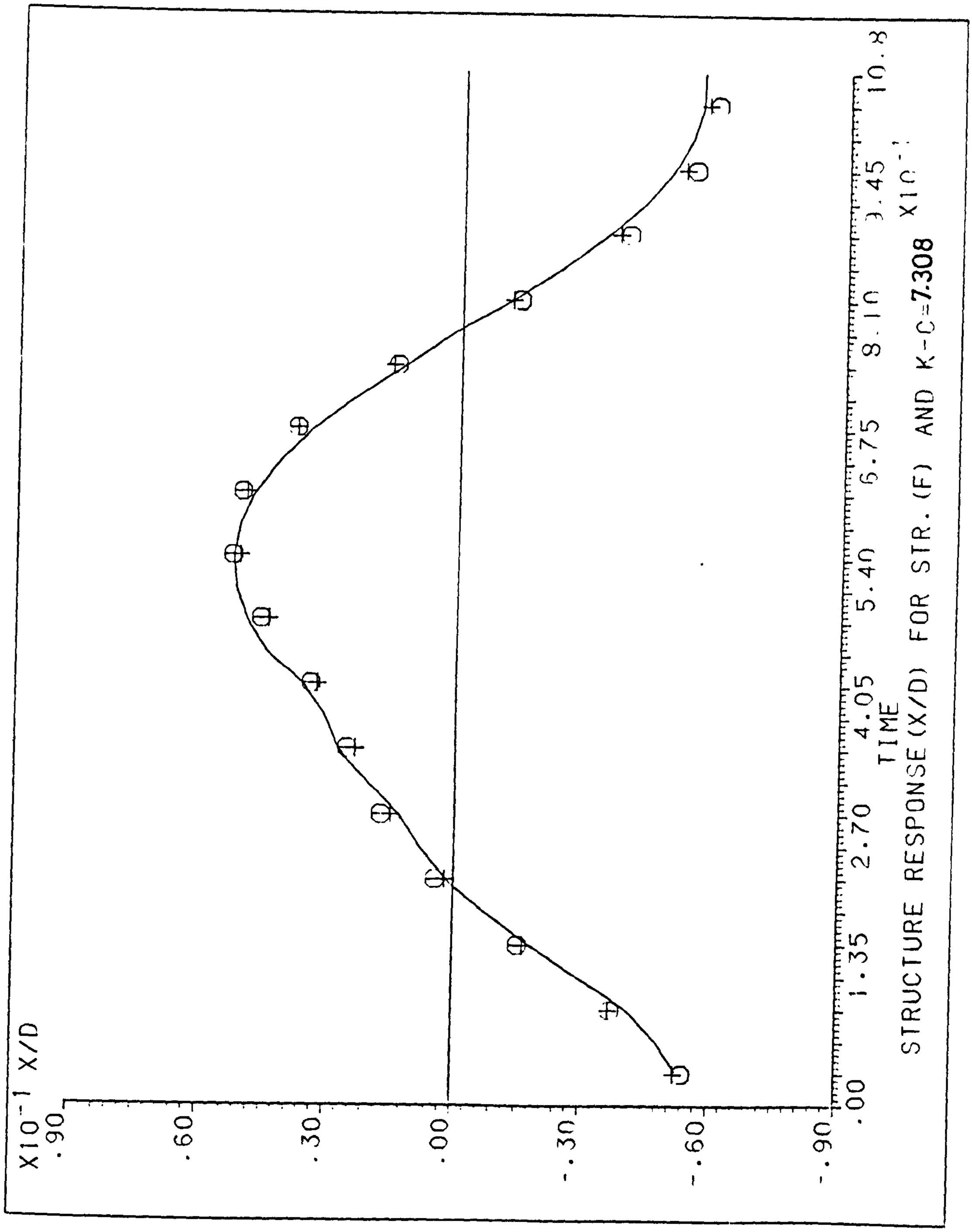


Figure (54D) - Structure's Response (X/D) during the Wave Cycle for Structure F with 1010 Grs at Top (K-C = 7.308).

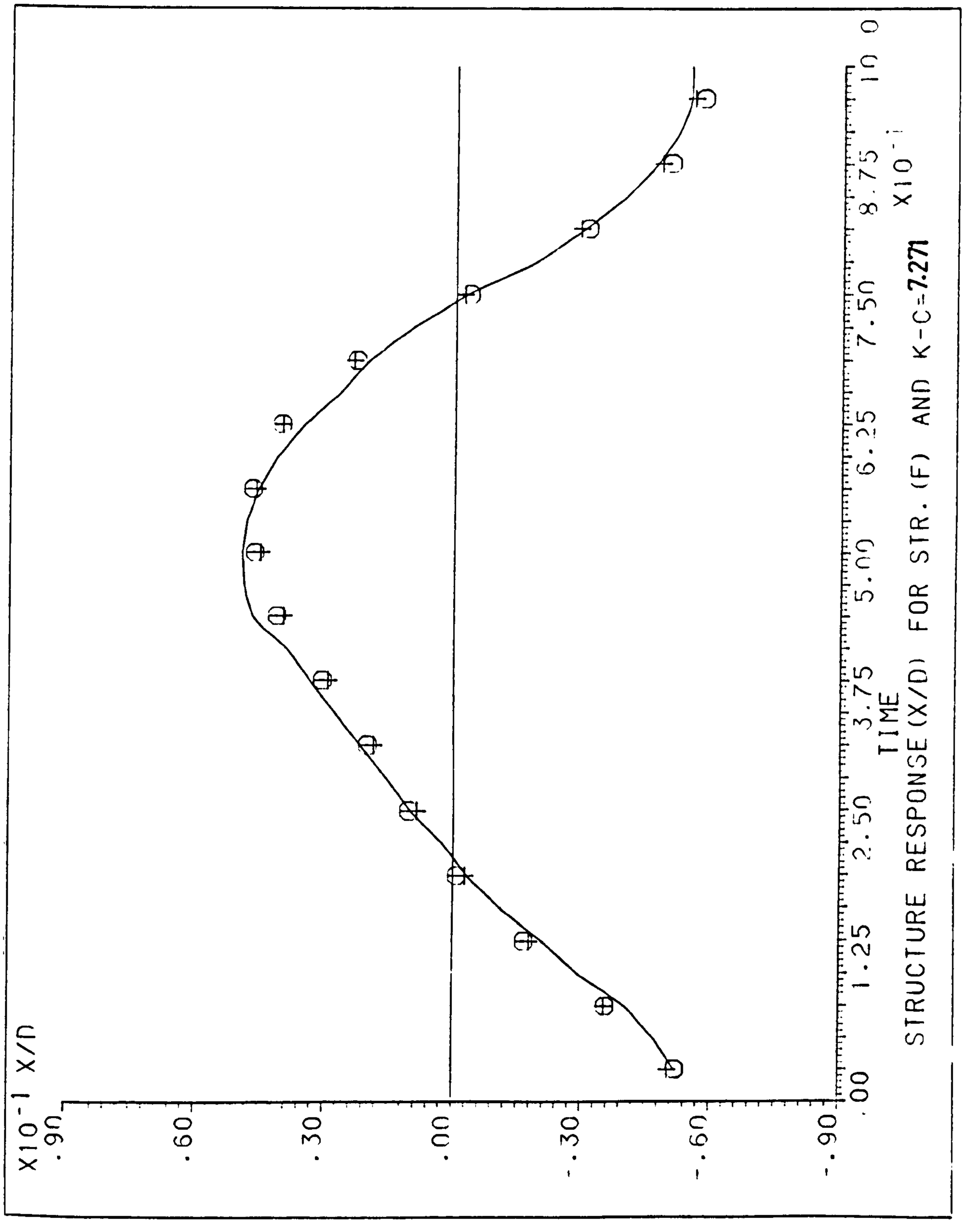


Figure (55D) - Structure's Response (X/D) during the Wave Cycle for Structure F with 1010 Grs at Top (K-C = 7.271).

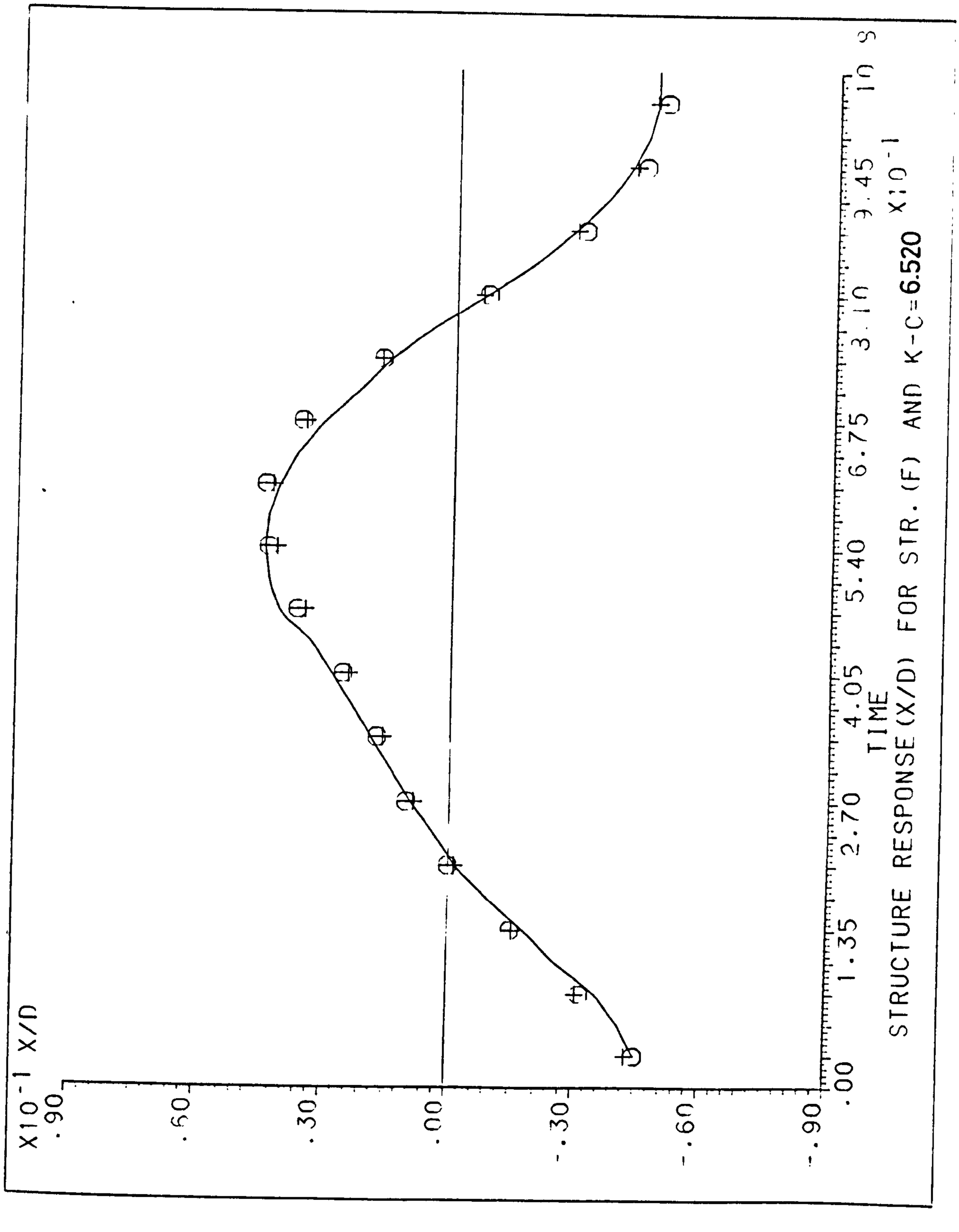


Figure (56D) - Structure's Response (X/D) during the Wave Cycle for Structure F with 1010 Grs at Top (K-C = 6.520).

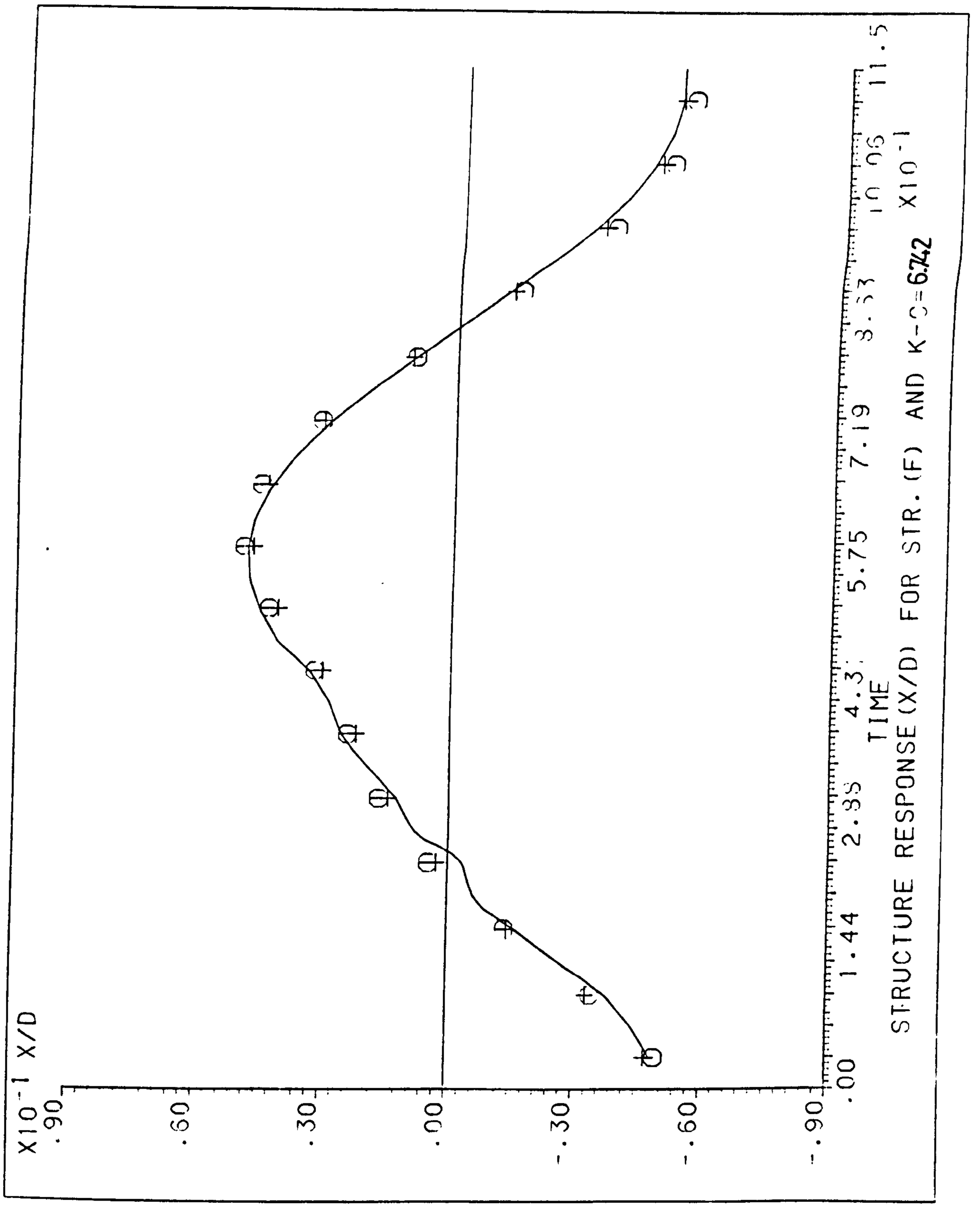


Figure (57D) - Structure's Response (X/D) during the Wave Cycle for Structure F with 2020 Grs at Top (K-C = 6.742).

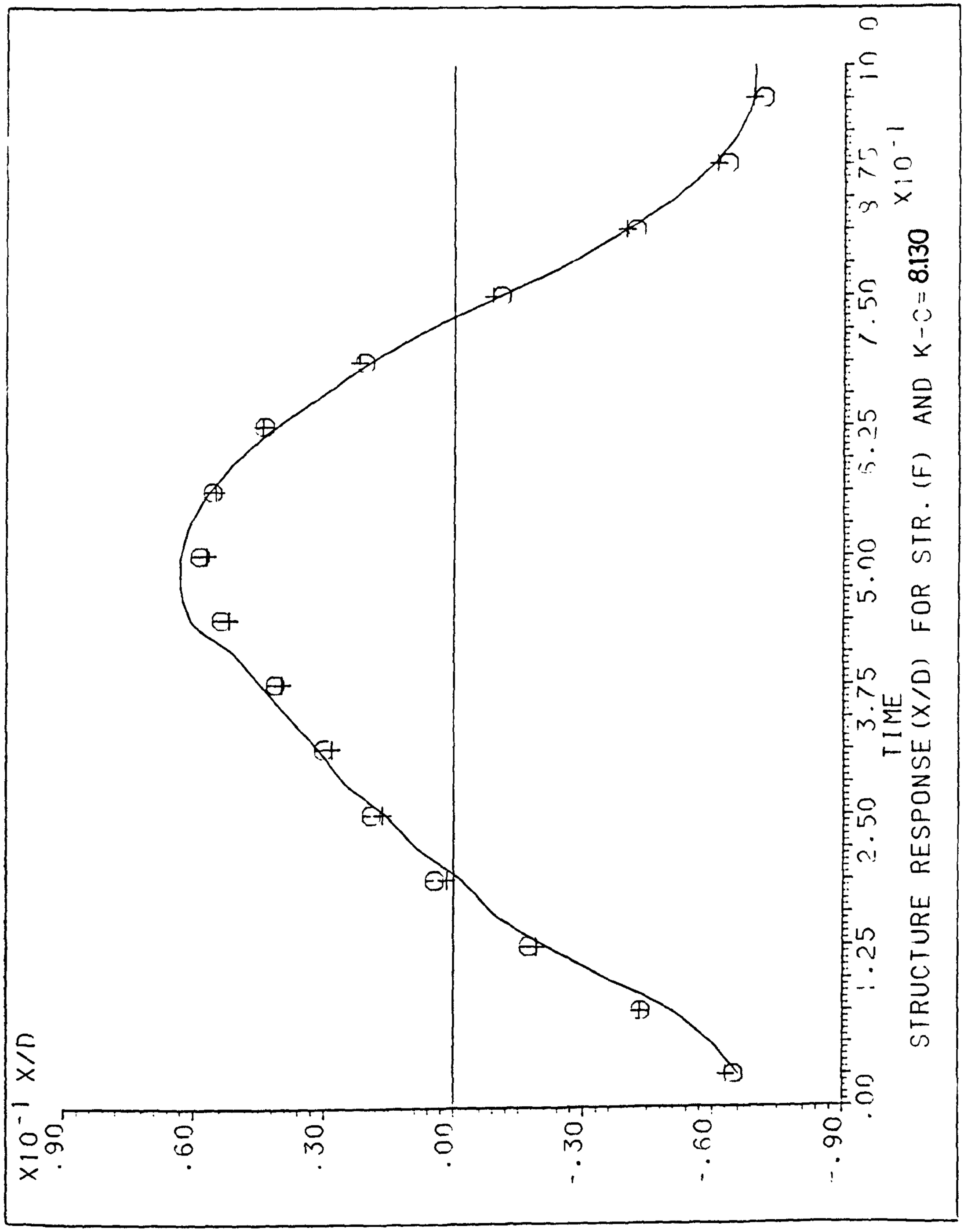


Figure (58D) - Structure's Response (X/D) during the Wave Cycle for Structure F with 2020 Grs at Top (K-C = 8.130).

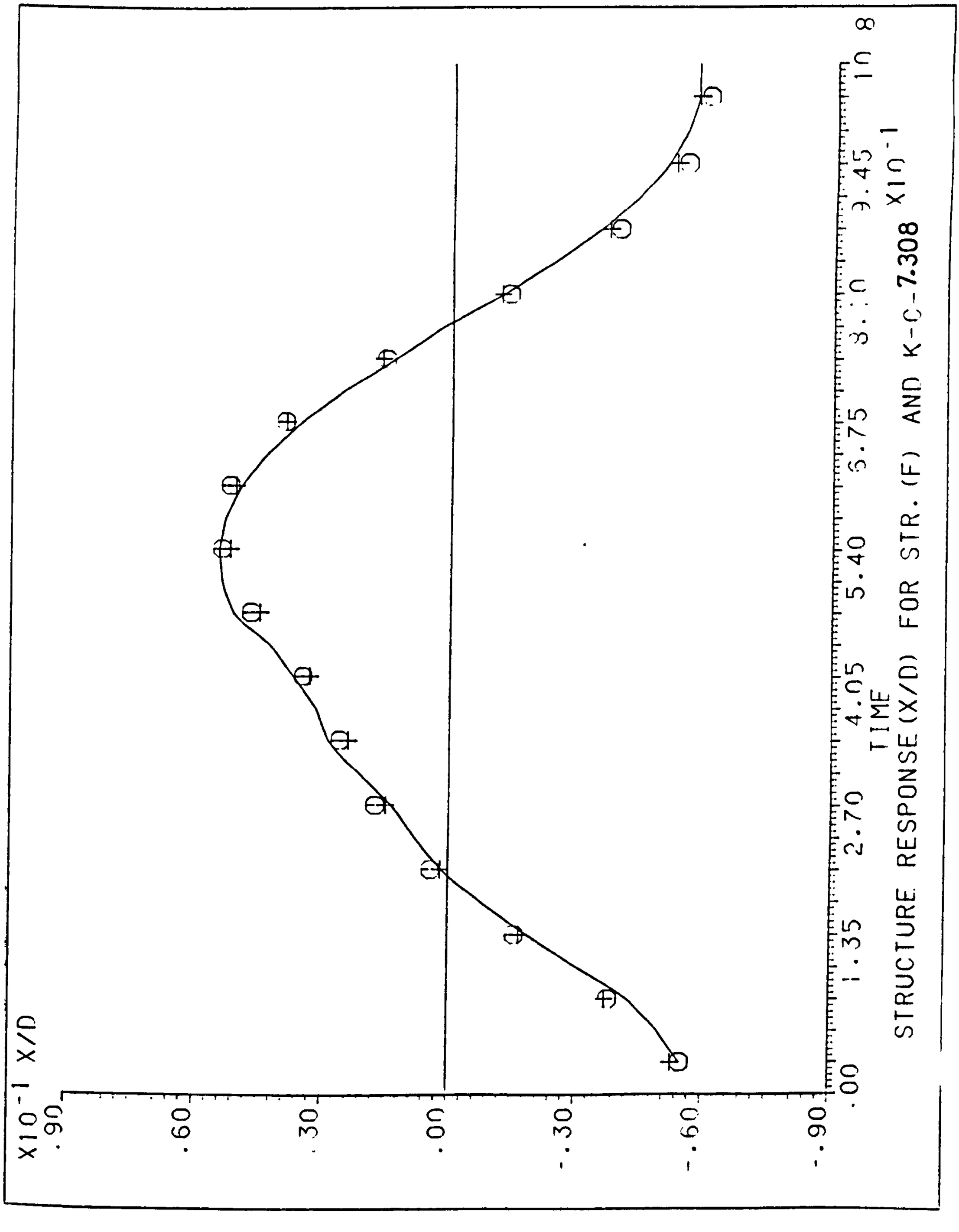


Figure (59D) - Structure's Response (X/D) during the Wave Cycle for Structure F with 2020 Grs at Top (K-C = 7.308).

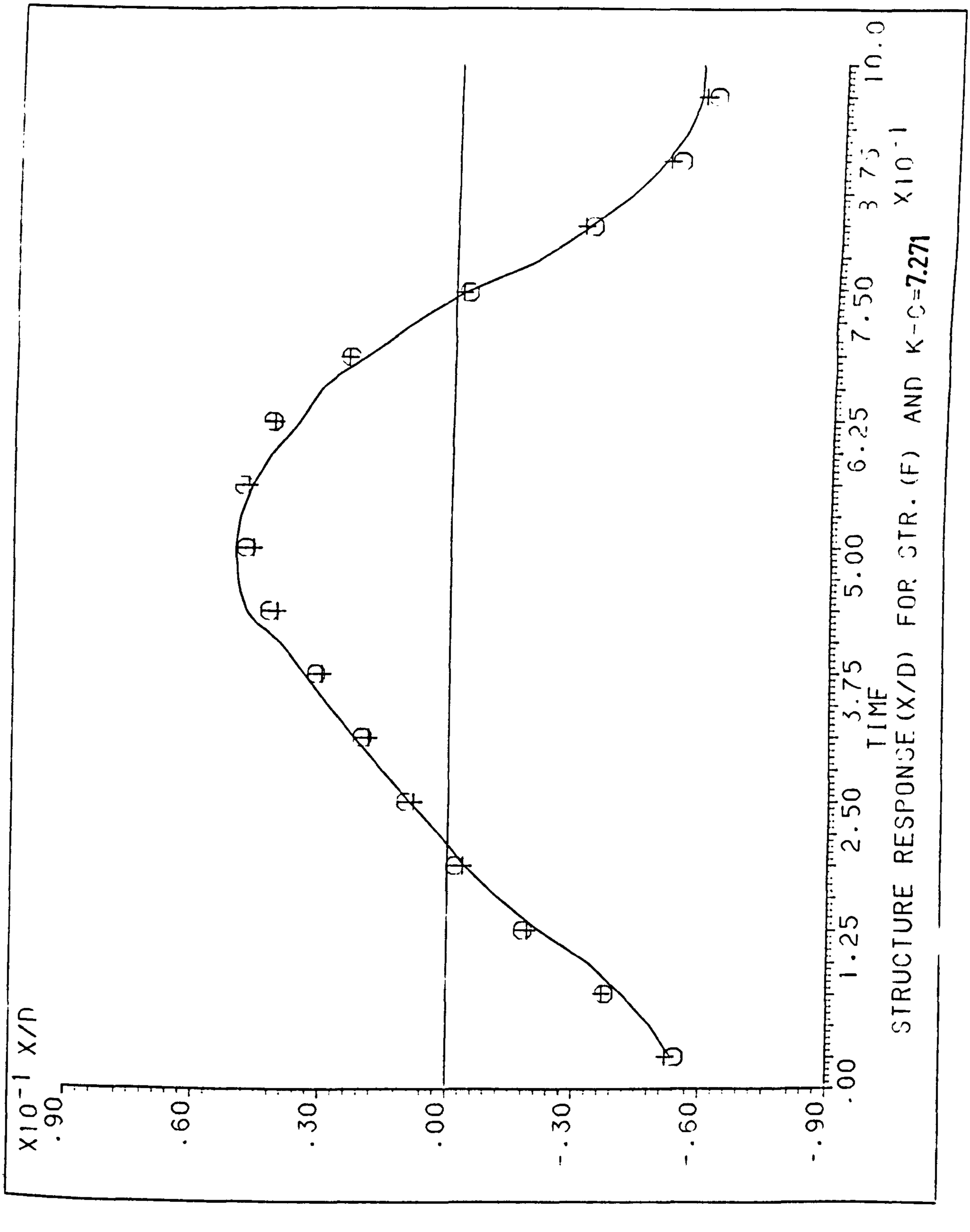


Figure (60D) - Structure's Response (X/D) during the Wave Cycle for Structure F with 2020 Grs at Top (K-C = 7.271).

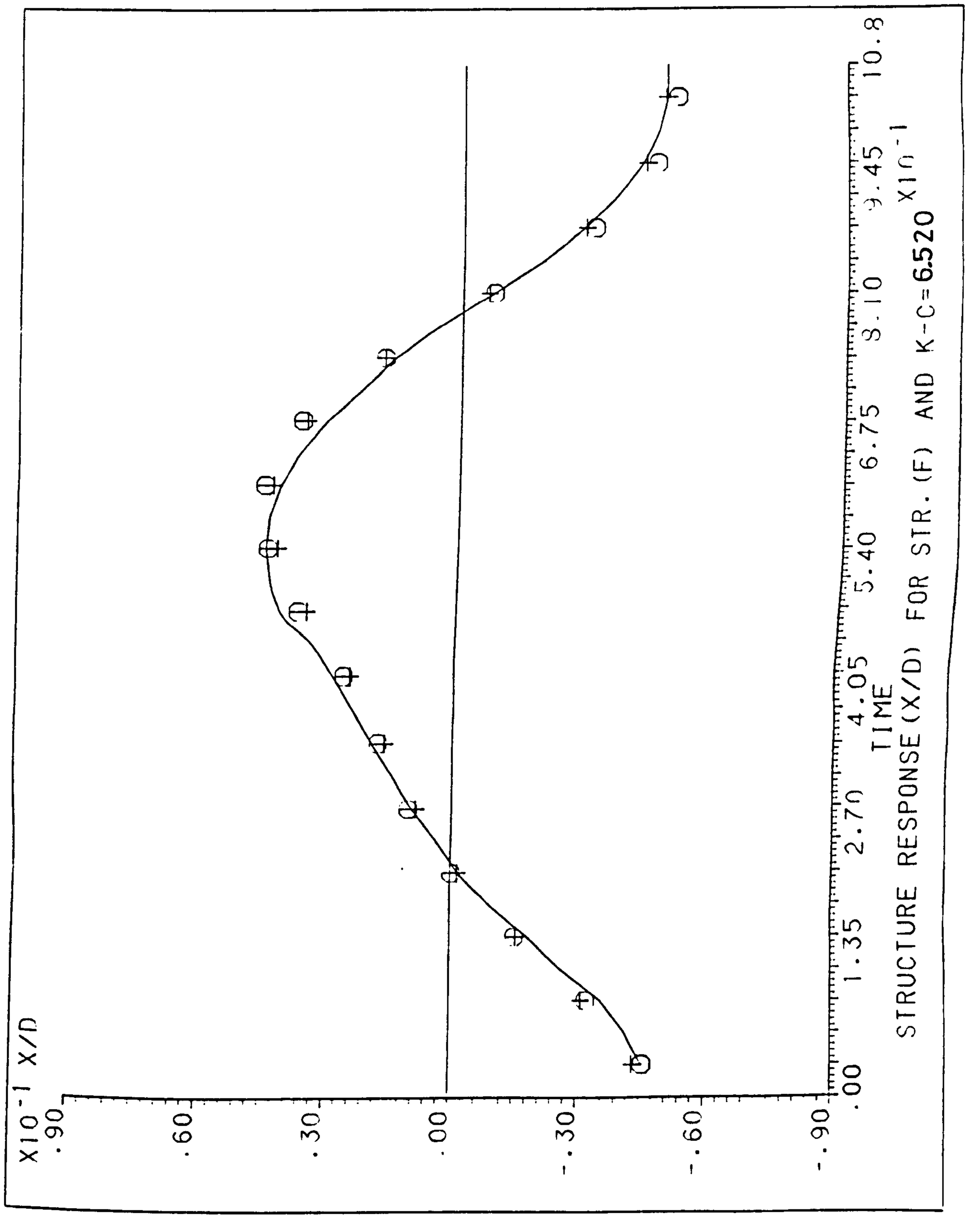


Figure (61D) - Structure's Response (X/D) during the Wave Cycle for Structure F with 2020 Grs at Top (K-C = 6.520).

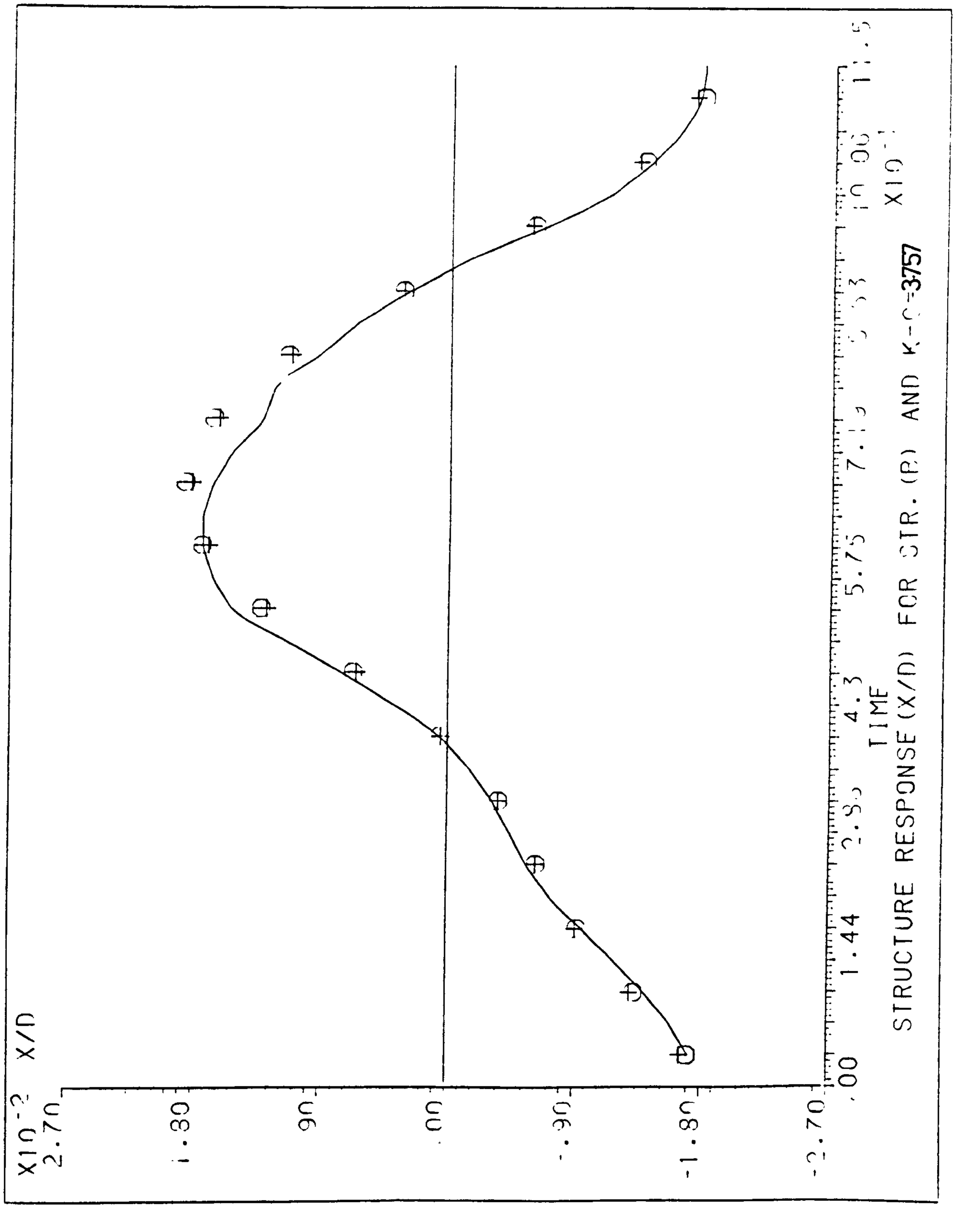


Figure (5.3.24) - Structure's Response (X/D) during the Wave Cycle for Structure B.

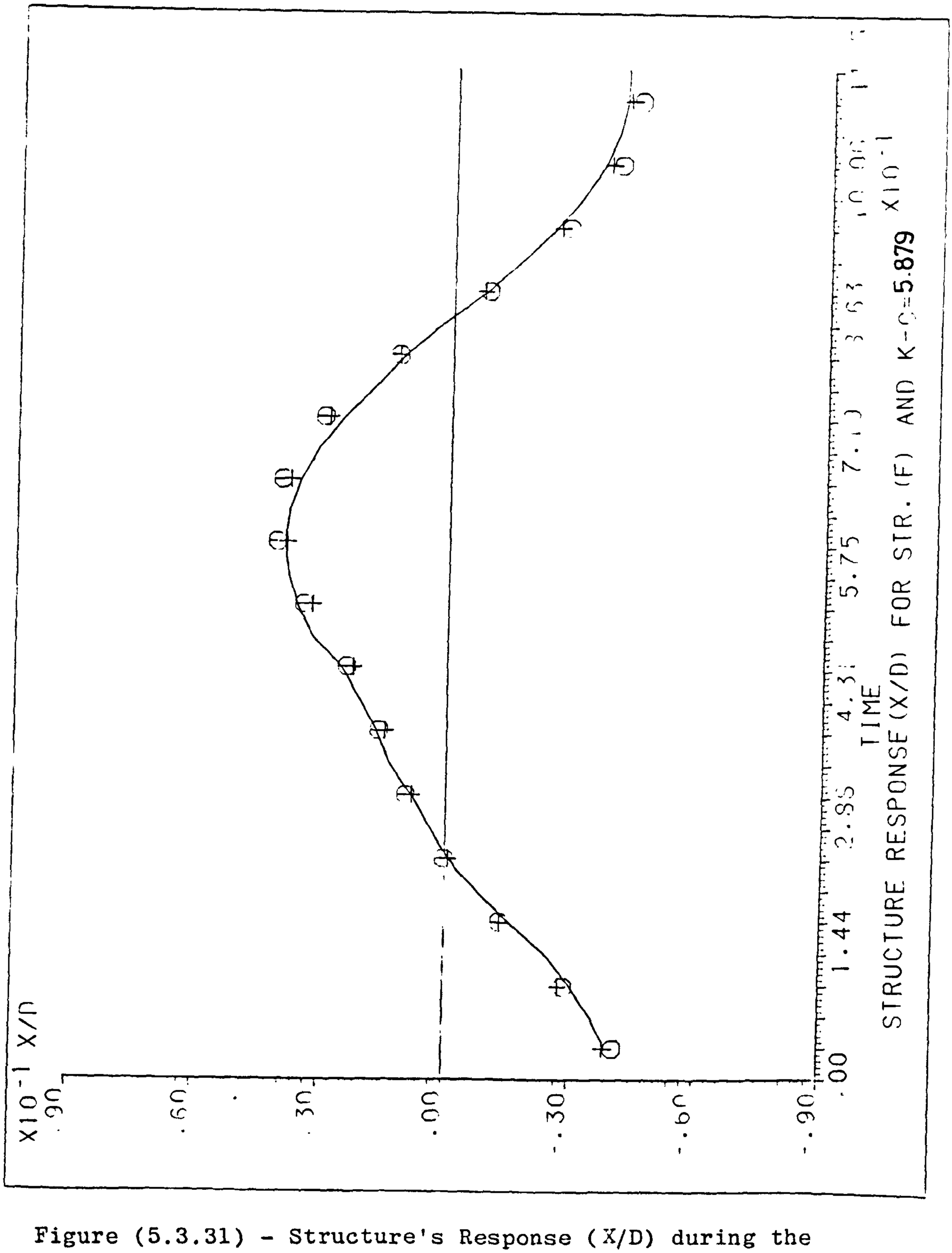


Figure (5.3.31) - Structure's Response (X/D) during the Wave Cycle for Structure F with 1010 Grs at Top.

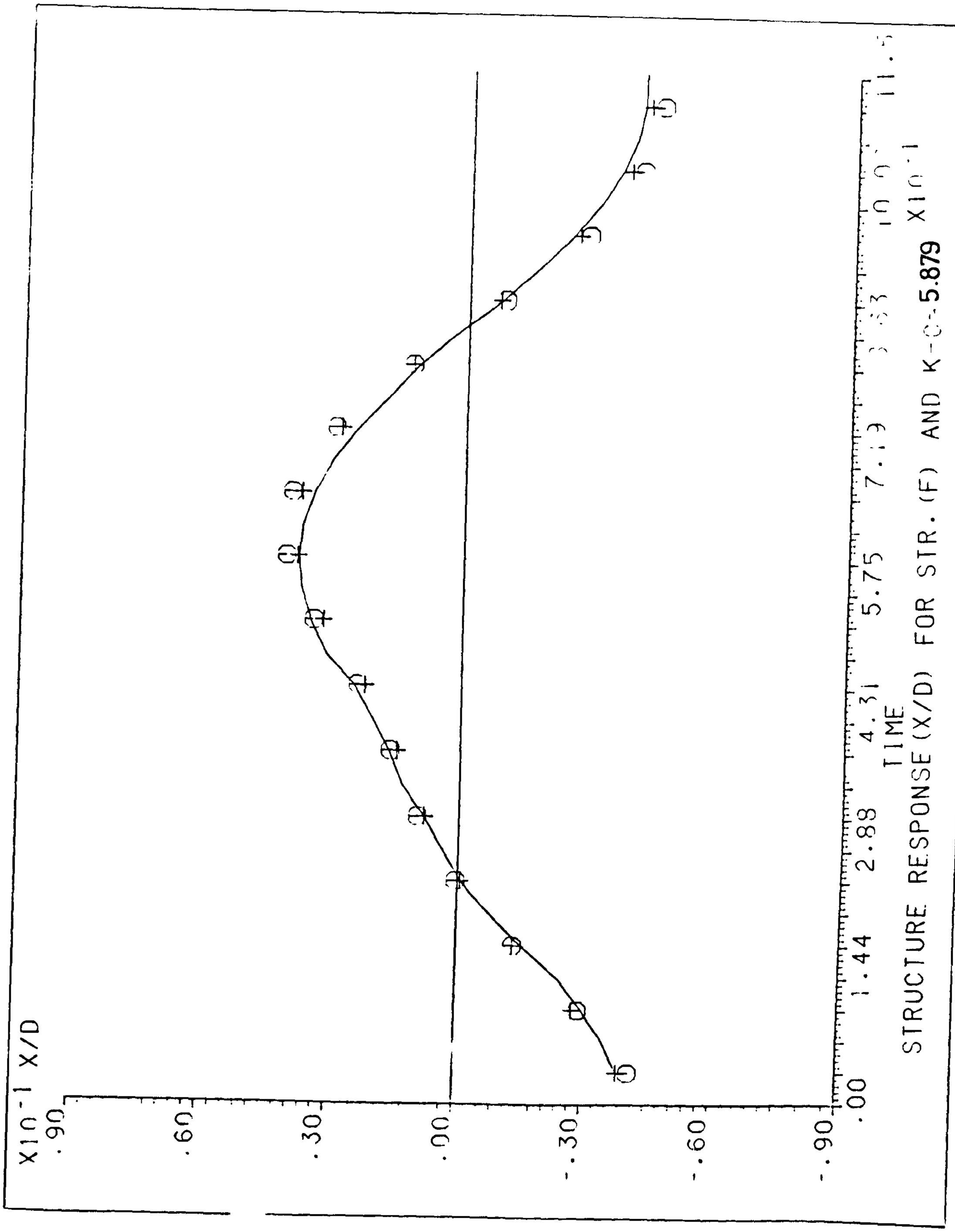


Figure (5.3.32) - Structure's Response (X/D) during the Wave Cycle for Structure F with 2020 Grs at Top.

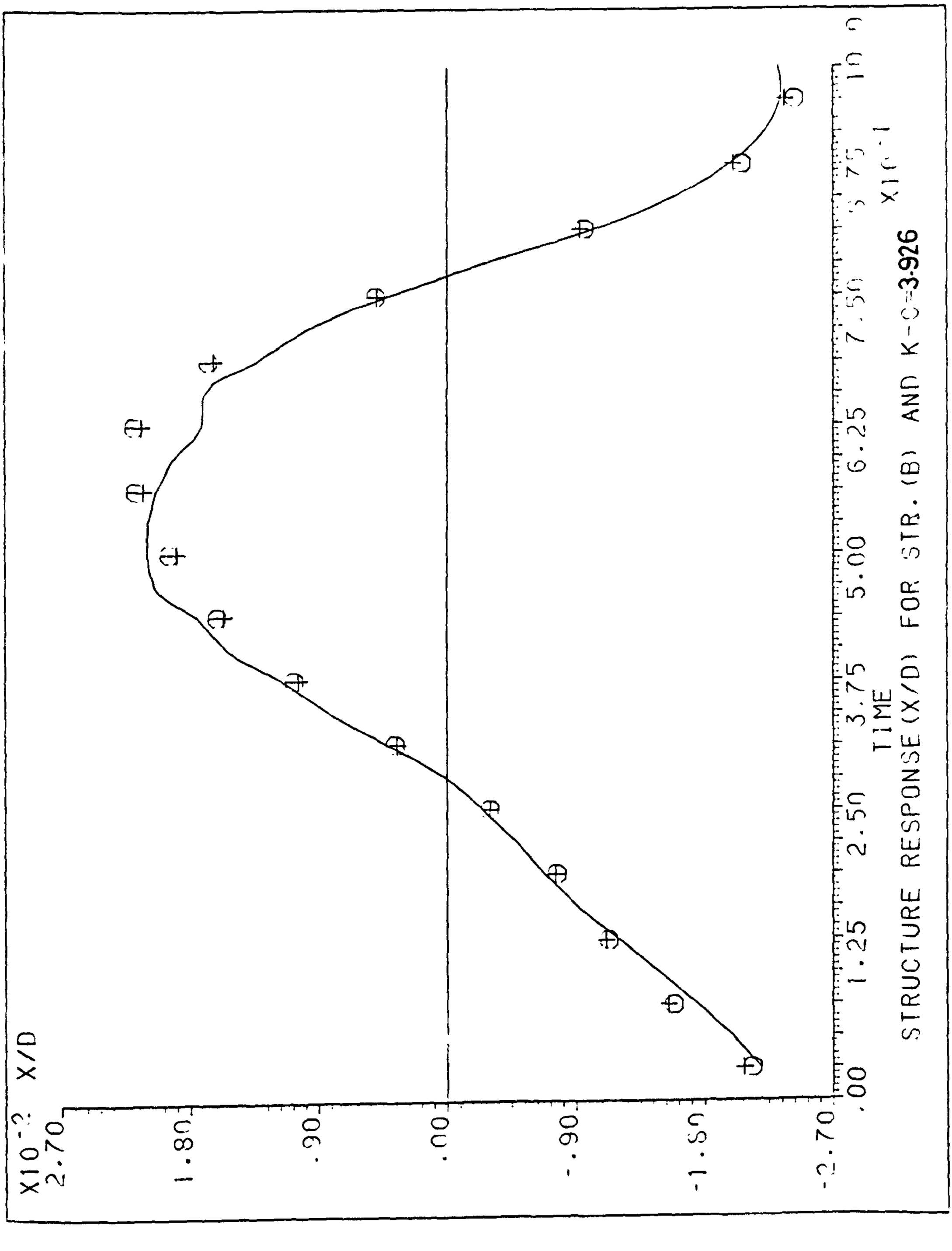


Figure (16D) - Structure's Response (X/D) during the Wave Cycle for Structure B (K-C = 3.926).